

## High energy storage density with ultra-high efficiency and fast charging–discharging capability of sodium bismuth niobate lead-free ceramics

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Received 25 March 2021; Revised 17 April 2021; Accepted 3 June 2021; Published 8 July 2021

Ceramics-based capacitors with excellent energy storage characteristics, fast charging/discharge rate, and high efficiency have received significant attention. In this work,  $\text{Na}_{0.73}\text{Bi}_{0.09}\text{NbO}_3$  (NBN) ceramics were processed through solid-state sintering route. The investigated ceramics were crystallized in a single perovskite phase. Dense microstructure, with small average grain size ( $\sim 0.92 \mu\text{m}$ ) is obtained for the investigated ceramics. A high dielectric constant  $>1000$  accompanied by a low dielectric loss was achieved for these ceramics at ambient temperature. A recoverable energy density  $\sim 0.92 \text{ J/cm}^3$  and ultra-high efficiency of 96.33% at 138 kV/cm were obtained at room temperature. Furthermore, a lower discharging time of 0.14  $\mu\text{s}$  was also achieved. This material is a suitable candidate for power pulsed applications.

**Keywords:** Polarization; energy density; capacitor; efficiency.

### 1. Introduction

Increased utilization of renewable energy requires improvement in advanced dielectric capacitors' efficiency and energy storage characteristics to broaden its application area. Ceramic dielectric capacitors possess high dielectric constant, fast charging and discharging capabilities, long life cycle, and good mechanical and thermal stabilities at high-temperatures.<sup>1–4</sup> The critical parameters by which the dielectric ceramics are evaluated for energy storage capability include high charged energy density ( $W_s$ ), high recoverable energy density ( $W_{\text{rec}}$ ), a high efficiency ( $\eta$ ), and a low discharge rate. The parameters' values are computed from the respective polarization versus electric field curve, i.e., the P–E loop via the following equations.<sup>5–7</sup>

$$W_s = \int_0^{P_{\text{max}}} E dP, \quad (1)$$

$$W_{\text{rec}} = \int_{P_r}^{P_{\text{max}}} E dP, \quad (2)$$

$$\eta = \frac{W_{\text{rec}}}{W_s} \times 100. \quad (3)$$

Here,  $P_{\text{max}}$  demonstrates the extreme (maximum) polarization,  $P_r$  the remnant polarization, and  $E$  the electric field.

Figure 1 displays the schematic polarization versus electric field (P–E) loops of pure  $\text{NaNbO}_3$  (NN) and partial Bi-substituted NN ceramics. This shows that partial Bi substitution for Na leads to slim polarization versus electric field (P–E) loop.

In the past, some linear dielectric materials such as mica and  $\text{TiO}_2$  had been reported to exhibit very high efficiency; however, their very low dielectric constant restricted further increase in their energy storage density.<sup>8–11</sup> Therefore,

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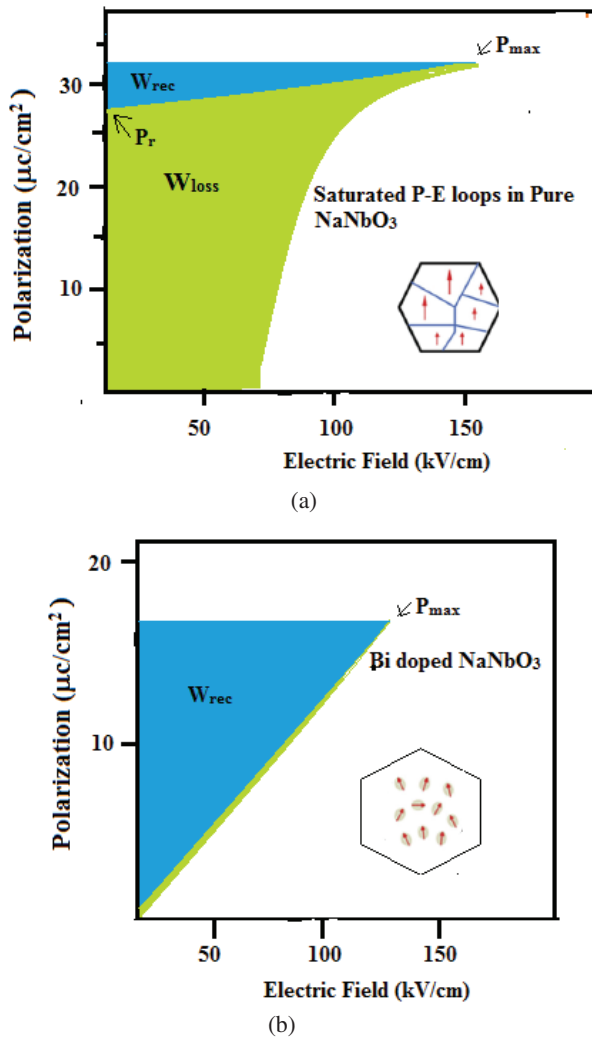


Fig. 1. Schematic of the polarization versus electric field (P-E) loops of (a) pure  $\text{NaNbO}_3$  (NN) and (b) Bi-substituted NN ceramics.

nonlinear dielectrics, such as antiferroelectrics (AFEs) and relaxor ferroelectrics (RFEs) are being investigated for their energy storage characteristics as these materials possess extensive lattice polarization results with a high dielectric constant.<sup>12–16</sup> AFEs have anti-parallel adjacent dipoles leading to a zero net polarization. These materials exhibit high  $W_s$  due to a double hysteresis loop. Moreover, AFEs possess some drawbacks, e.g., first, AFEs usually have low breakdown strength ( $\sim 10$  kV/mm). Second, due to the large hysteresis of AFEs, energy is mostly wasted as heat energy, leading to a decrease in efficiency, usually less than 80%. It increases the material's temperature due to thermal destruction.<sup>17,18</sup>

On the other hand, RFEs have a lack of long-range order of dipoles resulting in low remnant polarization ( $P_r$ ) and low coercive field ( $E_c$ ), causing slanted and slim P-E hysteresis loops.<sup>19,20</sup> These materials also possess low dielectric loss allowing these to be suitable candidate materials for high energy density applications.<sup>21</sup> The relaxor maintains a

relatively high dielectric polarization difference ( $\Delta P = P_{\max} - P_r$ ), resulting in high  $W_{\text{rec}}$  and high efficiency.<sup>22</sup>

Lead-free RFEs systems such as  $\text{BaTiO}_3$  (BT)-based,<sup>23–26</sup>  $\text{K}_{0.5}\text{Na}_{0.5}\text{NbO}_3$  (KNN)-based,<sup>27–30</sup>  $\text{AgNbO}_3$ -based,<sup>31–34</sup>  $\text{BiFeO}_3$ -based,<sup>21,35–38</sup>  $\text{Bi}_{0.5}\text{Na}_{0.5}\text{TiO}_3$  (BNT)-based,<sup>39,40</sup> have been investigated for their energy storage potential. KNN-based dielectrics are reported to possess  $W_{\text{rec}} \sim 4$  J/cm<sup>3</sup> but with low  $\eta$  values ( $< 65\%$ ).<sup>29</sup> However, the enhancement in energy storage characteristics still lags in that it is very difficult to get a high recoverable/discharge energy density ( $W_{\text{rec}}$ ) and a high efficiency ( $\eta$ ) simultaneously for a single material.

On the other hand,  $\text{NaNbO}_3$  (NN) has been investigated due to its good piezoelectric properties. NN is known to exhibit a relatively lower theoretical density (4.55 g/cm<sup>3</sup>), thus enabling light weight dielectric storage capacitors. Single crystal NN has been reported to possess AFE structure,<sup>41</sup> but its polycrystalline form displays a ferroelectric (FE) nature.<sup>42,43</sup> NN displays coexistence of FE P21ma (Q) and AFE Pbma (P) phases at room temperature. It is an easy way to enhance energy storage characteristics of NN-based materials via inducing relaxor characteristics through substitutions of suitable cations at the A site and/or B site of its perovskite structure. For example, Liu *et al.*, obtained a  $W_{\text{rec}} \sim 0.55$  J/cm<sup>3</sup> and  $\eta \sim 63\%$  for  $\text{NaNbO}_3\text{-CaZrO}_3$  system.<sup>44</sup> Similarly, in another study, a  $W_{\text{rec}}$  of 2.20 J/cm<sup>3</sup> and  $\eta$  of 82% were reported for the  $(1-x)\text{-NaNbO}_3\text{-Bi}(\text{Zn}_{0.5}\text{Ti}_{0.5})\text{O}_3$  ( $x = 0.09$ ) ceramic system.<sup>45</sup> For the  $0.90\text{NaNbO}_3\text{-}0.06\text{BaZrO}_3\text{-}0.04\text{CaZrO}_3$  ceramic system, a  $W_{\text{rec}}$  of 1.59 J/cm<sup>3</sup> and  $\eta \sim 30\%$  were obtained.<sup>46</sup> Bismuth-containing lead-free relaxors have been investigated for high energy density capacitor applications in recent years.<sup>47–52</sup>

Therefore, the energy storage characteristics of  $\text{Na}_{0.73}\text{Bi}_{0.09}\text{NbO}_3$  (NBN) ceramics were investigated for pulsed power applications in this study.

## 2. Experimental

NBN ceramics were fabricated through a solid-state mix oxide sintering process.  $\text{Na}_2\text{CO}_3$  (99.9%, Alladin China),  $\text{Bi}_2\text{O}_3$  (99%, Alladin China),  $\text{Nb}_2\text{O}_5$  (99.9%, Alladin China) were used as the initial ingredients. These materials were weighed in stoichiometric ratio and sand milling in isopropanol for 24 h with zirconia grinding balls to make a slurry. The slurry was dried at 95 °C in an electric oven overnight and then calcined for 2 h at 850 °C followed by re-ball milling for 24 h to get fine powders. The dried calcined powders were added with 5 wt% of PVA liquid solution as a binder and then grinded and sieved. The sieved powders were uniaxially pressed to make these pellets at a pressure of 100 MPa in a 12 mm diameter steel die using a hydraulic pellet press. These pellets were initially heated at 600 °C for two hours to expel the binder and finally sintered at 1275 °C for two hours. The phase constitution of the sintered samples was investigated using X-ray powder diffraction (XRD) with  $\text{CuK}\alpha$  radiation (Bruker AXS D4 Endeavor).

The dense samples' surfaces were finely polished and etched thermally at 1150 °C for 30 min, and Au coated for

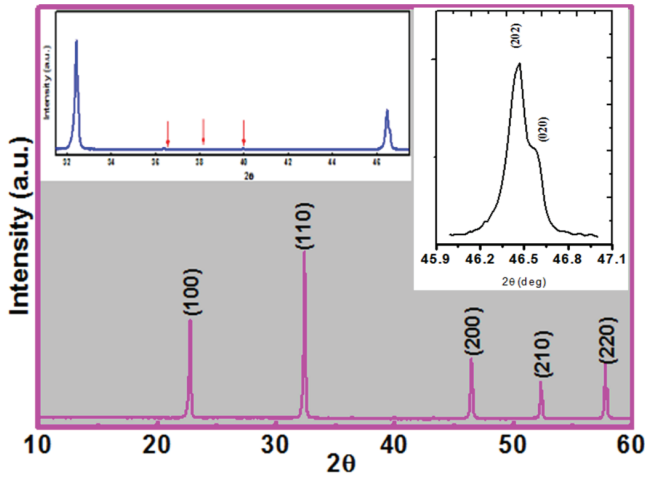


Fig. 2. XRD pattern of NBN ceramic recorded at room temperature, showing single perovskite phase crystallized into orthorhombic structure.

microstructural examination. For this purpose, a field emission scanning electron microscope (FESEM) (Jeol JXA 840A, Japan) was used.

To measure the dielectric characteristics, the opposite faces of ~0.7 mm thick sintered pellet was polished followed by silver pasting and heating at 800 °C for two hours. The dielectric loss ( $\tan\delta$ ) and dielectric constant ( $\epsilon_r$ ) of these samples were determined in the temperature range of 20–500 °C at 1 kHz, 10 kHz, 100 kHz, and 1 MHz via a computer coupled LCR meter (E4980A, Agilent made of USA).

The electric field-dependent polarization at 1 Hz for a silver pasted 0.3 mm thick, dense sample by a dielectric test system (model Premier II, Radiant, USA) at different electric fields up to its dielectric breakdown strength (DBS). The composition’s charging and discharging rates were tested with the help of a charge-discharge test system (model CPR1701–100, Ploy K) at an electric field of 43 kV/cm, 65 kV/cm, 87 kV/cm, and 108 kV/cm respectively.

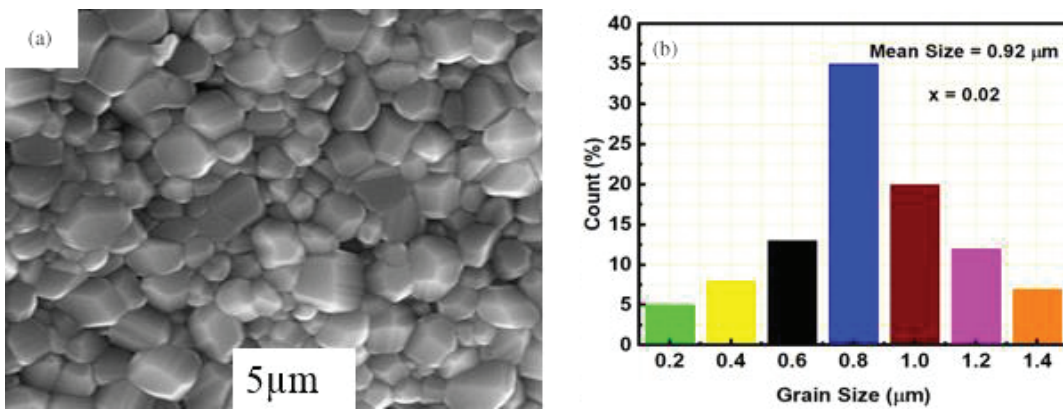


Fig. 3. (a) Scanning electron microscope image (SEM) and (b) the grain size distribution of NBN ceramics sintered at 1275 °C for two hours.

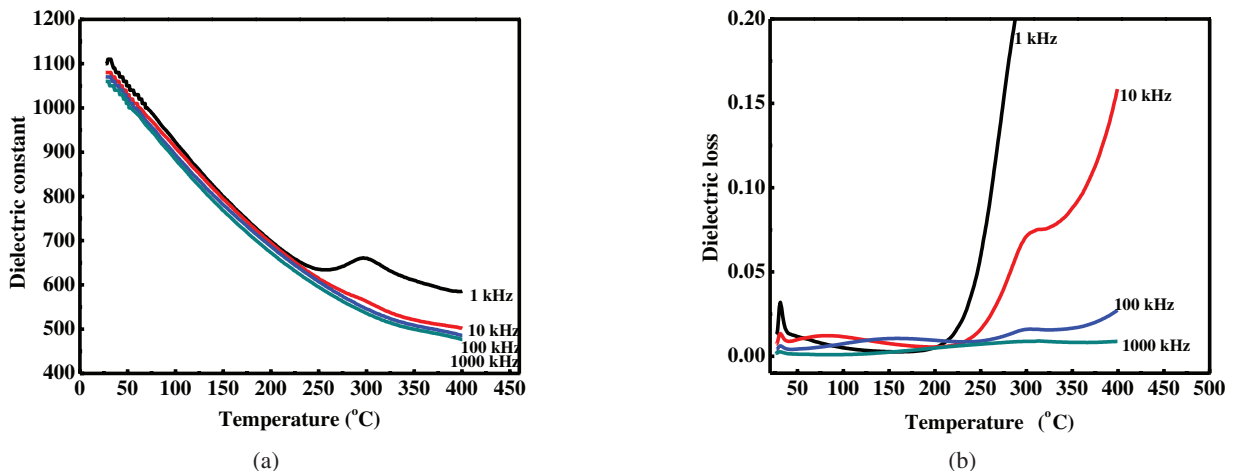


Fig. 4. Variation in (a)  $\epsilon_r$  and (b)  $\tan\delta$  of NBN ceramics with temperature.

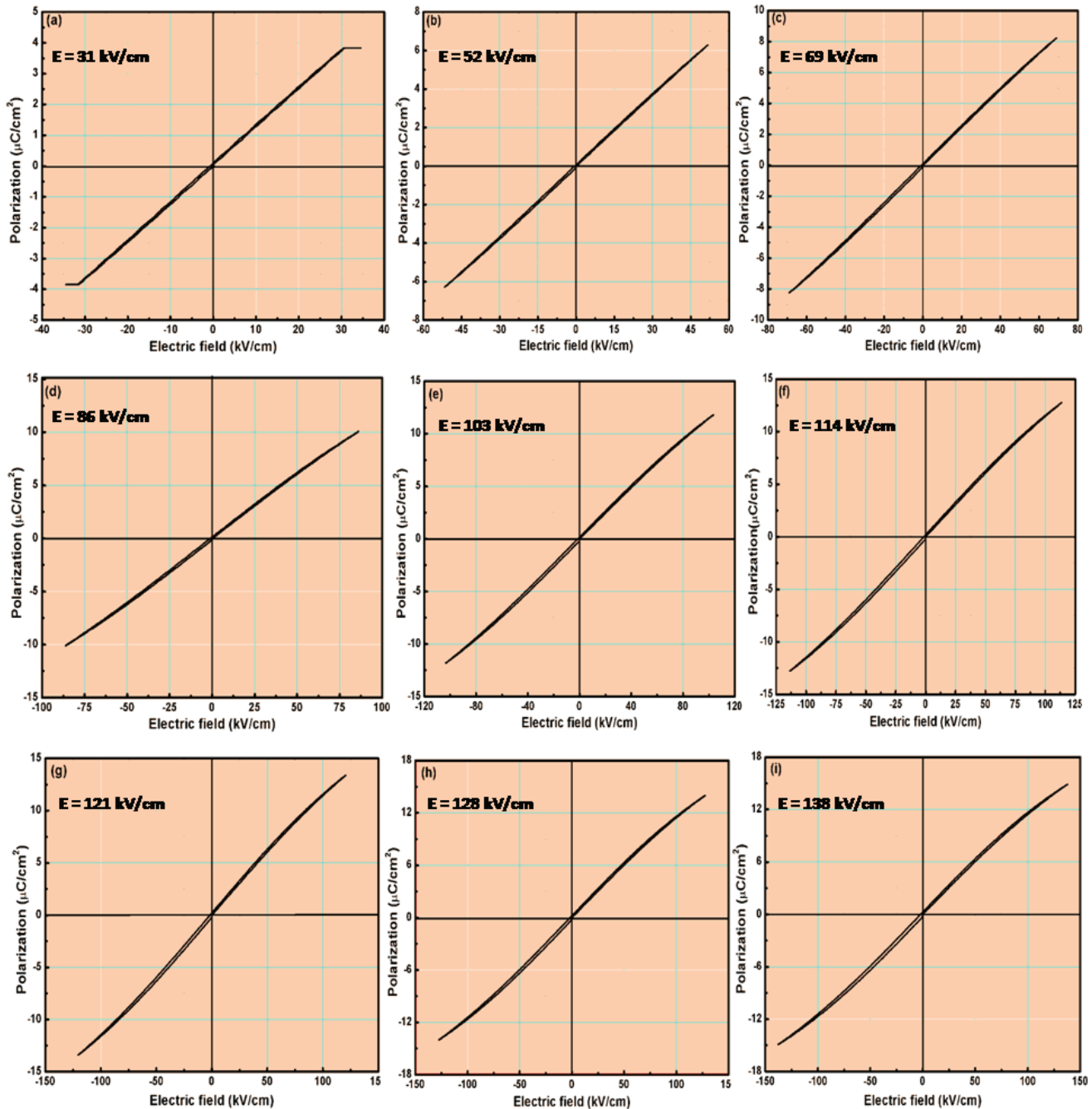


Fig. 5. (a)–(i) Room temperature P–E loops at various electric fields for NBN ceramics.

### 3. Results and Discussion

The XRD spectrum of NBN ceramic composition sintered at 1275 °C for two hours is shown in Fig. 2. Upon indexing the pattern with standard data, it could be confirmed that the sample comprised a single perovskite phase with no second phase formation within the in-house XRD detection limit. This shows that  $\text{Bi}^{3+}$  has occupied the  $\text{Na}^{+1}$  sites in the lattice of NN as its ionic radii ( $r_{\text{Bi}^{3+}} = 1.34 \text{ \AA}$ ) is matching that of the  $\text{Na}^{+1}$  ( $r_{\text{Na}^{+1}} = 1.39 \text{ \AA}$ ) for coordination no 12.<sup>53</sup> The enlarged

XRD pattern of NBN ceramic, in the  $2\theta$  range of 32.0–32.6° shows a single (110) peak at 32.4°, and a slight splitting of (200) peak at ~46.4° was also observed, suggesting orthorhombic structure of NBN ceramic system.<sup>54,55</sup> Furthermore the observation of some low intensity peaks marked by arrows around 36–40° may show orthorhombic super structure.

Figure 3(a) shows the secondary electron SEM image (SEI) of the polished and thermally-etched NBN ceramics sample sintered at 1275 °C for two hours. A highly dense

microstructure with almost no porosity can be observed in this image. Nanomeasurer analytical software was used to get grain size distribution, based on which the statistical grain size is measured to be  $\sim 0.92 \mu\text{m}$ , as shown in Fig. 3(b). It has been reported that several microstructural parameters, such as porosity, grain size, second phase formation, and other crystal defects affect the DBS of ceramics.<sup>56</sup> The grain size is considered to mainly affect the DBS of bulk ceramics.<sup>43,57</sup> The smaller grain size of the investigated ceramics is beneficial for DBS and high recoverable energy density. The relationship of DBS to grain size was studied by Tunkasiri *et al.*, and is given by the following equation<sup>57</sup>:

$$\text{DBS} \sim \frac{1}{\sqrt{G}}. \tag{4}$$

Figures 4(a) and 4(b) display the variation in dielectric constant ( $\epsilon_r$ ) and dielectric loss ( $\tan\delta$ ) of NBN ceramic with temperature at 1 kHz, 10 kHz, 100 kHz, and 1 MHz. The

$\epsilon_r$  graph shows that the investigated ceramics has a Curie temperature below  $0^\circ\text{C}$ , suggesting that NBN belongs to the group II  $\text{NaNbO}_3\text{-ABO}_3$  relaxors. In group II NN- $\text{ABO}_3$  relaxors, the dielectric maximum temperature ( $T_m$ ) rapidly decreases to a lower temperature region as the  $\text{ABO}_3$  end member's concentration increases before reaching a critical value.<sup>58,59</sup> The observed frequency dispersion in  $\epsilon_r$  of NBN ceramics may evident relaxor characteristics. The dielectric constants' values at room temperature for NBN ceramics at 1 kHz, 10 kHz, 100 kHz, and 1 MHz were measured to be  $\sim 1100, 1080, 1075,$  and  $1060,$  respectively, accompanied by a low dielectric loss. Upon increasing the temperature, the dielectric loss slightly increased but did not exceed 0.05 up to temperature range of  $200^\circ\text{C}$  at all measured frequencies. The low dielectric loss and the moderate  $\epsilon_r$  values are also promising for NBN ceramic's high DBS.<sup>4</sup>

The room temperature P-E loops of NBN ceramics sintered at  $1275^\circ\text{C}$  for two hours at various electric fields are shown in Fig. 5. NBN ceramics display almost linear P-E loops at low electric fields ( $<70 \text{ kV/cm}$ ). It exhibits a very

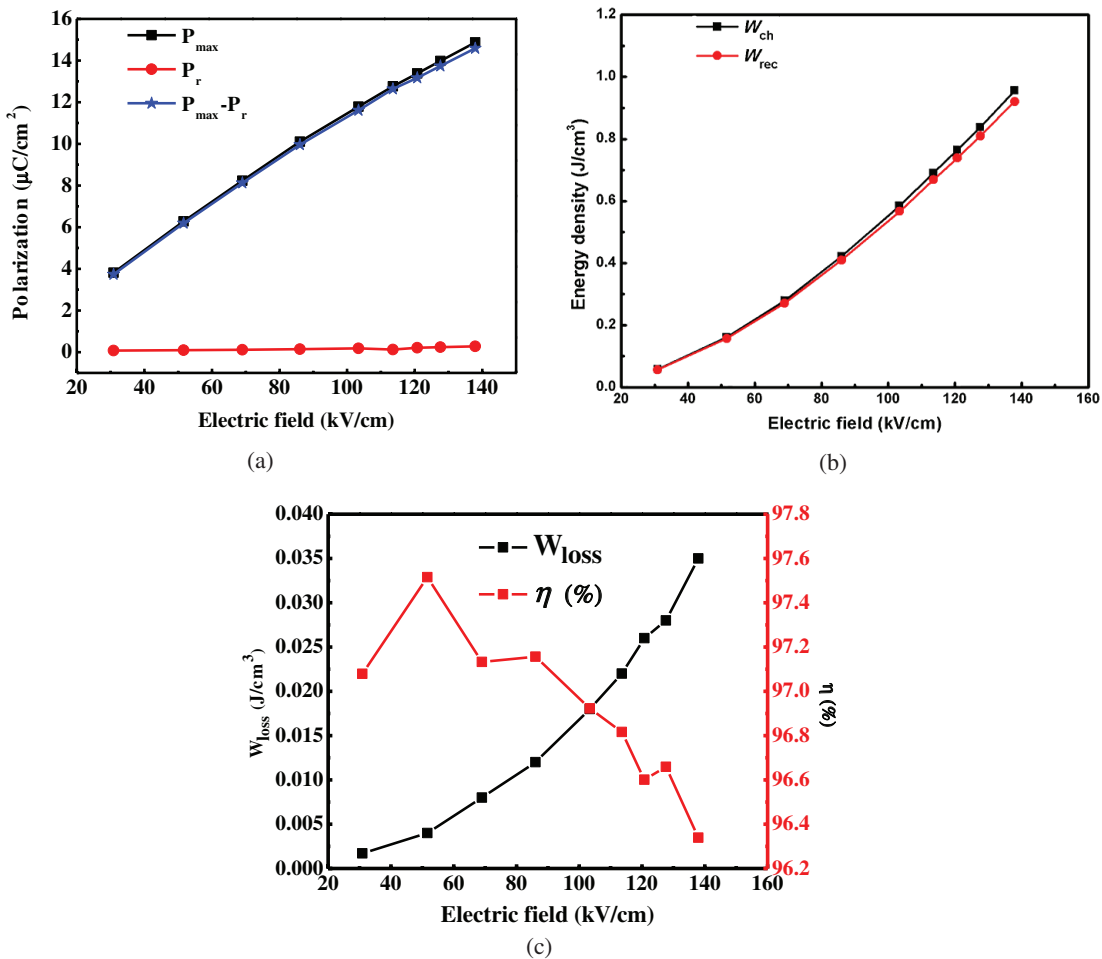


Fig. 6. (a) Variation in  $P_{\text{max}}$ ,  $P_r$ , and  $P_{\text{max}} - P_r$ , (b) energy storage density ( $W$ ), recoverable energy storage density ( $W_{\text{rec}}$ ), and (c) energy density loss ( $W_{\text{loss}}$ ), and energy storage efficiency ( $\eta$ ) of the NBN ceramics at room temperatures as a function of applied electric field.

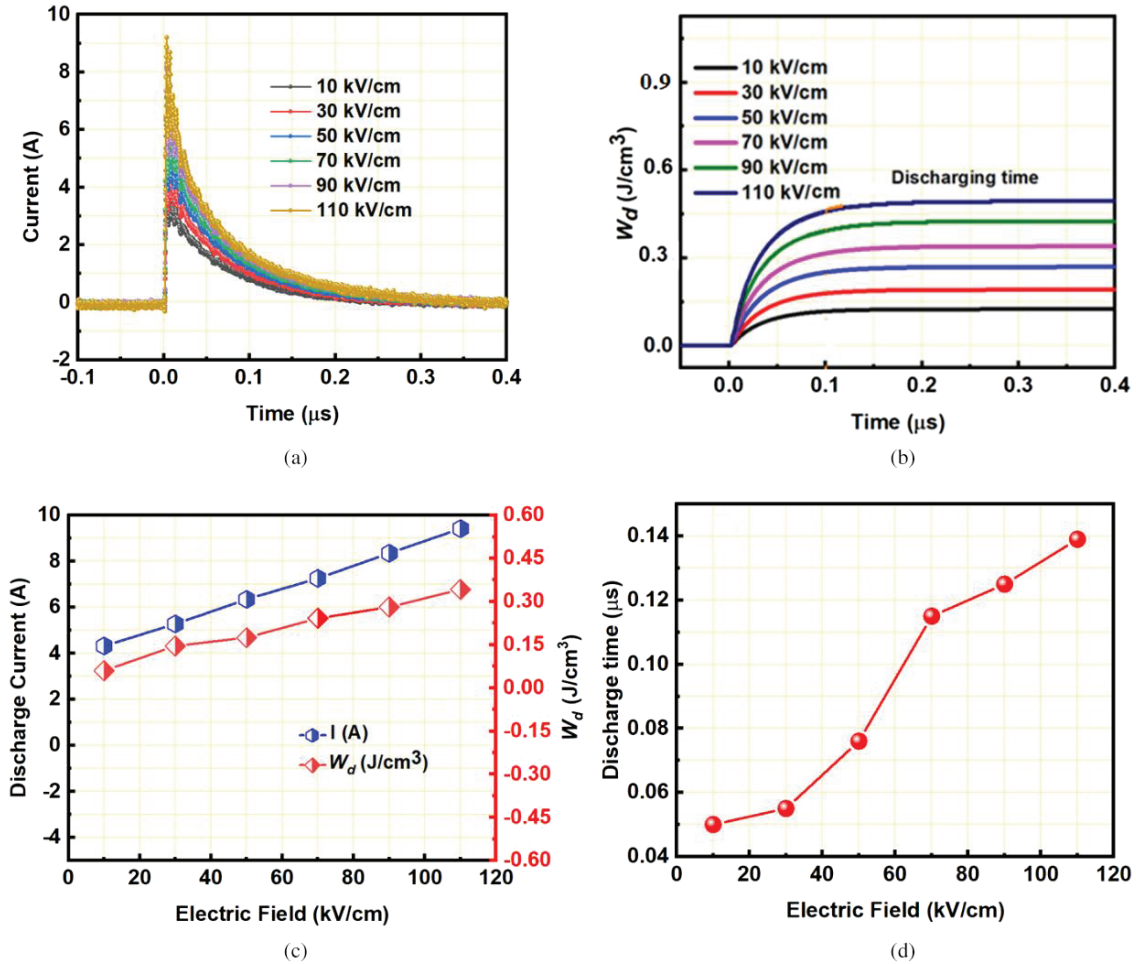


Fig. 7. (a) Variation of the discharge current of NBN ceramics with varying electric fields, (b) relation of  $W_d$  and discharge time of NBN ceramics, (c) variation of discharge current and  $W_d$  of NBN ceramics, and (d) variation of discharge time of NBN ceramics with the applied electric field.

slight deviation from linearity as the electric field increases to 138 kV/cm. The P–E loops of NBN ceramics possess minimal  $P_r$ , slim characteristic and high breakdown strength. Pure NN ceramics has a considerable  $P_r$  value showing its ferroelectric nature. Thus, Bi for Na’s partial substitution led to lower  $P_r$  with a slight decrease in  $P_{\text{max}}$  and an increase in  $P_{\text{max}}-P_r$  value. The significantly slimmer P–E loops reflect the typical relaxor characteristic of NBN ceramics. NBN ceramic’s relaxor characteristic confirms the existence of polar nano regions (PNRs) where the alignment and back-switching response with electric field is faster for micro-domains than the macroscopic domains. Such a behavior leads to slimmer P–E loops with negligible energy loss.<sup>60</sup> The large  $P_{\text{max}}$  and nearly zero  $P_r$  guarantee large value of  $W_{\text{rec}}$  and high  $\eta$ , enabling the NBN ceramics to be a good candidate material for energy storage capacitor applications.

Figure 6(a) display the variation in  $P_{\text{max}}$ ,  $P_r$ , and  $P_{\text{max}}-P_r$ ’s variation of NBN ceramics with applied electric field. The value of  $P_{\text{max}}$  increased from  $3.81 \mu\text{C}/\text{cm}^2$  to  $14.87 \mu\text{C}/\text{cm}^2$

upon increasing the electric field from 31 kV/cm to the breakdown value of 138 kV/cm. On the other hand, a very small increase in  $P_r$  was observed with an increase in the electric field. The  $P_r$ ’s value increased from  $0.096 \mu\text{C}/\text{cm}^2$  to  $0.283 \mu\text{C}/\text{cm}^2$  as the electric field was increased from 31 kV/cm to the breakdown value of 138 kV/cm. Similarly,  $\Delta P$  ( $P_{\text{max}}-P_r$ ) increased from  $3.71 \mu\text{C}/\text{cm}^2$  to  $14.58 \mu\text{C}/\text{cm}^2$  with an increase in the electric field from 31 kV/cm to 138 kV/cm. The significantly low value of  $P_r$  and considerable values of  $P_{\text{max}}$  and  $\Delta P$  help in attaining high  $W_{\text{rec}}$  and  $\eta$ . The values of storage energy density ( $W_s$ ), the recoverable energy density ( $W_{\text{rec}}$ ) are calculated from the P–E loops of NBN ceramics via Eqs. (1) and (2), while the energy density loss ( $W_{\text{loss}}$ ), and the efficiency ( $\eta$ ) are determined via Eq. (3) and are drawn as a function of electric field, as shown in Fig. 5. Upon increasing the electric field from 31 kV/cm to 138 kV/cm, both the energy densities, i.e.,  $W_s$  and  $W_{\text{rec}}$  of NBN ceramics increased. Furthermore, with increasing electric field, a slight increase in the energy density loss

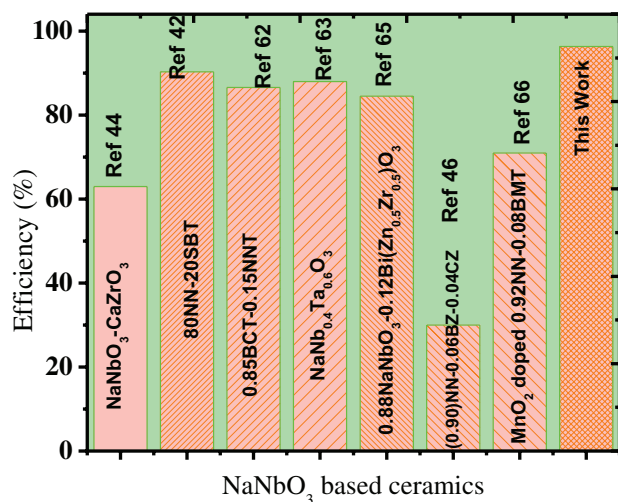


Fig. 8. Comparison of NBN ceramics' energy storage efficiency with some typical lead-free NN ceramics.

accompanied by a slight decrease in the energy storage efficiency was observed. In this study, a high storage energy density ( $W_s$ ) of  $0.96 \text{ J/cm}^3$ , a high recoverable energy density ( $W_{\text{rec}}$ ) of  $0.92 \text{ J/cm}^3$  accompanied by a high efficiency ( $\eta$ ) of 96.33% at DBS of 138 kV/cm were obtained for NBN ceramics.

The discharging rate of the energy by a ceramic capacitor is determined via a discharging experiment, where the measured discharge time ( $\tau_{0.9}$ ) is an important parameter. During this time interval ( $\tau_{0.9}$ ), a capacitor releases ~90% of its discharge energy. The discharged energy density ( $W_d$ ) and discharged current measured through a load resistance of  $1000 \Omega$  at different electric fields of 10 kV/cm, 30 kV/cm, 50 kV/cm, 70 kV/cm, 90 kV/cm, and 110 kV/cm, respectively, for NBN ceramics are shown in Fig. 6.

The discharge energy density ( $W_d$ ) through this experiment was calculated via the following equation<sup>61</sup>:

$$W_d = R \int i(t)^2 dt / V, \quad (6)$$

where  $R$  is the total load resistance ( $1000 \Omega$ ),  $V$  is the sample's volume, and  $i(t)$  is the discharged current. The measured discharged energy density from the discharged curve via this experiment was less than the recoverable energy density calculated from the P-E loop, as it has been performed at a low electric field (Fig. 7). However, it may also be slightly lower even if measured at the breakdown voltage in this experiment. Two factors are responsible for the lower values of the discharged energy density measured in this experiment. The first factor is the energy losses through the RC circuit, and the second is the variation in the thickness of different samples in the two measurements. The  $W_d$  value increased from  $0.09 \text{ J/cm}^3$  to  $0.4 \text{ J/cm}^3$  as the electric field was increased from 10 kV/cm to 110 kV/cm. The value of  $\tau_{0.9}$  increased from  $0.05 \mu\text{s}$  to  $0.13 \mu\text{s}$  when the electric field was increased from 10 kV/cm to 110 kV/cm. This lower discharging time

may be related to the increased polarization hysteresis with increasing electric field.

The electric-field-dependent (10–110 kV/cm) over damped discharge current curves of NBN ceramic are shown in Fig. 7, in which the current gained the maximum value rapidly in the overall discharge duration ( $<0.14 \mu\text{s}$ ).

Figure 8 shows a comparison of energy storage efficiency of NBN ceramics to other NN-based ceramics studied previously. It is evident that the energy storage efficiency obtained for NBN ceramic is higher in this study as compared to the previously reported NN-based ceramics.<sup>62–66</sup>

#### 4. Conclusion

Novel single perovskite phase NBN ceramics were successfully synthesized via a solid-state sintering process. In this study, good energy storage density ( $\sim 0.96 \text{ J/cm}^3$ ), good recoverable energy density ( $\sim 0.92 \text{ J/cm}^3$ ), and high energy storage efficiency ( $\sim 96.33\%$ ) at 138 kV/cm were obtained for the samples investigated in this study. A short pulse discharge time ( $\sim 0.14 \mu\text{s}$ ) was achieved at room temperature. The results illustrate that NBN ceramics can be a suitable candidate material for energy storage applications as pulsed power ceramic capacitors.

#### Acknowledgments

The Higher Education Commission (HEC) Islamabad Pakistan has financially supported this work through the National Research Program for University (7488/KPK/NRPU/R&D/HEC/ 2017). The financial support extended by HEC, Pakistan through Project No. NRPU-8148

#### Conflicts of Interest

There are no conflicts to declare this research work.

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