

# Growth behavior and electronic properties of $\text{Ge}_{n+1}$ and $\text{AsGe}_n$ ( $n = 1-20$ ) clusters: a DFT study

M. Benaïda<sup>1,†</sup>, K. E. Aiadi<sup>1</sup>, S. Mahtout<sup>2</sup>, S. Djaadi<sup>1</sup>, W. Rammal<sup>3</sup>, and M. Harb<sup>4,†</sup>

<sup>1</sup>Laboratoire de Développement des Energies Nouvelles et Renouvelables en Zones Arides, Université de Ouargla, 30000 Ouargla, Algeria

<sup>2</sup>Laboratoire de Physique Théorique, Faculté des Sciences Exactes, Université de Bejaïa, 06000 Begaïa, Algeria

<sup>3</sup>Faculty of Sciences, Lebanese University, Lebanon

<sup>4</sup>KAUST Catalysis Center (KCC), Physical Sciences and Engineering Division (PSE), King Abdullah University of Science and Technology (KAUST), Thuwal 23955-6900, Saudi Arabia

**Abstract:** We present a systematic computational study based on the density functional theory (DFT) aiming to highlight the possible effects of one As doping atom on the structural, energetic, and electronic properties of different isomers of  $\text{Ge}_{n+1}$  clusters with  $n = 1-20$  atoms. By considering a large number of structures for each cluster size, the lowest-energy isomers are determined. The lowest-energy isomers reveal three-dimensional structures starting from  $n = 5$ . Their relative stability versus atomic size is examined based on the calculated binding energy, fragmentation energy, and second-order difference of energy. Doping  $\text{Ge}_{n+1}$  clusters with one As atom does not improve their stability. The electronic properties as a function of the atomic size are also discussed from the calculated HOMO–LUMO energy gap, vertical ionization potential, vertical electron affinity, and chemical hardness. The obtained results are significantly affected by the inclusion of one As atom into a  $\text{Ge}_n$  cluster.

**Key words:** density functional theory; As–Ge clusters; structural properties; electronic properties

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## 1. Introduction

Studying clusters of various chemical elements has become a modern research topic in both physical and chemical communities over the last four decades<sup>[1, 2]</sup>, due to the size-dependent evolution of their fundamental properties and their technological applications in large variety of research fields, from catalysis to optoelectronics. Their particular structural, energetic, and electronic properties are fully understood and still constitute the subject of many research projects<sup>[3, 4]</sup>. Nano-scale materials (called clusters) with various sizes can provide different behaviors to that of the bulk material. The physical and chemical features of bimetallic clusters are dependent not only on the size and shape but also on the chemical composition and the atomic arrangement of the two metal elements<sup>[2]</sup>. Therefore, studying the changes in the structural and electronic properties of the cluster with its size has become important<sup>[4, 5]</sup>. Several theoretical calculations have been performed on pure and mixed neutral and charged clusters of group 14 elements<sup>[6, 7]</sup> especially silicon and germanium.

Numerous theoretical and experimental studies on Ge clusters have been published over the last decade<sup>[8]</sup> and different types of structures have been proposed<sup>[9, 10]</sup>. Theoretically, several computations have concluded that  $\text{Ge}_n$  cages can be stabilized through the encapsulation of guest atom inside the cage. This was seen before in  $\text{Si}_n$  cages. Matthias Brack *et al.*<sup>[11]</sup>

presented the global minimum of  $\text{CuGe}_{10}^+$  clusters as a magic number along with  $D_{4d}$  symmetry. Han *et al.*<sup>[12]</sup> have presented a theoretical investigation of very small  $\text{Ge}_n$  ( $n = 1-4$ ) clusters doped by Sn, and they found a charge transfer from Sn to Ge atoms. Singh *et al.*<sup>[13]</sup> have reported that the  $n$  capsulation can be utilized for stabilizing highly symmetric  $\text{Ge}_n$  cages having from 16 to 20 atoms. Wang and Han<sup>[14]</sup> have investigated  $\text{CuGe}_n$  ( $n = 2-13$ ) clusters and shown that Cu doping can decrease its binding energies, and so, the stability of  $\text{Ge}_{n+1}$  clusters. Zhao and Wang have studied in 2009 Mn-doped  $\text{Ge}_n$  clusters<sup>[15]</sup> and shown that Mn dopant can contribute to the stability increase of  $\text{Ge}_{n+1}$  clusters. Jaiswal and Kumar using studied the atomic and electronic structures of both neutral and negatively charged  $\text{ZrGe}_n$  ( $n = 1-21$ ) clusters using ab-initio calculations<sup>[16]</sup> and predicted cage-like stable geometries for  $n \geq 13$ . Siouani *et al.*<sup>[10]</sup> have investigated systematically the equilibrium geometries and electronic properties of  $\text{VGe}_n$  ( $n = 1-19$ ) clusters and found that V atom in  $\text{VGe}_n$  can make the stability stronger starting from  $n = 7$ . More recently, Mahtout *et al.*<sup>[17]</sup> have studied the structural, energetic, and electronic properties of  $\text{MGe}_n$  clusters with  $M = \text{Cu, Ag, Au}$  and  $n = 1-19$  using DFT approach. They have found the endohedral structures where the metal atom was incorporated inside the  $\text{Ge}_{n+1}$  cage appear at  $n = 10$  when the dopant is Cu and at  $n = 12$  for Ag or Au. Djaadi *et al.*<sup>[18]</sup> have investigated the structures and relative stability of pure  $\text{Ge}_{n+1}$ , neutral cationic and anionic  $\text{SnGe}_n$  ( $n = 1-17$ ) clusters. They found that the Sn atom occupied a peripheral position for  $\text{SnGe}_n$  clusters when  $n < 12$  and occupied a core position for  $n > 12$ .

Here, we report a systematic computational study based on the density functional theory (DFT) aiming to highlight the possible effects of one arsenic atom on the structural, energetic

Correspondence to: M Benaïda, [meriembenaïda@gmail.com](mailto:meriembenaïda@gmail.com); M Harb, [moussab.harb@kaust.edu.sa](mailto:moussab.harb@kaust.edu.sa)

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Table 1. Averaged bond length  $a$ , binding energy  $E_b$ , vertical ionization potential VIP, and vertical electron affinity VEA for  $\text{Ge}_2$ ,  $\text{Ge}_3$  and  $\text{As}_2$ .

Symmetry	Our work				Bibliographic data <sup>[24-40]</sup>			
	$a$ (Å)	$E_b$ (eV)	VIP (eV)	VEA (eV)	$a$ (Å)	$E_b$ (eV)	VIP (eV)	VEA (eV)
$\text{Ge}_2$	2.503	1.445	7.362	1.473	2.450	1.446	7.844	1.900
					2.413	~1.430	7.627	1.751
					2.540	1.620	7.934	1.549
					2.610	1.812		
					2.570	~1.350		
					2.420	1.410		
					2.440	1.320		
					2.421	1.230		
$\text{Ge}_3$	2.370	2.110	8.024	1.306	2.546	2.059	7.804	2.200
					2.400	2.240		
					2.476	2.150		
						2.040		
						1.860		
$\text{As}_2$	2.143	1.686	9.677	0.132	2.189	1.763		
					2.103			
					2.192			

ic, and electronic properties of different isomers of  $\text{Ge}_{n+1}$  in the atomic size range  $n = 1-20$  atoms. We believe this work is useful for deeply understanding the effects of incorporating one As atom into  $\text{Ge}_{n+1}$  clusters and can be considered as a guideline for future experiments. To the best of our knowledge, no systematic study has been addressed on neutral and charged  $\text{AsGe}_n$  clusters.

## 2. Computational methods

The electronic structure calculations of  $\text{AsGe}_n^q$  ( $n = 1-20$ ,  $q = 0, \pm 1$ ) clusters were performed using the density functional theory (DFT)<sup>[19]</sup> as implemented in the SIESTA program<sup>[20]</sup>. This code uses norm-conserving Troullier-Martins nonlocal pseudopotentials<sup>[2, 21]</sup> and employs flexible basis sets of localized Gaussian-type atomic orbitals<sup>[2]</sup>. The exchange correlation energy was evaluated using the generalized gradient approximation (GGA) parameterized by Perdew and Zunger<sup>[22]</sup> and by Perdew, Burke, and Ernserh of (PBE)<sup>[23]</sup>. The self-consistent field (SCF) calculations were carried out with convergence criterion of  $1 \times 10^{-4}$  a.u. for total energy. We used a double  $\zeta$  (DZ) basis with polarization function for As and Ge atoms. With energy shift parameter of 50 meV, the charge density was calculated in regular real-space grid with cut-off energy of 150 Ry. The simulated clusters were placed in a big cubic supercell with a parameter of 40 Å, including enough vacuums between neighboring clusters and periodic boundary conditions were imposed. To sample the Brillouin zone, only a single  $k$ -point centered at  $\Gamma$  was used because of the extended size of the supercell. The conjugated gradient method within Hellmann-Feynman forces was used and all the forces after structural relaxation were less than  $10^{-3}$  eV/Å.

We first searched for the lowest-energy structures of pure  $\text{Ge}_{n+1}$  clusters in the 1–20 atoms range by exploring various possibilities of isomers. Secondly, the most stable ground state structures obtained for  $\text{Ge}_{n+1}$  clusters were doped through substitution with one As atom. Then, the obtained  $\text{AsGe}_n$  clusters were optimized until reaching their ground states. In order to get lowest-energy structures of the  $\text{AsGe}_n$  clusters,

several initial isomeric structures, including some high and low symmetries, were optimized by placing one As atom in substitution in different possible sites of the pure corresponding  $\text{Ge}_{n+1}$  in order to get as close as possible to the low energy structures. Then, we cannot be sure that a more stable structure than those found in our calculations does not exist. We aim of our study is to highlight the variation of the properties of germanium cage clusters due to the As doping atom. We hope that this work would be useful to understand the influence of the As atom on the properties of germanium clusters and provide some guidelines for the probable future experimental studies. To check the validity of our computational method, benchmark tests have been done on  $\text{Ge}_2$ ,  $\text{Ge}_3$ , and  $\text{As}_2$  parameters. The values are reported in Table 1 together with available theoretical and experimental results. Our calculated results were found to be in line with the literature, confirming the reliability of our protocol to simulate small Ge clusters.

## 3. Results and discussion

### 3.1. Structural analysis

We report in Fig. 1 the lowest-energy structures obtained for  $\text{Ge}_{n+1}$  ( $n = 1-20$ ) and their corresponding isomers. Their energetic ordering is reported in Table 2. Our calculations reveal that almost all atoms are on the surface. Until  $n = 20$ , prolate-type geometries are in competition with the nearly spherical ones. The calculated results for the most favorable isomers are given in bold character. The most stable structures for  $n + 1 = 2, 3$ , and 4 adopt a planar disposition in line with previous works<sup>[10, 14, 18, 24, 25]</sup> using DFT different calculations. The triangular geometry with  $C_{2v}$  symmetry is found to be the lowest-energy structure for  $\text{Ge}_3$ . The lowest-energy isomer of the tetramer  $\text{Ge}_4$  has  $D_{2h}$  symmetry in line with previous findings<sup>[15, 18, 20, 26, 27]</sup>. The most favorable isomer for  $\text{Ge}_5$  cluster reveals a triangular bipyramid disposition with  $D_{3h}$  symmetry, which is also in line with the previous data<sup>[10, 14, 24, 26]</sup>. The lowest-energy  $\text{Ge}_6$  cluster has bicapped quadrilateral structure with  $C_{2v}$  symmetry, in good agreement with the previous data<sup>[10, 18, 26]</sup>. For  $\text{Ge}_7$  cluster, the most stable structure reveals

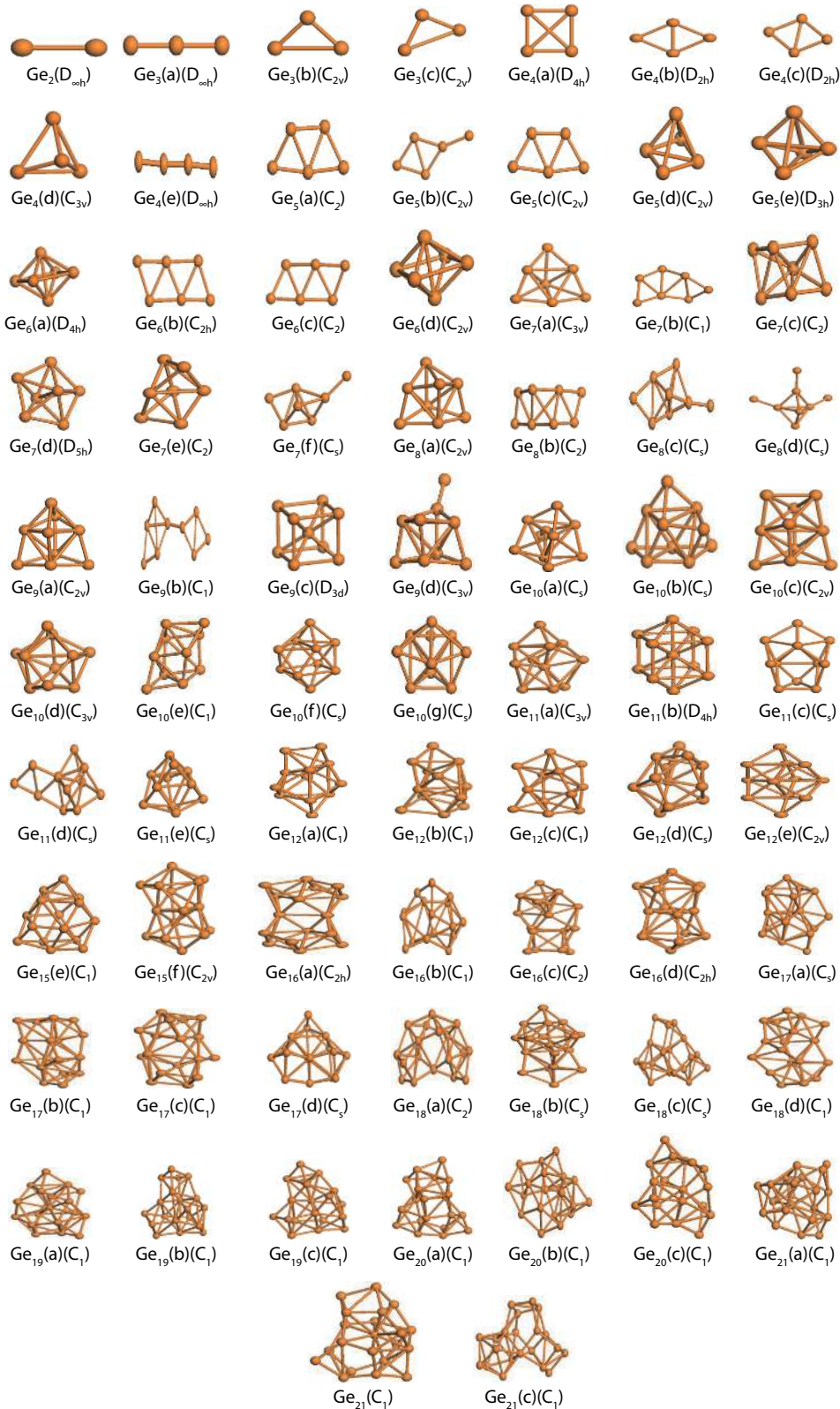


Fig. 1. (Color online) Most favorable structures together with their corresponding isomers for  $\text{Ge}_{n+1}$  ( $n = 1-20$ ) clusters.

a pentagonal bipyramid of  $D_{5h}$  symmetry. Other researchers also previously obtained similar results for  $\text{Ge}_7$ <sup>[10, 14, 24, 26]</sup>. For  $\text{Ge}_8$  clusters, the most stable isomer shows a capped pentagonal bipyramid disposition of  $C_{2v}$  symmetry, as obtained in earlier works<sup>[20, 28]</sup>.  $\text{Ge}_9$  is a capped pentagonal bipyramid structure of  $C_{2v}$  symmetry. The most favorable isomer of  $\text{Ge}_{10}$  cluster is a capped pentagonal basis structure and has  $C_{3v}$

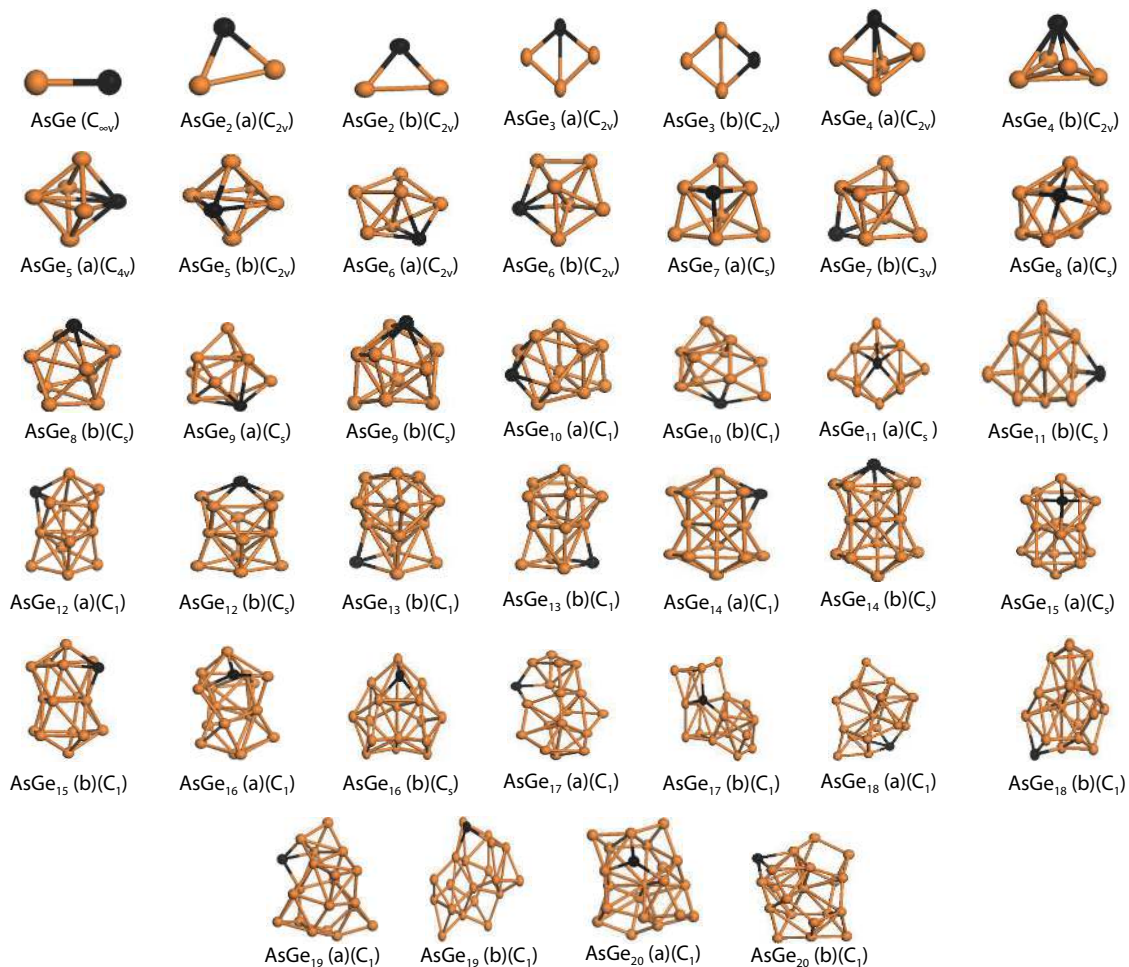
symmetry. The lowest-energy  $\text{Ge}_{11}$  cluster shows a compact near spherical geometry of  $C_3$  symmetry. As per  $\text{Ge}_{12}$  and  $\text{Ge}_{13}$  clusters, prolate-type structure with  $C_{2v}$  symmetry was always preferred. The shape of  $\text{Ge}_{14}$  is a prolate structure with  $C_3$  symmetry. For  $\text{Ge}_{15}$  and  $\text{Ge}_{16}$  clusters, prolate-type geometry was obtained with  $C_{2v}$  and  $C_{2h}$  symmetries, respectively. The lowest-energy isomer for  $\text{Ge}_{17}$  possesses a near spherical geo-

Table 2. Group of symmetry,  $E_b$ ,  $\Delta E$ , VIP, VEA,  $\eta$ , and  $a_{\text{Ge-Ge}}$  for pure  $\text{Ge}_{n+1}$  ( $n = 1 - 20$ ) clusters.

Cluster size ( $n$ )	Symmetry	$E_b$ (eV/atom)	$\Delta E$ (eV)	VEA (eV)	VIP (eV)	$\eta$ (eV)	$a_{\text{Ge-Ge}}$ (Å)
1	$D_{\infty h}$	1.445	0.265	7.362	1.473	5.889	2.503
2	$D_{\infty h}$	2.048	1.272	7.557	1.476	6.081	2.339
	$C_{2v}$	2.109	1.543	8.024	1.305	6.719	2.370
	$C_{2v}$	2.110	1.543	8.024	1.306	6.718	2.370
3	$D_{4h}$	2.228	0.455	6.834	1.599	5.235	2.550
	$D_{2h}$	2.556	1.180	7.734	1.758	5.976	2.576
	$D_{2h}$	2.557	1.179	7.733	1.758	5.976	2.597
	$C_{3v}$	2.078	0.660	7.056	1.717	5.339	2.479
	$D_{\infty h}$	2.061	1.169	7.114	2.102	5.012	2.346
4	$C_2$	2.518	1.089	6.911	2.149	4.762	2.606
	$C_{2v}$	2.450	1.026	6.766	2.207	4.559	2.564
	$C_{2v}$	2.499	0.808	6.774	2.070	4.704	2.618
	$C_{2v}$	2.504	0.577	7.588	2.560	5.028	2.775
	$D_{3h}$	2.707	2.036	8.672	0.218	8.454	2.547
5	$D_{4h}$	2.847	1.998	7.777	1.372	6.405	2.782
	$C_{2h}$	2.668	1.236	7.142	1.820	5.322	2.760
	$C_2$	2.672	1.386	7.172	1.790	5.382	2.606
	$C_{2v}$	2.848	1.957	7.742	1.359	6.383	2.710
6	$D_{3v}$	2.877	2.350	7.460	1.064	6.396	2.734
	$C_1$	2.687	0.164	7.304	1.673	5.631	2.647
	$C_2$	2.843	1.170	7.034	1.749	5.285	2.754
	$D_{5h}$	2.974	1.836	7.875	1.774	6.101	2.747
	$C_2$	2.843	1.170	7.034	1.748	5.286	2.754
	$C_s$	2.617	0.755	6.620	2.102	4.518	2.701
7	$C_{2v}$	2.866	0.980	7.216	2.232	4.984	2.776
	$C_2$	2.735	0.977	6.937	2.194	4.743	2.658
	$C_s$	2.739	0.708	6.556	2.212	4.344	2.730
	$C_s$	2.422	0.431	6.194	2.805	3.389	2.651
8	$C_{2v}$	2.985	1.654	7.126	1.570	5.556	2.782
	$C_1$	2.700	1.066	6.849	2.394	4.455	2.570
	$D_{3d}$	2.574	0.227	6.288	2.271	4.017	2.798
	$C_{3v}$	2.827	1.102	7.068	2.313	4.755	2.686
9	$C_s$	2.848	1.459	6.462	2.146	4.316	2.742
	$C_s$	3.002	1.753	7.137	1.662	5.475	2.779
	$C_{2v}$	2.968	1.379	7.070	1.981	5.089	2.785
	$C_{3v}$	3.082	1.812	7.432	1.857	5.575	2.775
	$C_1$	2.953	1.295	6.900	1.913	4.987	2.738
	$C_s$	3.013	1.015	7.152	2.393	4.759	2.794
	$C_s$	3.003	1.547	7.410	1.383	6.027	2.771
10	$C_{3v}$	2.792	1.063	6.402	2.364	4.038	2.714
	$D_{4h}$	2.715	0.793	6.326	1.975	4.351	2.816
	$C_s$	2.907	0.962	6.605	2.084	4.521	2.748
	$C_s$	2.892	0.990	6.820	2.318	4.502	2.724
	$C_s$	3.029	1.258	7.107	1.332	5.775	2.770
11	$C_s$	2.936	0.803	6.852	2.547	4.305	2.798
	$C_1$	2.955	0.916	6.759	2.345	4.414	2.784
	$C_1$	2.964	1.059	6.696	2.175	4.521	2.801
	$C_s$	2.979	0.370	6.656	2.780	3.876	2.793
	$C_{2v}$	3.032	1.793	7.402	1.258	6.144	2.744
12	$C_{2v}$	3.050	1.181	8.243	1.290	6.953	2.760
	$C_2$	3.007	1.153	7.153	2.390	4.553	2.792
	$C_1$	3.015	0.945	6.785	2.412	4.373	2.835
	$C_s$	2.990	1.130	6.608	2.102	4.506	2.769
	$C_1$	3.015	0.944	7.302	1.779	5.523	2.831
13	$C_{3v}$	2.922	0.987	5.967	2.633	3.334	2.696
	$C_s$	3.026	1.107	6.625	2.179	4.446	2.700
	$O_h$	2.977	1.007	6.915	3.095	3.820	2.659
	$C_1$	2.986	1.036	6.471	2.127	4.344	2.784
	$C_s$	3.092	1.628	7.486	1.539	5.947	2.797
14	$D_{3d}$	2.864	0.495	6.726	2.977	3.749	2.666
	$C_1$	2.999	1.234	6.800	2.311	4.489	2.829
	$C_s$	2.959	0.941	6.553	2.393	4.160	2.785
	$C_1$	3.022	1.149	6.579	2.178	4.401	2.812
	$C_1$	3.023	1.156	6.927	2.468	4.459	2.786
	$C_{2v}$	3.082	0.899	7.339	1.851	5.488	2.814

Continued from Table 2

Cluster size ( $n$ )	Symmetry	$E_b$ (eV/atom)	$\Delta E$ (eV)	VEA (eV)	VIP (eV)	$\eta$ (eV)	$a_{\text{Ge-Ge}}$ (Å)
15	$C_{2h}$	3.027	1.436	6.840	2.308	4.532	2.715
	$C_1$	3.020	0.883	6.555	2.494	4.061	2.755
	$C_2$	3.034	1.364	6.755	2.262	4.493	2.808
	$C_{2h}$	3.095	1.393	7.548	1.753	5.795	2.780
16	$C_s$	3.050	1.103	6.588	2.343	4.245	2.781
	$C_1$	3.043	1.254	6.696	2.428	4.268	2.772
	$C_1$	3.038	1.022	6.616	2.500	4.116	2.762
	$C_s$	3.077	0.858	7.075	1.716	5.359	2.823
17	$C_2$	3.014	0.864	6.493	2.590	3.903	2.869
	$C_s$	3.013	0.948	6.537	2.537	4.000	2.790
	$C_s$	3.009	0.555	6.400	2.815	3.585	2.782
	$C$	3.062	1.452	7.152	1.486	5.666	2.729
18	$C_1$	3.053	0.966	6.473	2.477	3.996	2.843
	$C_1$	3.019	0.779	6.516	2.728	3.788	2.771
	$C_1$	3.019	0.780	6.517	2.727	3.790	2.771
19	$C_1$	3.046	0.828	6.387	2.596	3.791	2.735
	$C_1$	3.033	0.962	6.365	2.455	3.910	2.768
	$C_1$	3.001	0.828	6.418	2.650	3.768	2.781
20	$C_1$	3.041	0.743	6.372	2.679	3.693	2.769
	$C_1$	3.061	3.061	6.403	2.845	3.558	2.735
	$C_1$	3.050	0.559	6.182	2.734	3.448	2.761

Fig. 2. (Color online) Most favorable structures and their corresponding isomers of AsGe<sub>n</sub> ( $n = 1-20$ ) clusters.

metry of  $C_s$  symmetry. From Ge<sub>18</sub> to Ge<sub>21</sub> clusters, prolate-type shape with  $C_1$  symmetry was always preferred.

The most favorable geometries of AsGe<sub>n</sub> ( $n = 1-20$ ) clusters and their corresponding isomers are summarized in

Fig. 2, whereas their energetic ordering is reported in Table 3. The AsGe<sub>n</sub> clusters adopt somehow similar structures to their corresponding Ge<sub>n+1</sub> except for  $n = 8, 10, 11$ , and 16. In all cases, the arsenic atom is always located on the surface. The

Table 3. Group of symmetry,  $E_b$ ,  $\Delta E$ , VEA, VIP,  $\eta$ , and  $a_{\text{Ge-Ge}}$ ,  $a_{\text{As-Ge}}$  for  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters.

Cluster size ( $n$ )	Symmetry	$E_b$ (eV/atom)	$\Delta E$ (eV)	VEA (eV)	VIP (eV)	$\eta$ (eV)	$a_{\text{Ge-Ge}}$ (Å)	$a_{\text{As-Ge}}$ (Å)
1	(a) $C_{\infty v1}$	1.426	0.171	2.175	8.142	5.967	-	2.350
2	(a) $C_{2v}$	2.139	1.110	0.969	7.425	6.456	2.775	2.445
	(b) $C_{2v}$	2.139	1.109	0.778	8.198	7.410	2.775	2.445
3	(a) $C_{2v}$	2.380	0.589	1.440	6.704	5.264	2.603	2.543
	(b) $C_{2v}$	2.418	1.266	1.726	7.377	5.651	2.661	2.473
4	(a) $C_{2v}$	2.575	0.208	0.451	8.640	8.189	2.738	2.661
	(b) $C_{2v}$	2.641	0.917	0.124	8.809	8.685	2.692	2.734
5	(a) $C_{4v}$	2.705	1.116	0.928	6.657	5.729	2.807	2.679
	(b) $C_{2v}$	2.702	1.058	1.088	6.951	5.863	2.782	2.706
6	(a) $C_{2v}$	2.837	1.279	1.632	6.614	4.982	2.786	2.703
	(b) $C_{2v}$	2.837	1.278	1.632	6.615	4.983	2.786	2.704
7	(a) $C_s$	2.814	0.503	2.060	7.179	5.119	2.788	2.635
	(b) $C_{3v}$	2.835	0.743	1.701	7.022	5.321	2.821	2.567
8	(a) $C_s$	2.908	0.571	1.279	6.661	5.382	2.808	2.678
	(b) $C_s$	2.909	0.571	1.279	6.661	5.382	2.808	2.678
9	(a) $C_s$	2.977	1.066	1.604	6.271	4.667	2.774	2.819
	(b) $C_s$	2.977	1.064	1.603	6.271	4.668	2.774	2.819
10	(a) $C_1$	2.949	0.759	1.051	6.663	5.612	2.793	2.697
	(b) $C_1$	2.949	0.898	0.884	6.953	6.069	2.774	2.716
11	(a) $C_s$	2.945	1.026	0.783	6.609	5.826	2.740	2.813
	(b) $C_s$	2.940	0.773	1.305	6.481	5.176	2.775	2.596
12	(a) $C_1$	2.993	0.393	1.019	8.052	7.033	2.788	2.693
	(b) $C_s$	3.003	0.396	1.044	7.736	6.692	2.814	2.667
13	(a) $C_1$	3.016	0.696	1.518	6.578	5.060	2.826	2.606
	(b) $C_1$	3.016	0.697	1.520	6.574	5.054	2.780	2.607
14	(a) $C_1$	3.042	0.581	1.248	7.125	5.877	2.821	2.681
	(b) $C_s$	3.046	0.797	1.174	6.946	5.772	2.828	2.783
15	(a) $C_s$	3.045	0.782	1.482	6.984	5.502	2.801	2.736
	(b) $C_1$	3.043	0.846	1.300	6.917	5.617	2.817	2.719
16	(a) $C_1$	3.045	0.538	1.394	6.977	5.583	2.813	2.595
	(b) $C_s$	3.035	0.810	1.547	6.823	5.276	2.830	2.593
17	(a) $C_1$	3.031	0.623	0.763	7.031	6.268	2.766	2.566
	(b) $C_1$	3.030	0.390	0.453	7.310	6.857	2.747	2.556
18	(a) $C_1$	3.017	0.627	1.867	6.346	4.479	2.777	2.629
	(b) $C_1$	3.024	0.594	2.000	6.361	4.361	2.803	2.592
19	(a) $C_1$	3.023	0.559	1.916	6.098	4.182	2.769	2.581
	(b) $C_1$	3.029	0.519	1.667	6.834	5.167	2.732	2.580
20	(a) $C_1$	3.065	0.667	1.961	6.558	4.597	2.739	2.573
	(b) $C_1$	3.059	0.579	1.870	6.937	5.067	2.745	2.601

$\text{AsGe}_2$  cluster shows a triangular geometry of  $C_{2v}$  symmetry with two equivalent As–Ge bonds of 2.445 Å and one Ge–Ge bond of 2.775 Å. The As–Ge bond distance of 0.09 Å is larger than that in  $\text{AsGe}$  dimer. The most stable structure of  $\text{AsGe}_3$  cluster presents a planar  $C_{2v}$  symmetry with a binding energy of 2.418 eV/atom, which is smaller than that of tetramer  $\text{Ge}_4$  (2.557 eV/atom). For  $\text{AsGe}_4$ , a distorted rectangular pyramid with  $C_{2v}$  symmetry is found with a binding energy of 0.066 eV/atom, which is also smaller than  $\text{Ge}_5$ . The Ge–Ge and As–Ge bond lengths are 2.692 and 2.734 Å, respectively. For  $\text{AsGe}_5$ , the As atom is located at the convex site of a quasi-rectangular bipyramid structure of  $C_{4v}$  symmetry, As–Ge bond distance of 2.679 Å, and an average Ge–Ge bond distance of 2.807 Å. The lowest energy isomer for  $\text{AsGe}_6$  cluster is a structure with  $C_{2v}$  point group symmetry, As–Ge bond length of 2.704 Å, and an average Ge–Ge bond distance of 2.786 Å. For  $\text{AsGe}_7$  cluster, the lowest-energy isomer reveals a low-lying structure with a planar  $C_{3v}$  symmetry and a binding energy of

2.835 eV/atom, which is smaller than that for tetramer  $\text{Ge}_8$  (2.866 eV/atom). For  $\text{AsGe}_8$  cluster, its binding energy of only 0.076 eV/atom is also smaller than that obtained for  $\text{Ge}_9$  cluster with  $C_s$  symmetry of the ground state isomer. The lowest-energy structure of  $\text{AsGe}_9$  cluster has  $C_s$  symmetry combining two irregular hexagonal prisms with As atom on top of one of them. The ground state geometry of  $\text{AsGe}_{10}$  has  $C_1$  point group symmetry. The As atom tends to be stabilized on the surface. For  $n = 11, 12, 13, 14, 15, 16,$  and  $17$ , prolate structures were found to be the most stable in their ground state. Its binding energies are much smaller than  $\text{Ge}_{n+1}$ .  $\text{AsGe}_{18}$  has a lowest-energy structure with  $C_1$  point group symmetry. The As atom tends to be stabilized on the surface. The most favorable isomer for  $\text{AsGe}_{19}$  cluster shows prolate-like and cage-like structures with  $C_1$  symmetry and a calculated binding energy of 3.029 eV/atom, which is close to that of tetramer  $\text{Ge}_{20}$  (3.046 eV/atom). For  $n = 20$ , the lowest-energy isomer combines a prolate-like structure with the cage-like one. The binding energy

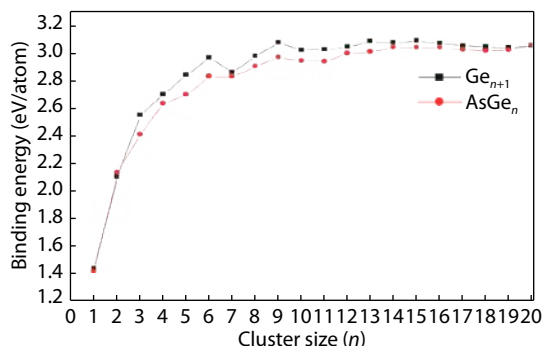


Fig. 3. (Color online) Evolution of the binding energy per atom for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

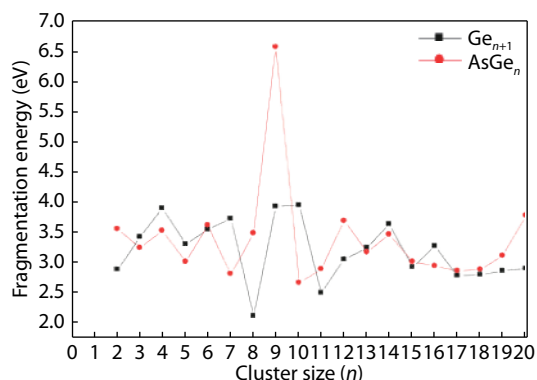


Fig. 4. (Color online) Evolution of the fragmentation energy of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

of  $\text{AsGe}_{20}$  (0.004 eV/atom) is almost the same than that obtained for the ground state structure of the pure  $\text{Ge}_{21}$  cluster.

## 3.2. Relative stability

### 3.2.1. Binding energy

The size dependence on the binding energies per atom for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters are shown in Fig. 3. As expected, the bonding energy gradually increases with increasing size, and this can be associated with the increasing average number of neighbors per atom. For  $\text{AsGe}_n$ , we observe that the binding energies are lower than those for  $\text{Ge}_{n+1}$ . This means that doping with As atoms has no immediate effects on enhancing the stability of germanium cluster at small size. In most of  $\text{AsGe}_n$  clusters, the final structures do not differ from that of the corresponding pure germanium cluster. This may be due to the equivalence in the nature of bonding, the size and the atomic mass between the two metalloids arsenic and germanium used in this study. However, for  $n = 2$  and  $n = 20$  we observe that the binding energy per atom of  $\text{AsGe}_n$  clusters is larger than those of corresponding pure  $\text{Ge}_{n+1}$  clusters. Then, the substituting a Ge atom by a As atom increases the stability these two clusters. An increase in the binding energy is obtained with 1.426 eV for  $n = 2$  to 2.837 eV for  $n = 6$ , and then non-monotonic and slow growth could be reached until  $n = 20$ .

### 3.2.2. Fragmentation energy

Fig. 4 shows the plot of the size dependence on the fragmentation energies of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters. An oscillating behavior is observed. The clusters with large values of fragmentation energy are relatively stronger in thermodynamic stability than neighboring clusters. Consequently, the

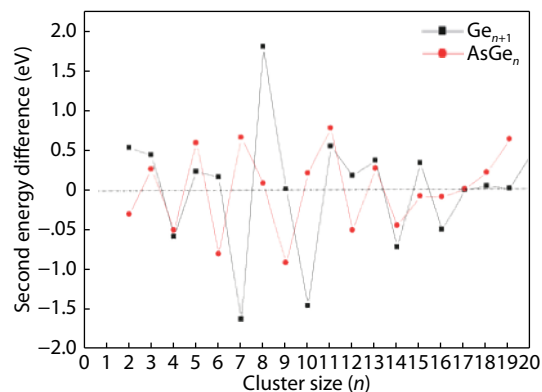


Fig. 5. (Color online) Evolution of the second-order difference of energy for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

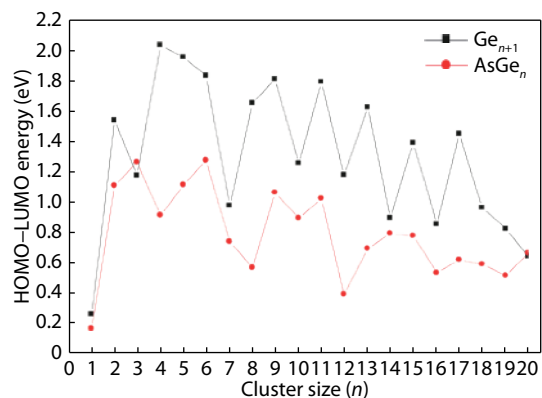


Fig. 6. (Color online) Evolution of the HOMO-LUMO gap for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

thermodynamic stabilities of  $\text{Ge}_5$ ,  $\text{Ge}_8$ ,  $\text{Ge}_{10}$ ,  $\text{Ge}_{11}$ ,  $\text{AsGe}_6$ ,  $\text{AsGe}_9$ ,  $\text{AsGe}_{12}$ , and  $\text{AsGe}_{20}$  clusters are relatively strong.

### 3.2.3. Second-order difference

The evolution of the second-order difference of energies for the most favorable structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of the cluster size is plotted in Fig. 5. The curve shows pronounced peaks for  $\text{AsGe}_n$  at range size  $n = 3, 5, 7, 10, 11, 13$ , and 19 atoms. This suggests these clusters to be more favorable than their neighbors. In cluster physics, if the values of  $\Delta_2 E$  are positive this means that the dissociation of As atom is an unfavorable process and the clusters are particularly stable. It can also be seen that the curve of  $\text{Ge}_{n+1}$  clusters with range size  $n = 2, 3, 5, 6, 8, 11, 12, 13, 15$ , and 20 exhibit higher stability than their neighbors. As a consequence, the stability of  $\text{AsGe}_n$  structures with  $n = 3, 5, 11, 13$  atoms correlates with the stability of the corresponding  $\text{Ge}_{n+1}$  structures, where the  $\text{AsGe}_n$  structure was maintained the same upon the incorporation of As dopant.

## 3.3. Electronic properties

### 3.3.1. HOMO-LUMO gap

In order to obtain insight into the kinetic stability of  $\text{AsGe}_n$  clusters, we calculated and analyzed the HOMO-LUMO gap. In general, the reactivity of the cluster decreases with increasing the HOMO-LUMO gap<sup>[28]</sup>. Fig. 6 reports the size dependence of HOMO-LUMO gap for the most favorable structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters. The decreasing behavior with

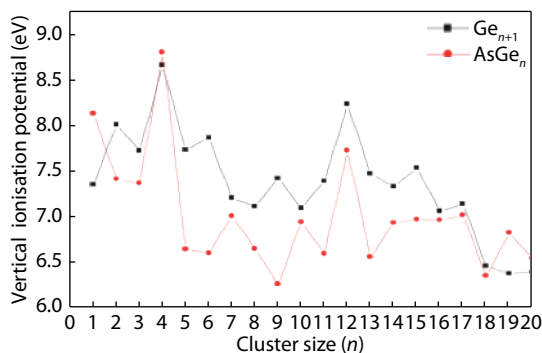


Fig. 7. (Color online) Evolution of the vertical ionization potential (VIP) for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

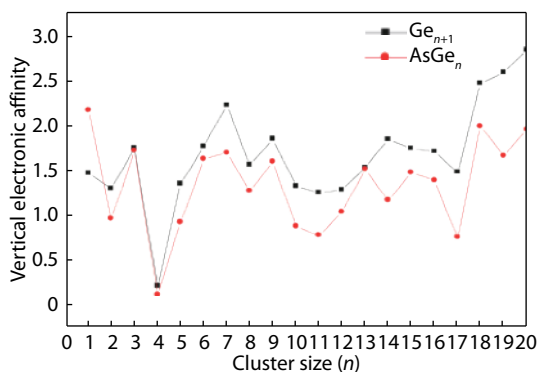


Fig. 8. (Color online) Evolution of the vertical electron affinity (VEA) for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

the size is important for  $\text{Ge}_{n+1}$  clusters, while is less pronounced for  $\text{AsGe}_n$  clusters. Overall, the gaps of  $\text{AsGe}_n$  are much lower than those obtained for  $\text{Ge}_{n+1}$  clusters, except for  $n = 3$  and  $20$ . The value for  $\text{AsGe}_n$  roughly oscillates in between  $0.171$  and  $1.278$  eV, which indicates a chemical activity increase of  $\text{Ge}_{n+1}$  clusters when doped with As. Doping  $\text{Ge}_{n+1}$  cages with an As atom leads to a significant HOMO–LUMO gap reduction in  $\text{AsGe}_n$  clusters. This means that the chemical activity of  $\text{AsGe}_n$  is higher than that of  $\text{Ge}_{n+1}$  clusters and the inserted As atom highlights the metallic character of  $\text{AsGe}_n$  clusters. It should also be noted that  $\text{Ge}_4$  cluster possesses the largest HOMO–LUMO gap of  $2.036$  eV, which indicates that  $\text{Ge}_4$  cluster is expected to have an enhanced chemical stability. As a consequence, the substitution of an As atom would affect the chemical features of pure  $\text{Ge}_{n+1}$  clusters.

### 3.3.2. Vertical ionization potential (VIP) and vertical electronic affinity (VEA)

The size dependence on the vertical ionization potential (VIP) for the most favorable geometries of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters are displayed in Fig. 7. For  $\text{AsGe}_n$  clusters, the VIP reveals an oscillating trend up to  $n = 14$ . All values are in the  $6.2-8.8$  eV range and decreases slowly as the cluster size increases and it is well known that when the VIP becomes smaller, the cluster will be more close to a metallic system. This means that the clusters of  $\text{AsGe}_n$  with size more than 6 atoms exhibit high metallic character which, consequently, these clusters can more easily lose one electron comparatively to the clusters of smaller size. The smallest VIP values are observed for  $\text{AsGe}_5$ ,  $\text{AsGe}_6$ ,  $\text{AsGe}_8$ ,  $\text{AsGe}_9$ ,  $\text{AsGe}_{11}$ ,  $\text{AsGe}_{13}$ ,  $\text{AsGe}_{18}$  and  $\text{AsGe}_{20}$  indicating that these clusters are more readily ion-

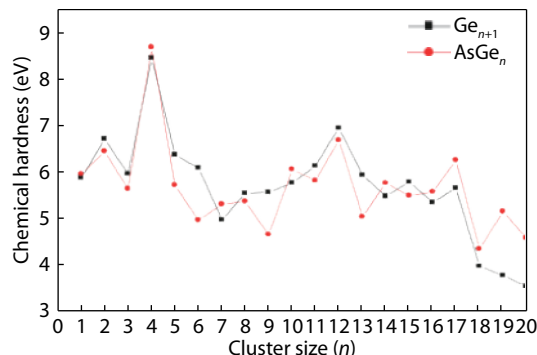


Fig. 9. (Color online) Evolution of the chemical hardness  $\eta$  for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size.

ized than the others. The cluster  $\text{AsGe}_4$  has large VIP value ( $8.809$ ). In Fig. 8, we plotted the cluster size-dependent VEA for  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  clusters. It can be seen that the electron affinity reveal also an oscillating trend with an increasing behavior with the size, which means the larger clusters are expected to capture more easily electrons more easily. This means that the small  $\text{AsGe}_n$  clusters will become gradually unstable after they acquired an electron. The calculated values of VEA for the most stable  $\text{AsGe}_n$  clusters are much lower than the VIP values which indicating that these clusters can easily accept one electron.

### 3.3.3. Chemical hardness

Fig. 9 shows the evolution of the chemical hardness for the lowest energy structures of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters as a function of cluster size. Our calculations reveal that  $\text{AsGe}_4$  clusters have the largest chemical hardness of  $8.454$  eV, confirming the better stability of this cluster as compared to the neighboring ones. Other local peaks are also observed for  $n = 12$  and  $17$ , leading to the conclusion that  $\text{AsGe}_{12}$  and  $\text{AsGe}_{17}$  will be less reactive than other cluster sizes. These clusters are very inert and can be considered as good candidates to the fabrication assembled cluster materials for application in nano-electronics and nanotechnologies. It has been established that chemical hardness is an electronic parameter that may characterize the relative stability of small clusters through the principle of maximum hardness (PMH) proposed by Pearson<sup>[39,41]</sup>. The clusters with high values of hardness are less reactive and more stable.

## 4. Conclusion

We have systematically investigated the structural, energetic and electronic properties of  $\text{Ge}_{n+1}$  and  $\text{AsGe}_n$  ( $n = 1-20$ ) clusters by means of DFT-based first principles quantum computations. The  $\text{AsGe}_n$  clusters adopted somehow similar structures as those obtained for  $\text{Ge}_{n+1}$  except for  $n = 8, 10, 11$ , and  $16$ , which significantly differed from their corresponding  $\text{Ge}_{n+1}$ . In all cases, the As-doping atom was found to always be located on the surface. Their relative stabilities have been examined through the calculated binding energies, fragmentation energies, and second-order difference of energies. Their electronic features such as HOMO–LUMO energy gaps, vertical ionization potentials, vertical electron affinities, and chemical hardness were also examined.

Our theoretical study could give detailed and relevant information to deeply understand the possible effects of doping one single As atom on the properties of  $\text{Ge}_{n+1}$  clusters. We



believe this work will provide guidelines for future experimental work.

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