# REVIEWS

# Growth and fundamentals of bulk $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals

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**Abstract:** The rapid development of bulk  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals has attracted much attention to their use as ultra-wide bandgap materials for next-generation power devices owing to its large bandgap (~ 4.9 eV) and large breakdown electric field of about 8 MV/cm. Low cost and high quality of large  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single-crystal substrates can be attained by melting growth techniques widely used in the industry. In this paper, we first present an overview of the properties of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals in bulk form. We then describe the various methods for producing bulk  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals and their applications. Finally, we will present a future perspective of the research in the area in the area of single crystal growth.

**Key words:**  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>; crystal structure; bulk crystal growth; applications

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# 1. Introduction

Beta gallium oxide has attracted rapidly growing interest in the past decade and is considered a promising candidate for power device applications. Gallium oxide is a significant semiconductor oxide and has a large bandgap energy (~4.9 eV) with a decent mobility (~100 cm<sup>2</sup>/Vs), a high breakdown strength (8 MV/cm), and good thermal stability. It was first studied in the 1950s, and its crystal structure was determined by Geller (1960)<sup>[1]</sup>. Ga<sub>2</sub>O<sub>3</sub> has considerable potential in high-electric-field-strength and high-temperature applications owing to its wide bandgap. There exists an empirical relationship between the electrical breakdown value and semiconductor bandgap. A slight increase in the bandgap value can greatly improve the electric field strength (discussed in section 3.2). Therefore, Ga<sub>2</sub>O<sub>3</sub> is an important candidate for high breakdown voltage devices. Bulk  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals have been successfully growth to large and high-quality crystals at a low production cost; this can make them competitive in the advancement technological applications. In this article, we briefly review the growth, doping, and application of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulk single crystals. This article is divided into three main sections. The first section presents crystal structures and the electronic band structure. The second section presents the growth technique. The final section describes the role of the defects and impurities.

# 2. Crystal structure of $\beta$ -Ga<sub>2</sub>O<sub>3</sub>

We discuss the crystal structure and electronic band structure of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and summarize the main points. Gallium oxide has a melting point of approximately 1800 °C, and five polymorphs named  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\delta$ , and  $\varepsilon$  exist in nature depending strictly on the preparation conditions<sup>[2]</sup>. Among these phases,

Correspondence to: C T Xia, xia\_ct@siom.ac.cn Received 30 SEPTEMBER 2018; Revised 7 DECEMBER 2018. ©2019 Chinese Institute of Electronics  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is the most thermally and chemically stable, while the other phases are in a metastable state and transform into  $\beta$ -phase at adequately high temperatures.

 $\beta$ -Ga<sub>2</sub>O<sub>3</sub> has a wide bandgap of ~ 4.9 eV<sup>[3, 4]</sup> at RT. The lattice parameters of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> were firstly estimated by Kohn *et al.*<sup>[5]</sup>, and its structure was determined by Geller<sup>[1]</sup> in space group C2/m. Walton and Chase<sup>[6]</sup> clarified that the morphological symmetry appeared to be lower than that of the accepted point group 2/m.

As noted by Geller<sup>[1]</sup>, there were several other interpretations regarding  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> but none had indicated a deviation from a monoclinic structure. Furthermore, an accurate investigation of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal structure was presented by Ahman *et al.*<sup>[7]</sup> and Janowitz *et al.*<sup>[8]</sup>. The X-ray diffraction symmetry and refinement of the structure clearly indicted a C-centered monoclinic cell with space group C2/m. Fig. 1 presents a unit cell that contains two nonequivalent Ga atoms: one with tetrahedral Ga(I) and the other with octahedral Ga(II). The oxygen atoms are arranged in a distorted cubic closely packed array and have three non-equivalent oxygen ions [O (I), O (II), and O (III)] that are situated at 4i(*x*, 0, *z*) positions<sup>[9]</sup>.

Thus far, the density functional theory (DFT) is the most suitable for calculating the electronic band structure of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. Although, several theoretical studies on the electronic structure of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> have been published<sup>[10–15]</sup>, they presented a similar calculation for DFT (Fig. 2).

## 3. Ga<sub>2</sub>O<sub>3</sub> growth methods

In this section, we present an overview of the bulk crystal growth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. The following are the most common methods of bulk Ga<sub>2</sub>O<sub>3</sub> crystal growth:

·Verneuil method
·Floating zone (FZ) method
·Czochralski (CZ) method
·Edge-defined film-fed growth (EFG) method
·Bridgman method



Fig. 1. (Color online) Unit cell of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, which possesses two nonequivalent Ga sites: Ga (I), Ga (II), and three nonequivalent O-sites. Projection of the unit cell of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> along the *c*- (A), *a*- (B) and *b*-axes (C). The figure was adopted from Ref. [9].



Fig. 2. (Color online) Electronic band structure of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub><sup>[15]</sup>.

#### 3.1. Verneuil method

The Verneuil method was developed by Auguste Verneuil in 1902 and is also known as flame fusion. This process is used mainly to fabricate the gem varieties of sapphire and ruby. The steps of the operation include melting the powder material using an oxyhydrogen flame, and thereafter, the melted material crystallizes into boule (as shown in Fig. 3).

In 1964, Chase<sup>[17]</sup> prepared the first  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystal boules using the Verneuil method. The crystals grow to the desired size by increasing the gas flow and powder feed; the obtained boules are 3/8 inch in diameter and 1 inch in length. The growth direction of the boule was mostly parallel to the crystallographic *b* axis.  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> has two cleavage planes (100) and (001). Thin slabs of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> can be prepared by slitting the crystal on these planes. In addition,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals in the direction [100] have a twin plane and the majority of the crystals exhibit fine laminar twinning.

In addition, Lorenz *et al.*<sup>[18]</sup> found that the Verneuil method permits a redox reaction in the growth process.  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals were insulating and colorless when grown in an oxidizing condition. In contrast, crystals grown under reducing conditions were bluish. Mg- and Zr-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals were grown by Harwig *et al.*<sup>[19, 20]</sup> using the same method. They demonstrated that the crystals were light blue and blue when doped with Mg and Zr, respectively<sup>[19]</sup>.

#### 3.2. Floating zone (FZ) method

In 1955, Theurer<sup>[21]</sup> modified the method developed by Pfann at Bell Lab for germanium, which is known as the FZ method. This method relies on moving a liquid zone through



Fig. 3. (Color online) (a) Sketch of an early furnace used by Verneuil. (b) Simplified diagram of Verneuil process for synthesizing Ga<sub>2</sub>O<sub>3</sub><sup>[16]</sup>.



Fig. 4. (Color online) Schematic of float zone single crystal growth.

the material, and if properly seeded, a single crystal may be obtained, as shown in Fig. 4. The most prominent feature of this method is that no crucible is required, which facilitates the growth of congruent as well as incongruent melting materials.

Villora *et al.*<sup>[22]</sup> grew a single crystal  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> of approximately 1 inch in three crystallographic directions: [100], [010], and [001] (Fig. 5).

Moreover, Zhang *et al.* <sup>[23]</sup> grew crystals with a 6-mm diameter and 20-mm length using the same method and also found that the preferred crystallographic growth direction is <010>.

The major drawback of this method is that it is difficult to prevent the liquid from collapsing. It can only be held in place by surface tension. Cracks occur in the cooling crystal owing to the stress caused by high thermal gradients.

#### 3.3. Czochralski method

The Czochralski method is a process of growing single crystals of metals (Pd, Pt, Ag, Au, etc.), semiconductors (Si, Ge, and GaAs), and artificial gemstones (Fig. 6).

In 1915, Czochralski invented this method while investigating the rate of metal crystallization<sup>[24]</sup>. The most prominent feature of this method is the growth of large boules or cylindrical ingots of a single crystal. Some scientists believe the Czochralski method is the "inverse" of the Verneuil technique. The following steps briefly describe the work process involved. High-purity material is melted within a crucible, and a seed on a rod is then immersed in the liquid. The rod is pulled upwards and rotated simultaneously.

The diameter of the crystal increased to approximately 2 inch. The growth direction is always along the <010> crystallographic direction, which is parallel to both cleavage planes (100) and (001).



Fig. 5. (Color online) As-grown crystals along the crystallographic axis (a- <100>), (b- <010>), (c- <001>]).



Fig. 6. Schematic of Czochralski method.

The first Czochralski growth of a  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystal was reported by Tomm *et al.*<sup>[26]</sup>. They utilized 10% CO<sub>2</sub> in Ar instead of O<sub>2</sub> to decrease the evaporation of molten Ga<sub>2</sub>O<sub>3</sub>. Galazka *et al.*<sup>[25]</sup> utilized similar growth conditions to grow  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals. They found that the CO<sub>2</sub> concentration in the growth atmosphere may be used for adjusting the crystal perfection and growth stability in addition to tuning the optical and electrical properties of the crystals. Galazka *et al.*<sup>[27]</sup> presented a new approach for scaling up the growth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals grown from the melt. They found that the desired larger crystal required a higher oxygen concentration in the growth of a cylindrical crystal with a 2-inch diameter (as shown in Fig. 7). When using an oxygen concentration

Furthermore, the same co-workers of Galazka in 2018<sup>[28]</sup> experimentally studied the segregation of Ce, Cr, Al-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals grown using the Czochralski method as well as the effect of doping on the optical properties. They observed that both the segregation and incorporation of Cr and Ce into the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals strongly depended on the O<sub>2</sub> concentration during the growth and on their charge state.

### 3.4. Edge-defined film-fed growth (EFG) method

Intensive experiments performed by La Belle and Mlavsky<sup>[29-31]</sup> on the growth of sapphire tubes resulted in the discovery of the edge-defined film-fed growth (EFG) technique. They observed that if the top surface of the shaper (also referred to as a die) was flat, the molten materials spread out to the edge and stopped, which implied that the melt wetted a die shaper. Hence, the diameter of the die, and not of the melt column, dictated the diameter of the filament produced (Fig. 8). Consequently, the technique was named "edge-defined film-fed growth". The major problems of this method are the geometry and the material of the die, while the major advantages of this method are the growth of a near-net shape and other complicated shapes and a higher growth rate. The growth direction was the [010] crystallographic direction, which is parallel to both the cleavage planes (100) and (001).



Fig. 7. (Color online) Semiconducting  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals of 2 cm (a, c) and 5 cm (b, d) diameter grown at low- (a, b) and high-oxygen concentrations (c, d)<sup>[25]</sup>.



Fig. 8. (Color online) Edge-defined film-fed growth method.

The first successful growth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals using the EFG method was reported by Aida<sup>[32]</sup> *et al.* wherein the rate of growth was 10 mm/h for crystal ribbons with 50-mm width, 70-cm length, and 3-mm thickness. Mu *et al.*<sup>[33]</sup> grew bulk  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals having a width of 1 inch using an optimized EFG method under an Ar and 50% CO<sub>2</sub> atmosphere. The full width at half maximum of the X-ray rocking curve was only 43.2 arcsec, indicating high-quality single crystals.

Kuramata *et al.*<sup>[34]</sup> reported on the growth of large and high-quality  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals using the EFG method (Fig. 9(a)). They obtained a semiconductor substrate containing no twin boundaries with a diameter of up to 4 inches (Fig. 9(b)). The etch pit density revealed that the dislocation density on the ( $\overline{2}01$ ), (010), and (001) oriented wafers was in the order of 10<sup>3</sup> cm<sup>-2</sup>. Moreover, Kuramata *et al.*<sup>[35]</sup> also showed the growth of large crystals of up to 6 inches by the EFG method.

Chinese scientists also used this method to grow bulk crystals, Zhang *et al.*<sup>[36]</sup> grew 2-inch crystal with an atmosphere of 20% argon (Ar) and 80% carbon dioxide (CO<sub>2</sub>) (Fig. 9(c)), and FWHM of rocking curve was only 19.06 arcsec. Our group also reported 2-inch crystal growth (Fig. 9(d) published in CCCG-18 conference).

### 3.5. Bridgman method (vertical or horizontal)

The Bridgman–Stockbarger method is named after the Harvard physicist, Percy Williams Bridgman (1882–1961) and MIT physicist, Donald C. Stockbarger (1895–1952)<sup>[37, 38]</sup>. This method depends on the directional solidification process controlled by the movement of the crucible containing the molten to the axial temperature gradient in a vertical/horizontal furnace. Initially, the polycrystalline material is required to be

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Fig. 9. (Color online) (a) Photograph of EFG-grown  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulk crystal. (b) Photographs of processed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> substrates. (c) 2-inch wafer by Electronic Material Research Institute of Tianjin. (d) 2-inch plate by Shanghai Institute of Optics and Fine Mechanics.



Fig. 10. (Color online) Bridgman technique.

melted completely in the hot zone, and then, it should be slowly cooled from the end of a seed crystal. A single crystal is gradually formed along the direction of the container length (Fig. 10).

Part of the seed will be re-melted after contact with the melt. One feature of this method of growing  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is the utilization of a platinum–rhodium (70%–30%) crucible<sup>[39, 40]</sup>, which can weather a high oxygen concentration.

Hoshikawa *et al.*<sup>[39]</sup> studied the effect of the crucible material on the shapes of the grown  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals on using the vertical Bridgman (VB) method. They found that  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals were grown in the direction perpendicular to the (100) faceted plane using a seedless growth method in ambient air with a diameter of 25 mm and without adhering to the wall of the vertical cylinder crucible or a conical crucible (Fig. 11).

Ohba *et al.*<sup>[40]</sup> studied the structural defects on (100) wafers in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals approximately 25 mm in diameter grown by directional solidification in a vertical Bridgman furnace. They observed no regions of high dislocation density near the periphery of the (100) wafers cut as cross sections from the crystal ingots. Moreover, they detected line-



Fig. 11. (Color online)  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals were grown using the VB method in ambient air. (a–a') Crystal was grown in a full-diameter crucible, and (b–b') crystal was grown in a conical crucible. The figure was adopted from Ref. [38].



Fig. 12. (Color online) Schematics of line-shaped defects. The figure was adopted from Ref. [40].

shaped defects that were  $20-150 \,\mu$ m in length in the <010> direction using optical microscopy, as shown in Fig. 12. These defects are considered to be the same as those observed in the previously used EFG method.

In summary, the crucible-free FZ and Verneuil methods facilitate the use of a high oxygen concentration and a reduced amount of residual impurities. The Czochralski method provides a large crystal size of approximately 2 inch in diameter, while the EFG method provides a 6-inch-width crystal of the highest quality.

 $Ga_2O_3$  substrates underwent a massive breakthrough owing to their low cost and low energy consumption. The following organization chart (Fig. 13) depicts the comparison of the melt growth methods for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulk single crystals.

In addition, the best method for the mass production of substrates is the EFG method, which has a good track of the availability of large area, low defect density, high crystal quality.

# 4. Doping of $\beta$ -Ga<sub>2</sub>O<sub>3</sub>

The conductivity control through doping and the curtailing of the trap states is the solution to obtaining  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> for technological applications<sup>[41]</sup>. In this section, we present an overview of high-quality doping.

## 4.1. Donor doping of $\beta$ -Ga<sub>2</sub>O<sub>3</sub>

#### 4.1.1. Unintentional self-doping

A  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal exhibits n-type conductivity with no intentional doping. The n-type semiconductivity is commonly attributed to oxygen vacancies, which are ionized and form donors. Therefore, it is assumed that there is a strong correlation



Fig. 13. The comparison of the melt growth methods for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulk single crystals.

between the conductivity of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals and oxygen vacancies. This assumption has been debated by Varley *et al.*<sup>[15]</sup> who performed first-principle calculations for various impurities and oxygen vacancies in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. Based on calculations, they suggested that the oxygen vacancy serves as a deep donor with an ionization energy of more than 1 eV and, thus, cannot contribute to n-type. They proposed that hydrogen is the probable cause of the electrical conductivity.

Zhang *et al.*<sup>[23]</sup> produced  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals using the FZ method. They found that the 300-nm emission can be ascribed to the potential wells formed by acceptor clusters, and the 692-nm emission was attributed to the recombination of holes and electrons via the acceptors and donors in the vicinity.

#### 4.1.2. Intentional n-type doping

The fabrication and characterization of Sn-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals were studied by Suzuki et al.<sup>[42]</sup> and Ueda et al.[43] via the FZ method. High-conductivity crystals were obtained using the FZ method using Ga<sub>2</sub>O<sub>3</sub> rods doped with 2-10 mol% SnO<sub>2</sub>. The low incorporation efficiency causes a massive amount of Sn atoms to evaporate during the growth process. Ueda et al.<sup>[43]</sup> reported that Sn-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> had a blue coloration, and the value of their conductivity was 0.96 Ω<sup>-1</sup> cm<sup>-1</sup> (electrical resistivity = 1.04 Ω·cm). Suzuki *et al.*<sup>[42]</sup> reported that the optimized specimen was 32 ppm. The conductivity value is 23.4  $\Omega^{-1}$ ·cm<sup>-1</sup> (electrical resistivity = 4.27  $\times$  $10^{-2} \Omega \cdot cm$ ), and the carrier concentration is approximately  $2.26 \times 10^{18}$  cm<sup>-3</sup>. Zhang *et al.*<sup>[23]</sup> studied the effect of Sn doping in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals prepared using the FZ method and that of the annealing process on their conductivity and optical properties. They showed that the forbidden bandwidth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is influenced by Sn<sup>4+</sup> doping. Both the absorption coefficient and conductivity increase on increasing the Sn doping. A reduction in the absorption coefficient after the annealing refers to the decrease in the free carriers. Thus, the annealed process changes the sample into an insulator. Moreover, Ohira et al.[44] examined the effect of the annealing process on the surface

morphology, electrical and optical properties of Sn-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> produced using the FZ method. It was observed that these properties did not change, however the lattice parameters increased slightly after annealing at 1100 °C. Furthermore, Sn was uniformly dispersed in the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> matrix.

VÍIlora *et al.*<sup>[45]</sup> obtained Si-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> using the FZ method. It was found that the Si concentration is in the range of 10<sup>16</sup>–10<sup>18</sup> cm<sup>-3</sup>, and the conductivity increases contentiously from 0.03 to 50  $\Omega^{-1}$ cm<sup>-1</sup>. At a low doping rate, nearly all the Si is incorporated into the crystals, while at a high doping rate, the Si concentration in the crystal is approximately 5% that of the feed rod. While the mobility fluctuated around 100 cm<sup>2</sup>/V, the free carrier concentration increased by three orders of their magnitude. Therefore, the electrical conductivity depends sustainably on the main impurity. This is in agreement with the theoretical calculations<sup>[15]</sup>, which shows that the oxygen vacancies in the crystal are deep donors and can hardly contribute to the electron conductivity.

Sasaki et al.<sup>[46]</sup> published good results on Si-ion implantation into Ga<sub>2</sub>O<sub>3</sub> crystals fabricated using the FZ method. The Si concentration varies from  $1 \times 10^{19}$  to  $1 \times 10^{20}$  cm<sup>-3</sup> owing to the varying in the total dose from  $2 \times 10^{14}$  to  $2 \times 10^{15}$  cm<sup>-2</sup>. The sample with a Si concentration of  $5 \times 10^{19}$  cm<sup>-3</sup> had a resistivity of 1.4 mΩ·cm. A high electrical activation efficiency of over 60% was obtained in the samples with a Si concentration of  $\leq 5 \times 10^{19}$  cm<sup>-3</sup> through the annealing process in the range of 900–1000 °C. Kuramata et al.[34] reported the growth of Sidoped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulk crystals using the EFG and FZ methods and investigated the residual impurities, controllable doping, crystal defect, and the effect of thermal annealing on the electrical properties of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulk crystals. They observed that the main residual impurity in the EFG crystals is Si and that the Si content controlled the effect donor concentration  $(N_{\rm d} - N_{\rm a})$  of unintentionally doped crystals. Therefore, obtaining an intentionally n-doped crystal was found to be possible by doping with SiO<sub>2</sub>.  $N_{\rm d} - N_{\rm a}$  was approximately the same as the Si concentration in the samples annealed in nitrogen,



Fig. 14. (Color online) Theoretical ideal performance limits of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> power devices in comparison with those of other major semiconductors. The figure was adopted from Ref. [58].

while it was independent of the Si concentration under oxygen ambient.

Zhou *et al.*<sup>[47]</sup> experimentally confirmed that Nb element is a good choice to control the conductivity of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. Nbdoped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals were grown using the FZ method, and their structure was that of monoclinic  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. The electrical resistivity decreased from 3.6 × 10<sup>2</sup> to 5.5 × 10<sup>-3</sup>  $\Omega$ ·cm and the carrier concentration increased from 9.55 × 10<sup>16</sup> to 1.8 × 10<sup>19</sup> cm<sup>-3</sup> on increasing the doping.

### 4.2. Acceptor doping of $\beta$ -Ga<sub>2</sub>O<sub>3</sub>

A major challenge for of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is its lack of an adequate method for fabricating p-type  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals. Few reports on p-type Ga<sub>2</sub>O<sub>3</sub> have been published<sup>[48]</sup>.

#### 4.2.1. Unintentional self-doping

Varley *et al.*<sup>[49]</sup> studied the effect of holes in many bandgap oxides and found that the holes in  $Ga_2O_3$  can form self-trap holes with an energy of 0.53 eV. This indicates that free holes will be precarious and tend to localize as small polarons. Kananen *et al.*<sup>[50]</sup> observed self-trapped holes in  $Ga_2O_3$ using the electron paramagnetic resonance technique. These defects are thermally stable below ~80–100 K.

#### 4.2.2. Intentional p-type doping

Onuma et al.<sup>[51]</sup> obtained Mg-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals using the FZ method and found the samples' behavior to be semi-insulating with a resistivity of  $6 \times 10^{11} \Omega \cdot cm$  for Mg concentration of 4  $\times$  10<sup>18</sup>–2  $\times$  10<sup>19</sup> cm<sup>-3</sup>. The Blue luminescence intensity depended on the Mg doping, and there was a correlation between the intensity and the formation energy of  $V_0$ . Mgdoped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals were prepared by Galazka *et al.*<sup>[27]</sup> via the CZ method in a very low range of Mg of 6–28 ppm at 100% CO<sub>2</sub>, which indicates that Mg very efficiently compensates acceptors. Thus far, only Liu et al.[52] have reported on N-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> nanowires with a p-type electrical conductivity. As in the case of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, p-type doped ZnO is difficult to acquire, and previous studies have stated that N acceptors are strongly compensated by intrinsic defects, and N atoms cannot be effective acceptors owing to the deep acceptor levels in the bandgap of ZnO. Dong et al.[53] studied the compensation mechanism between N acceptors and native defects in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> using first-principle calculations based on DFT. They found that the N dopant serves as a deep acceptor with an acceptor level of 1.33 eV above the valence band maximum, which cannot be an effective p-type dopant. In addition, the N dopant formed four types of defect complexes including oxygen and Gallium vacancies and interstitials. Kyrtsos *et al.*<sup>[54]</sup> used DFT to investigate the various cation substitutional dopants in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> for the possibility of p-type conductivity. They found that these dopants act as deep acceptor levels with ionization energies of more than 1 eV. They can also trap an extra hole at a very deep donor level, which inhibits p-type conductivity.

## 5. Applications

In this section, we present an overview of the applications of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>.

## 5.1. Gas sensors

The working principle of oxygen sensor devices is a negative relationship between the conductivity of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and the oxygen partial pressure in the surrounding atmosphere. The conductivity of the sensor changes with varying the proportion of oxygen or the concentration of gases in the atmosphere. Bartic *et al.*<sup>[55]</sup> studied the oxygen sensitivity of Ga<sub>2</sub>O<sub>3</sub> single crystals and thin films at high temperatures. They found that the difference in response times between a single crystal and poly material was small.

#### 5.2. Power and high-voltage devices

The critical electric field can be determined from the bandgap using an empirical relationship<sup>[56]</sup>:

 $E_{\rm c} = 1.73 \times 10^5 \, (E_{\rm q})^{2.5}$ 

The estimated breakdown voltage is 8 MV/cm, which is approximately three times larger than that of SiC or GaN. The corresponding Baliga's Fig. of merit (BFOM  $\simeq \epsilon \mu E_c^3$ ), which is used to estimate the potential for power device applications, is at least four times larger than those of SiC or GaN. Furthermore, the conduction loss of Ga<sub>2</sub>O<sub>3</sub> is lower than that of other semiconductor materials<sup>[57, 58]</sup>. Thus, Ga<sub>2</sub>O<sub>3</sub> is one of the most suitable materials for power-device applications. Fig. 14 shows a comparison of the theoretical performance limits of Ga<sub>2</sub>O<sub>3</sub> with those of other major semiconductors.

#### 5.3. Schottky diodes

High-performance Schottky barrier diodes were fabricated by Oishi *et al.*<sup>[59]</sup> on an EFG-grown  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystal. The highest theoretical value of electron mobility reported thus far is 886 cm<sup>2</sup>/V·s at 85 K. The current density was 70.3 A/cm<sup>2</sup> at 2.0 V for a forward voltage, and an ideality factor of 1.01 was obtained. Moreover, the Schottky barrier diodes fabricated using the FZ technique offered good characteristics, a high reverse breakdown voltage of approximately 150 V, and a barrier height of the Pt/ $\beta$ -Ga<sub>2</sub>O<sub>3</sub> interface of 1.3–1.5 eV<sup>[60]</sup>.

Recently, the highest performance of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> Schottky barrier diode (SBD) was reported by Xidian University<sup>[61]</sup>, with a reverse blocking voltage of more than 3 kV and a low DC specific ON-resistance (R<sub>ON</sub>,sp) of 24.3 m $\Omega$ ·cm<sup>2</sup>.

#### 5.4. Field effect transistors (FET)

Higashiwaki *et al.*<sup>[62]</sup> first demonstrated an interesting design of a metal semiconductor field-effect transistor (MES-FET) using Ga<sub>2</sub>O<sub>3</sub> single crystals. They fabricated n-channel Ga<sub>2</sub>O<sub>3</sub> MESFETs on a single-crystal Mg-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> (010) substrate. A Sn-doped (100) semi-insulating  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> grown onto a Mg-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal was fabricated by Green *et al.*<sup>[63]</sup> using Metal organic chemical vapor deposition(MOCVD). They reported on Ga<sub>2</sub>O<sub>3</sub> metal-oxide-semiconductor field effect transition of the semi-insulation of the semiconductor field effect transition of the semic



Fig. 15. (Color online) Schematic illustration: (a) cross-section and (b) optical micrograph of  $Ga_2O_3$  MESFET. According to the description from Ref. [62].

sistors (MOSFETs) with a gate-to-drain electric field corresponding to 3.8 MV/cm, which is the highest reported for any transistor surpassing bulk GaN and SiC theoretical limits<sup>[62]</sup>. A cross-section of the Ga<sub>2</sub>O<sub>3</sub> MESFET structure is shown in Fig. 15.

## 5.5. UV detectors

Photodetectors are one of the simplest devices to produce as they require only Schottky or Ohmic contacts. The major application of these devices is solar-blind ultraviolet (UV) radiation detection. Owing to its wide bandgap,  $Ga_2O_3$  provides a unique opportunity for the detection of UV radiation of wavelengths <280 nm.

Suzuki *et al.*<sup>[64]</sup> fabricated  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> photodiodes with an Au Schottky contact on a single crystal substrate. They also investigated the effect of post heat treatment on the electrical and optical properties of photodiodes. The ideality factor improved to nearly 1 subsequent to annealing at temperatures greater than 200 °C; however, the reverse leakage current remained nearly unchanged. The high responsivity and biasvoltage dependence of the photocurrent implies the presence of an internal gain. Yang *et al.*<sup>[65]</sup> fabricated a Schottky barrier diode solar-blind photodetector based on a single crystal  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. They found that the device has a clear response to the solar-blind wavelength at a zero bias, which confirms that it can be used as a self-powered solar-blind photodetector.

#### 5.6. Radiation detection

Wide bandgap semiconductors feature strong bonds that provide them with an intrinsically high radiation resistance<sup>[41]</sup>. Patrick *et al.*<sup>[66]</sup> investigated Ga<sub>2</sub>O<sub>3</sub> as a radiation detection material for fast (14 MeV) neutrons, using the <sup>16</sup>O (n,  $\alpha$ ) <sup>13</sup>C reaction. Yang *et al.*<sup>[67]</sup> exposed vertical rectifiers fabricated on epi Ga<sub>2</sub>O<sub>3</sub> on bulk  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> to 1.5-MeV electron irradiation at fluences from 1.79 × 10<sup>15</sup> to 1.43 × 10<sup>16</sup> cm<sup>-2</sup> at a fixed beam current of 10<sup>-3</sup> A. The electron irradiation led to a reduction in the carrier concentration of Ga<sub>2</sub>O<sub>3</sub> with a carrier removal rate of 4.9 cm<sup>-1</sup>. Szalkai *et al.*<sup>[68]</sup> illustrated the capability of Mg-doped  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystals in detecting fast neutrons. They observed a gamma signal from a Ga decay process when a neutron-irradiated  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal sample was tested using a highpurity germanium detector.

## 6. Conclusions and future perspectives

This review summaries the state of the art of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and achievements with respect to  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> ranging from preparation and characterization to applications. Although significant progress has been made in the growth of bulk crystals, defect and impurity control require further studies. As can be observed from this article, the published data on Ga<sub>2</sub>O<sub>3</sub> are still incomplete and controversial to some extent. We look forward to realizing the following objectives that may be important research and development paths with respect to  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> bulkbased devices.

(1) Identification of the dominant defects in bulk crystals and their effect on device performance

(2) Experimental efforts to address open issues such as ptype conductivity and realizing the full potential of  $Ga_2O_3$  materials in device applications

(3) New concepts and designs for improving the thermal management and reliability of power electronics devices

(4) Further progress in growth of large-diameter high-quality crystals

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