

# Modeling the degradation mechanisms of AlGaIn-based UV-C LEDs: from injection efficiency to mid-gap state generation

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In this work, we analyze and model the effect of a constant current stress on an ultraviolet light-emitting diode with a nominal wavelength of 285 nm. By carrying out electrical, optical, spectral, and steady-state photoluminescence (SSPL) analysis during stress, we demonstrate the presence of two different degradation mechanisms. The first one occurs in the first 1000 min of stress, is ascribed to the decrease in the injection efficiency, and is modeled by considering the defect generation dynamics related to the de-hydrogenation of gallium vacancies, according to a system of three differential equations; the second one occurs after 1000 min of stress and is correlated with the generation of mid-gap defects, for which we have found evidence in the SSPL measurements. Specifically, we detected the presence of deep-level states (at 1.6 eV) and mid-gap states (at 2.15 eV), indicating that stress induces the generation of non-radiative recombination centers. © 2020 Chinese Laser Press

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## 1. INTRODUCTION

Over the last years, AlGaIn-based UV-C LEDs have shown impressive advancements, thanks to the high volume of research done by several research and industrial laboratories and to the broad range of expected applications. UV-C LEDs can be used in different fields, including medical devices, disinfection, sterilization, water purification and UV curing [1–3], plant lighting [4], as insect trap [5], and in phototherapy and for bioagent detection and identification [6,7]. Moreover, UV-C LEDs present many advantages with respect to conventional UV lamps, such as higher spectral purity, smaller size, smaller drive voltage and operating current, and environmental friendliness.

The interest in UV-C technology is now further pushed by the COVID-19 emergency: specific antiviral treatments are being developed [8] (often with mercury or excimer lamps), and the development of high-efficiency UV-C light sources could ease the spread of compact disinfection systems.

The reliability and the efficiency of AlGaIn-based LEDs were significantly improved thanks to recent studies, resulting in higher injection efficiency, internal quantum efficiency, and light extraction efficiency. Nevertheless, several topics are still the subject of investigation, including the increase in drive

voltage during aging, the generation of mid-gap states in the active region [9], the incomplete activation of p-dopant [10,11], the migration of hydrogen [12] and, in some cases, catastrophic failures [13].

The aim of this paper is to analyze the different degradation mechanisms of 285 nm UV-C LEDs, through electrical, optical, spectral, and steady-state photoluminescence (SSPL) measurements. Through our investigation we demonstrate that degradation is due to two different mechanisms: the decrease in the injection efficiency and the generation of mid-gap Shockley–Read–Hall (SRH) recombination centers. The underlying mechanisms are discussed in detail in the paper.

## 2. EXPERIMENTAL RESULTS

The samples under investigation are AlGaIn-based UV-C LEDs with a nominal wavelength of 285 nm at the current of 350 mA. We performed a constant current stress at room temperature, at 250 mA; the stress test was repeatedly interrupted to carry out electrical ( $I$ - $V$ ), optical ( $L$ - $I$ ), and power spectral density (PSD) characterization at different temperatures, from 15°C to 75°C with a step of 10°C, in order to evaluate the behavior of the main parameters of the devices. At each stage of the stress test we also performed an SSPL scan at room

temperature, to detect the presence of defects and to investigate their evolution during the stress test. The current-voltage measurements were carried out from  $-4$  V to  $7.5$  V; the optical power measurements were carried out from  $10$   $\mu$ A to  $250$  mA, with a logarithmic step. The spectral characterization was performed at four different current levels ( $100$   $\mu$ A,  $1$  mA,  $10$  mA,  $100$  mA), in a range of wavelength from  $180$  nm to  $875$  nm. SSPC measurements were carried out by using a monochromatic light in a range from  $1.1$  eV to  $3.65$  eV, while keeping the LEDs biased at  $0$  V.

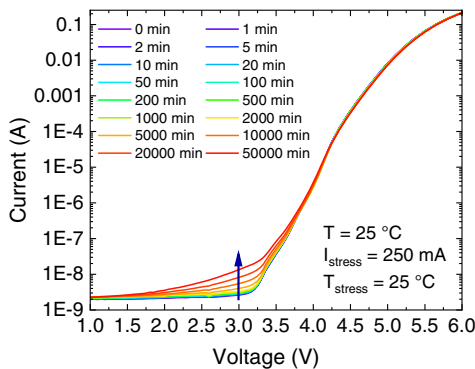
Figure 1 reports the  $I$ - $V$  measurements carried out during the constant current stress experiment. Two regions can be identified: in the first one, below the turn-on voltage ( $V < 4.5$  V), conduction is mediated dominantly by defects. Before stress, the current is smaller than  $2$  nA. This is the sensitivity limit of the experimental setup used for the investigation, that employs a specific circuitry to switch between electrical, optical, and capacitive measurements. For longer stress time ( $>1000$  min), a significant increase in current conduction is observed for  $1$  V  $< V < 3$  V, i.e., below the turn-on voltage of the pn junction. This effect is ascribed to the increase in the density of defects in the active region of the devices.

These results are confirmed also by the data reported in Fig. 2, which show a significant increase in the sub-turn-on current after  $1000$  min of stress [Fig. 2(a)] and a small decrease in the series resistance of the device [Fig. 2(b)].

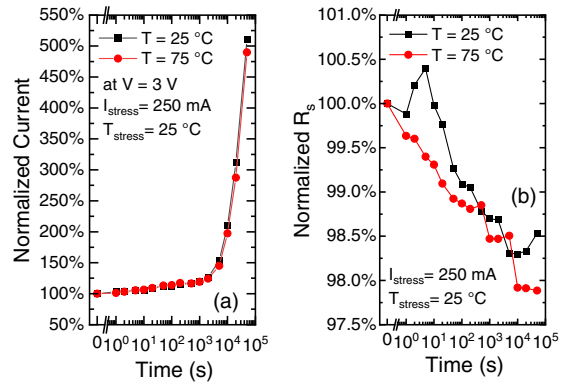
We ascribed the increase in the sub-turn-on current at the generation of mid-gap states during the stress, confirmed also by the calculation of the ideality factor (Fig. 3), which is an indicator of the increase in the defectiveness of the material from the  $I$ - $V$  result [14,15]. This generation of mid-gap states leads to an increase in non-radiative recombination events through trap-assisted mechanisms [16], which cause the reported increment in current.

With regard to the optical degradation, Fig. 4 reports the optical power versus current characteristics measured during stress, and the relative variation in optical power at three different current levels. Different behaviors are observed, depending on the conduction regime.

We start considering the behavior of the devices as observed for high current levels ( $100$  mA), i.e., in the regime where the behavior of the devices is not significantly influenced by SRH recombination. We suggest that the optical degradation



**Fig. 1.** Electrical characterization during the stress, in semi-logarithmic scale, carried out before and during the stress experiment at  $250$  mA. All measurements were taken at  $25^\circ\text{C}$ .



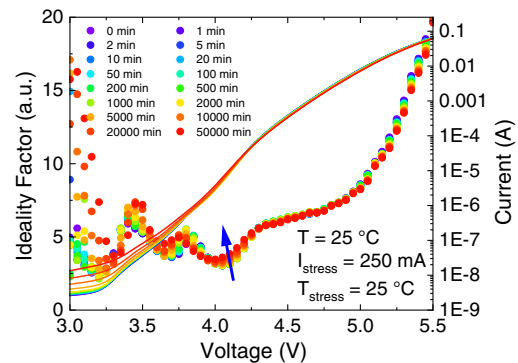
**Fig. 2.** (a) Normalized current at the voltage of  $3$  V, and (b) normalized series resistance ( $R_s$ ) at the temperatures of  $25^\circ\text{C}$  and  $75^\circ\text{C}$  during the stress.

detected at high current levels is due to a worsening of the injection efficiency. This is consistent with the fact that up to  $1000$  min of stress the degradation rate is almost independent on the measuring current level [see Fig. 4(b) and Ref. [17]].

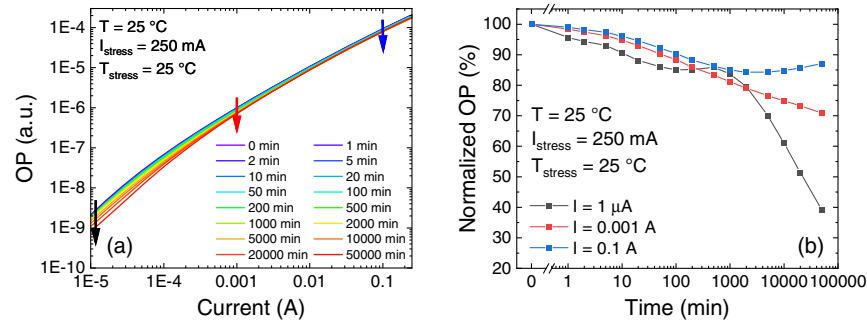
Several mechanisms can contribute to this variation in injection efficiency. A first process is the decrease in hole concentration at the p-side, due to a partial compensation of the acceptor dopant through the formation of Mg-H bonds [18,19]. This process is not supposed to play a major role, since it would result in a significant increase in the series resistance, which is not observed here [as shown in Fig. 2(b), series resistance decreases with stress time].

A second mechanism that impacts on the injection efficiency is carrier escape [20]. However, escape is a thermionic process, following an equation like  $J_n = -q \frac{D_n}{L_n} \cdot N_C \cdot \exp\left(-\frac{E_B - E_{F_n}}{kT}\right)$ , where  $D_n$  is the diffusion constant,  $L_n$  is the diffusion length,  $N_C$  is the effective density of states in the active region,  $E_B$  is the height of the barrier, and  $E_{F_n}$  is the Fermi level [20]. The rate of escape mainly depends on the height of the confinement barrier at the p-side of the quantum wells, which is not supposed to change as a consequence of stress.

A third mechanism that can lower the injection efficiency is the generation of charged defects in the active region [21]. Such charged defects may modify the band bending near the quantum wells (QWs), thus making carrier injection more difficult (Fig. 5).



**Fig. 3.** Ideality factor at the temperature of  $25^\circ\text{C}$  during the stress.



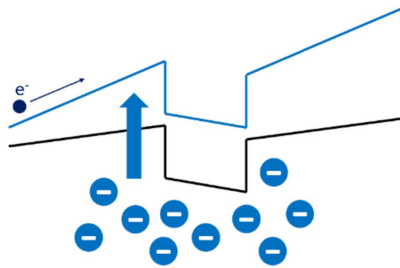
**Fig. 4.** (a) Optical power (OP) during the stress at the temperature of 25°C. (b) Normalized optical power at three different current levels: 10 μA, 1 mA, and 100 mA.

As a plausible model, we consider that stress induces the generation of negatively charged defects near/within the active region. A possible process can be the de-hydrogenation of (previously) hydrogenated gallium vacancies, which are naturally present in MOCVD-grown (metal organic chemical vapor deposition) gallium nitride.

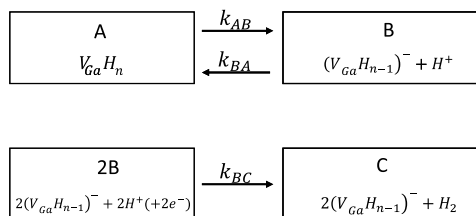
As discussed in Ref. [22], de-hydrogenation of Ga vacancies,



leads to an increase of acceptor concentration. This reaction may be promoted by the flow of carriers through the active region of the devices, and/or by the energy released by non-radiative recombination events (for instance, Auger recombination [23]). Since  $H^+$  ions have a high mobility, they may then re-bond with a gallium vacancy or find another hydrogen atom to form molecular hydrogen (with the contribution of two electrons). The processes could be described like in Fig. 6.



**Fig. 5.** Simplified representation of the increase in the injection barrier due to the presence of a distributed negative charge near/within the active region.



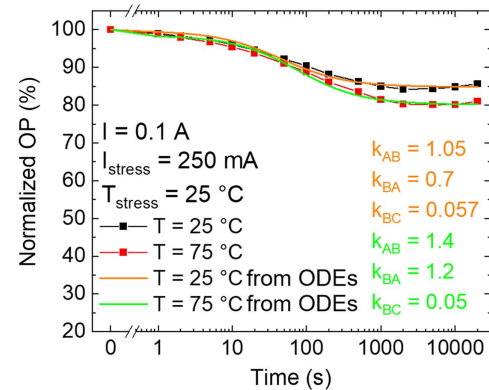
**Fig. 6.** Schematic representation of the reactions.

The Ga vacancies  $V_{Ga}H_n$  (state A) could de-hydrogenate into  $[V_{Ga}H_{n-1}]^- + H^+$  (state B, rate coefficient  $k_{AB}$ ). State B is metastable: this means that the hydrogen ion is free, and could either re-bond with a de-hydrogenated vacancy (going back to state A, with rate coefficient  $k_{BA}$ ), or find a second hydrogen atom to form molecular hydrogen  $H_2$  (going to state C, with rate coefficient  $k_{BC}$ ). In this second case, the hydrogen finds a stable form, and leaves behind a de-hydrogenated vacancy  $[V_{Ga}H_{n-1}]^-$ , that modifies the injection efficiency, as described above. These processes could be described by the following system of ordinary differential equations (ODEs):

$$\begin{cases} \frac{dA(t)}{dt} = -k_{AB} \cdot A(t) + k_{BA} \cdot B(t) \\ \frac{dB(t)}{dt} = +k_{AB} \cdot A(t) - k_{BA} \cdot B(t) + k_{BC} \cdot B^2(t). \\ \frac{dC(t)}{dt} = +k_{BC} \cdot B^2(t) \end{cases}$$

Here  $A$  represents the number of hydrogenated vacancies,  $B$  is the number of de-hydrogenated vacancies with neighboring H-atom ( $[V_{Ga}H_{n-1}]^- + H^+$ ), and  $C$  is the number of de-hydrogenated vacancies without H-atom (acting as acceptors).

After defining this system of equations, we solved them numerically, after imposing as a boundary condition that the initial value of  $A$  is equal to 1. The coefficients  $k_{ij}$  were obtained numerically, as fitting parameters. As shown in Fig. 7, we could



**Fig. 7.** Optical degradation measured at 25°C and 75°C during stress at 250 mA. Solid lines represent the solution of the system of ODEs reported above, showing a good agreement with the experimental data.

obtain a very good correspondence between the estimated variation in de-hydrogenated vacancies  $[V_{Ga}H_{n-1}]^-$ , described by the variation of term  $[C(t)]$  and the overall variation in optical power.

Thus, we conclude that the proposed model is able to predict the variation in optical power during stress time for high measuring current levels, i.e., in the regime where the variation in injection efficiency is the dominant degradation process.

For low current levels, a different trend is observed. Specifically, after the first phase ( $t < 1000$  min) has ended, a second mechanism is observed, for  $t > 1000$  min. This second process has a stronger effect at low measuring current levels, and it is ascribed to an increase in the non-radiative recombination rate, due to the generation of defects.

We can explain the behavior at low current density, with the hypothesis formulated by Ruschel *et al.* [24], where the authors indicated that the optical power has the following behavior for long operation time:

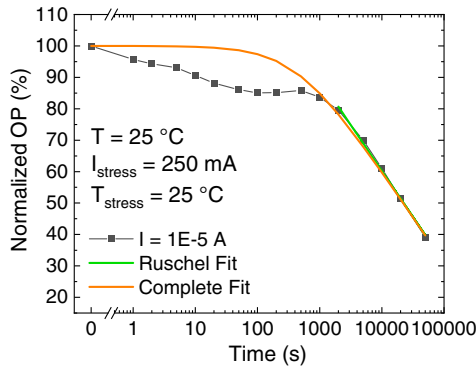
$$OP(t) = -\beta \cdot \ln\left(\alpha \cdot J^3 \cdot t + e^{-\frac{1}{\beta}}\right),$$

where  $J$  is the current density,  $t$  is the stress time, and  $\alpha$  and  $\beta$  are two fitting parameters. This is in agreement with the results in Fig. 8 that report the optical degradation and the related fit for long time. Ruschel *et al.* [24] also indicated that the degradation rate scales with the cube of the current density, suggesting a possible contribution of Auger recombination in the degradation.

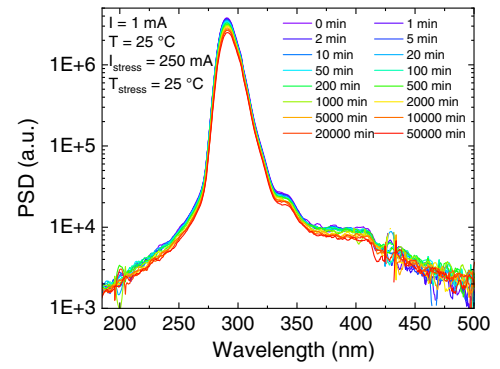
A possible scenario for long stress time considers the generation of non-radiative recombination centers due to the energy released by Auger recombination events. The consequent increase in SRH recombination coefficient has an impact, which is stronger for low measuring current levels, as predicted by the well known ABC rate equations. The contribution of this mechanism is negligible in the first phase of the stress, compared to the effect given by the drop in injection efficiency described above.

So, the long-term degradation measured at low current levels can be ascribed to the generation of non-radiative centers, possibly caused by the energy released by Auger recombination events.

In Fig. 9 we report the PSD of the device during the aging. We can see the presence of a parasitic shoulder peak in the range from 320 nm to 350 nm, probably due to recombination



**Fig. 8.** Fitting of the optical power data at low current level with the function proposed in Ref. [24].



**Fig. 9.** Power spectral density during the aging at the current of 1 mA and at the temperature of 25°C.

in the proximity of the electron blocking layer, caused by the carrier escape from the QWs and the carrier overflow [25]. No major change in the electroluminescence spectrum is observed after stress, apart from a decrease in the emitted signal. In particular, no additional parasitic emission was noticed after aging. This is a further confirmation of the non-radiative nature of the process responsible for degradation.

According to the SRH theory, the defects that have a stronger effect on the optical characteristics of LEDs are those located near the mid-gap [20]. Similarly, recent papers indicated that the defects that have the stronger impact on the sub-turn-on leakage components are located close to the mid-gap [15,16]. For this reason, we carried out an extensive analysis to identify the generation of traps near the active layer of the devices, by capacitance spectroscopy. To this aim, normal deep-level transient spectroscopy cannot be used, since the time constant of the capacitance transients would be too long for practical implementation of the technique.

In fact, the time constant for emission can be calculated as

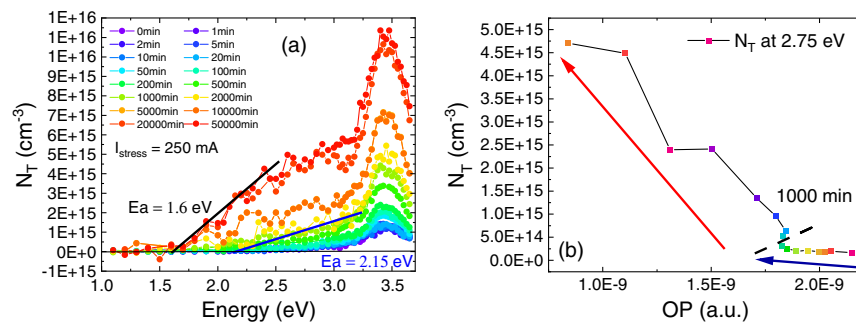
$$\frac{1}{\tau} = e_n = \sigma_n v_{th} N_c \exp\left(-\frac{E_C - E_T}{kT}\right).$$

With  $\sigma_n = 10^{-15} \text{ cm}^2$ ,  $v_{th} = 1.3 \times 10^7 \text{ cm/s}$  and  $N_c = 4.8 \times 10^{18} \text{ cm}^{-3}$ , time constants in excess of  $1.2 \times 10^{11} \text{ s}$  would be obtained even at high temperature (500 K), for mid-gap defects.

For this reason, we decided to opt for SSPC measurements. This method uses monochromatic light excitation to favor a fast ionization of defects, thus particularly effective for the study of very deep defects in wide bandgap semiconductors.

In Fig. 10 are reported the results of steady-state photocapacitance measurements during the stress. The results show two different slope increments, correlated to the presence and the generation of two different defects. The first one, visible at 1.6 eV, is present at each step of stress. The second, at 2.15 eV, is more evident in the first steps of stress and corresponds to defects placed at the mid-gap for the dominant QW bandgap. This supports the hypothesis of a generation of mid-gap states in the active region, which cause an increase in the SRH recombination. In fact, as reported in Fig. 10(b), there is a good correlation (red arrow) between the decrease in the optical





**Fig. 10.** (a) SSPC measurement during the aging, and (b) correlation between the second defect from SSPC and the optical power at low current levels.

power at low current levels and the increase in the concentration of mid-gap defect (proportional to the amplitude of the SSPC signal) after 1000 min of stress, probably ascribed to Auger processes.

### 3. CONCLUSIONS

In this work, we presented an extensive analysis of the behavior of UV-C LEDs during the constant current stress test. We found the presence of two different degradation mechanisms.

The first one is dominant in the first 1000 min of stress, and it is ascribed to a decrease in the injection efficiency. It causes a decrease in the optical power and in the amplitude of the main peak of the emission spectra, which is independent of the measuring current level. The related kinetics were modeled by considering the defect-reaction dynamics responsible for the de-hydrogenation of gallium vacancies, through a system of three differential equations. The results show a good agreement with the experimental data, supporting the validity of the model.

The second mechanism is dominant after 1000 min of stress, and it is correlated to the generation of mid-gap defects. It leads to an increase in the sub-turn-on current, and a decrease in the optical power and in the PSD main peak, especially at low current levels. Moreover, we found evidence of this defect generation from the slope of the optical power characteristics at low current levels, from the ideality factor, and from the results of SSPC measurement, which indicates the generation of defects at  $E_C$  of 2.15 eV.

In a future work, microscopy techniques (transmission electron microscopy, scanning electron microscopy, energy dispersive X-ray analysis, positron annihilation spectroscopy) could be used to experimentally evaluate the impact of electrical stress on the microscopic properties of the material. Also, positron annihilation spectroscopy could be used to evaluate the role of vacancies in the degradation process.

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