PHOTONICS Research

Generation of 84-fs pulses from a mode-locked Tm:CNNGG disordered garnet crystal laser

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A mode-locked laser based on a Tm:CNNGG disordered crystal as an active medium and a single-walled carbon nanotube saturable absorber is demonstrated, operating at a central wavelength of 2018 nm. Transform-limited 84 fs pulses are generated with an average output power of 22 mW at a repetition rate of ~90 MHz. A maximum output power of 98 mW is obtained at a slightly longer pulse duration of 114 fs. © 2018 Chinese Laser Press

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1. INTRODUCTION

Ultrashort pulse laser sources emitting near 2 μm are under rapid development nowadays due to their potential in various applications, e.g., remote sensing, material processing, medical diagnostics, and laser surgery. They present efficient tools for the investigation of the dynamics of physical systems on very short time scales [1], and can be used for synchronous pumping of mid-IR optical parametric oscillators and seeding of near-degenerate optical parametric chirped-pulse amplifiers [2] employed for high-order harmonics generation [3]. However, so far, using Tm- or Ho-doped bulk materials (crystals, ceramics, or glasses), sub-100 fs pulses have been reported only with a Tm:MgWO₄ crystal (87 fs) and a Tm:(Lu_{2/3}Sc_{1/3})₂O₃ "mixed" ceramic (63 fs) [4,5]. Therefore, exploring novel solid-state laser media for the generation of ultrashort pulses near 2 μ m down to the few-cycle limit is of great importance.

In recent years, disordered host crystals have drawn much attention due to their superior spectral bandwidths and moderate thermal properties. Doped with Tm or Ho ions, they have been proved to be excellent ultrafast laser materials in the 2 μm spectral range [6–12]. In such disordered crystals, different cations can occupy randomly one or more crystallographic sites, resulting in additional inhomogeneous spectral line broadening for the optical transitions of the dopant ion.

Among such disordered hosts, calcium niobium gallium garnets Ca₃(NbGa)_{2-x}Ga₃O₁₂ (CNGG) and calcium lithium gallium garnets $Ca_3Li_xNb_{1.5+x}Ga_{3.5-x}O_{12}$ (CLNGG), with Li⁺ added to partly overcome the nonstoichiometry, have been widely investigated. Active rare-earth ions such as Nd³⁺, Yb³⁺, and Tm³⁺, can be used as dopants to substitute for the Ca²⁺ ions in dodecahedral sites. The lattice disorder in these type of crystals results from the random distribution of Nb5+ and Ga3+ ions in the octahedral and tetrahedral sites [13]. In the spectral range near 1 µm, pulse durations of 45 and 55 fs were obtained from mode-locked lasers based on Yb-doped CNGG and CLNGG [14,15], respectively, and 534, 660, and 900 fs from Nd-doped CLNGG-CNGG, CNGG, and CLNGG garnets, respectively [16–18]. In the 2 µm spectral range, such disordered crystals doped with Tm3+ have shown broad and smooth continuouswave (CW) laser tuning, which is very promising for sub-100-fs pulse generation [19]. However, only Tm:CLNGG has been investigated for mode locking: pulse durations of 354 and 479 fs were obtained applying graphene or a semiconductor saturable absorber mirror as saturable absorbers [10,11]. Considering that CNGG and CLNGG have similar spectral properties to other dopants, we decided to study the modelocked laser performance of Tm:CNGG. We additionally

introduced Na⁺ into the crystal (hereafter Tm:CNNGG) to decrease the temperature of crystal growth and promote Tm³⁺ incorporation [13,20].

In this work, we present the CW and mode-locked laser performance of Tm:CNNGG. Applying a single-walled carbon nanotube saturable absorber (SWCNT-SA) and chirped mirrors for dispersion compensation, the mode-locked Tm:CNNGG laser generated pulses as short as 84 fs at 2018 nm.

2. EXPERIMENTAL DETAILS AND DISCUSSION

The Tm:CNNGG disordered crystal used in the experiment was grown by the conventional Czochralski method with a Tm³+ doping of 3.2 at. % and Na+ doping of 5.5 at. %, corresponding to a concentration of 4.47×10^{20} and 7.59×10^{20} ions/cm³ (measured in the crystal), respectively. For the growth process, Ca₃Nb_{1.6875}Ga_{3.1875}O₁₂ was selected as the composition to maintain charge neutrality. The maximum absorption cross-section was 0.38×10^{-20} cm² at 785.8 nm (³H₆ \rightarrow ³H₄ transition) and the full width at half maximum (FWHM) of this band amounted to 29 nm. The emission due to the ³F₄ \rightarrow ³H₆Tm³+ transition extended from 1600 to 2050 nm. The highest emission cross-section, $\sigma_e = 0.58 \times 10^{-20}$ cm², was at 1866.8 nm [20].

Figure 1 shows the calculated gain cross-section σ_g . It can be seen that the gain maximum of Tm:CNNGG is located at wavelengths below 2000 nm for $\beta > 0.08$, where strong water vapor absorption in air will prevent stable femtosecond mode locking. Only for low population inversion levels, i.e., $\beta \leq 0.08$, is lasing in Tm:CNNGG expected around 2000 nm. Such wavelengths are long enough to intrinsically avoid the strong water vapor absorption. This, together with the broad and smooth gain characteristics of the spectrum, is a prerequisite for femtosecond pulse generation.

In order to investigate the mode-locked laser performance of Tm:CNNGG, a standard X-shaped cavity as shown in Fig. 2 was employed. The sample used was 3.5 mm long, with a cross-section of 3 mm \times 3 mm. Both end faces of the crystal, which were perpendicular to the (111) crystallographic direction, were optically polished. The sample was placed at Brewster's angle between two folding concave mirrors M_1 and M_2 (radius of curvature, RoC = 100 mm). To remove the generated heat during laser operation, the crystal was wrapped with indium foil and tightly mounted in a water-cooled copper block.

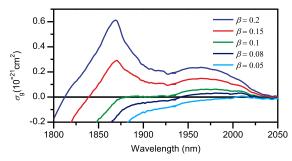


Fig. 1. Gain cross-section $\sigma_g = \beta \sigma_e - (1 - \beta) \ \sigma_a$ for the ${}^3F_4 \rightarrow {}^3H_6$ transition of Tm³⁺ in CNNGG for different inversion levels β .

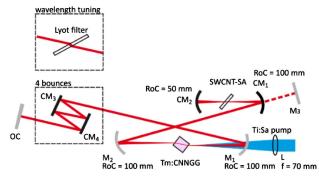


Fig. 2. Scheme of the mode-locked Tm:CNNGG laser (L, lens; M_1 and M_2 , dichroic folding mirrors; M_3 , plane mirror; CM_1 – CM_4 , chirped mirrors; OC, output coupler).

The cooling water was maintained at a temperature of 14.5° C. The pump source applied was a narrow-band CW Ti:sapphire laser with a maximum output power of 3.6 W at 785.8 nm. The output beam of the Ti:sapphire laser was focused into the Tm:CNNGG crystal with a spot radius of $\sim 30~\mu m$ using a focusing lens with a focal length of 70~mm. Four chirped mirrors (Layertec GmbH, Germany), CM_1 – CM_4 , were employed for intracavity dispersion compensation, each of them providing a group delay dispersion GDD = $-125~fs^2$ per bounce. The intracavity GDD was varied by the number of bounces on the two plane chirped mirrors ($CM_3~and~CM_4$).

Initially the CW laser performance was investigated with four different output couplers (OCs) having a transmission T of 0.2%, 0.5%, 1.5%, and 3.0%, employing mirror M_3 and without the chirped mirrors (see Fig. 2). The total cavity length was \sim 143 cm. Figure 3(a) shows the output power

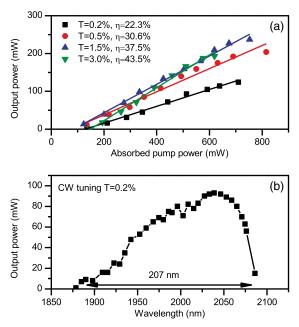


Fig. 3. (a) Output power versus absorbed pump power with different OCs in CW operation, and (b) wavelength tuning using a Lyot filter and OC with transmission T=0.2%. In (a), η denotes the slope efficiency with respect to the absorbed pump power.

versus the absorbed pump power. Maximum output powers of 124, 204, 237, and 193 mW were achieved with 0.2%, 0.5%, 1.5%, and 3.0% OCs, respectively. The maximum slope efficiency reached 43.5%. It should be noted that the output power increases slower than the linear dependence at the maximum applied pump levels, indicating an onset of thermal effects; thus, the full available pump power was not employed in this experiment.

Wavelength tuning of the Tm:CNNGG laser was studied by inserting a Lyot filter into the cavity, close to the output coupler (see Fig. 2). The Lyot filter was a 3.2 mm thick quartz plate with the optical axis at 60° to the surface. Figure 3(b) shows the tuning performance obtained with the T=0.2% OC: a wavelength tuning range of 207 nm (from 1879 to 2086 nm) was achieved, indicating the great potential of the Tm:CNNGG disordered crystal for the generation of ultrashort laser pulses.

To study the mode-locked laser performance, a transmission-type SWCNT-SA was inserted under Brewster's angle in the vicinity of the second waist formed by CM₁ and CM₂. This SA was characterized by a modulation depth of <0.50%, a saturation fluence of $<10 \mu J/cm^2$, and an interband carrier relaxation time τ_2 of ~1.16 ps around 2.0 μm [21]. Based on the ABCD propagation matrix method, the waist radius of the laser mode on the SWCNT-SA was calculated to be about 100 and 65 µm in the tangential and sagittal planes, respectively. With two bounces on each of the two plane chirped mirrors CM₃ and CM₄ (cf. Fig. 2), the total number of bounces on the chirped mirrors amounted to 11 per round trip. Taking into account the contribution of the 3.5 mm thick crystal under Brewster's angle and the 1 mm thick SA fused silica substrate, the total round trip GDD was calculated and details are shown in Fig. 4. Around 2000 nm, the total round trip GDD is about -1800 fs².

At first, an OC with T=1.5% was used. The pump power was set at 3.11 W, corresponding to an absorbed pump power of 1.15 W. After carefully optimizing the separations of M_1 - M_2 and CM_1 - CM_2 and adjusting the position of the SWCNT- CM_2 and Tm:CNNGG crystals, stable and self-starting modelocked operation was achieved. Pulses as short as 114 fs (FWHM intensity under sech² shape assumption) were obtained. Figure 5(a) shows the corresponding autocorrelation trace measured with a resolution of 8.4 fs. The corresponding optical spectrum shown in Fig. 5(b) was recorded with an

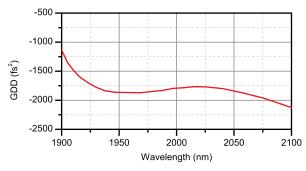


Fig. 4. Total round trip GDD introduced by chirped mirrors, Tm:CNNGG, and fused silica substrate of the SA.

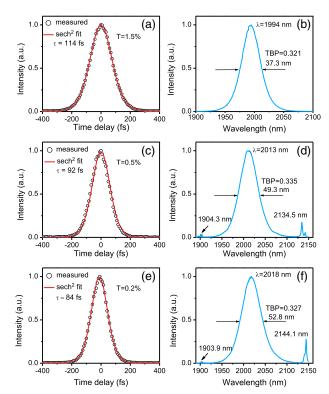
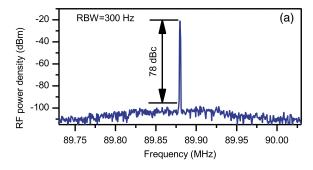


Fig. 5. Recorded autocorrelation traces (raw data and fits) and optical spectra with (a), (b) T=1.5%, (c), (d) T=0.5%, and (e), (f) T=0.2% (e-f).

optical spectrum analyzer (resolution: 0.3 nm); it has an FWHM of 37.3 nm centered at 1994 nm. This leads to a time-bandwidth product (TBP) of 0.321, corresponding to nearly Fourier-limited pulses. The measured average output power of the mode-locked laser was 98 mW. From the calculated intracavity power, an average fluence of $\sim 360~\mu J/cm^2$ on the SA was estimated.

With the aim of achieving shorter pulses, OCs with lower transmission, T = 0.5% and 0.2%, were further applied. Pulses as short as 92 and 84 fs, respectively, were obtained with these OCs [see Figs. 5(c) and 5(e)]. The spectral FWHM bandwidths of 49.3 nm centered at 2013 nm [Fig. 5(d)] and 52.8 nm centered at 2018 nm [Fig. 5(f)] resulted in TBPs of 0.335 and 0.327, respectively. These TBPs also indicate that the mode-locked pulses were nearly bandwidth limited at shorter durations. The corresponding average output powers were 40 and 22 mW, and the average fluence on the SA was \sim 450 and \sim 600 μ J/cm². No damage to the SWCNT-SA was observed at this fluence level. The applied pump level was the same for all OCs (corresponding to ~1.15 W of absorbed pump power) because the pulses tended to shorten with the pump level even in the region of slight thermal roll-off in Fig. 3(a), as long as the pump beam mode quality was excellent.

Compared to the previous results obtained with a Tm-doped CLNGG [11], much shorter pulses were achieved in the present experiment. The central wavelength in the mode-locked regime is similar to Ref. [11], which is an indication of similar overall cavity losses. Using a Ti:sapphire pump laser with high brightness and beam quality normally creates better



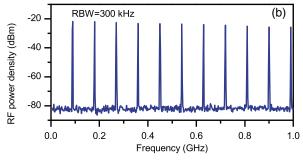


Fig. 6. Radio frequency spectra of the mode-locked Tm:CNNGG laser with the 1.5% OC: (a) fundamental beat note and (b) 1.0 GHz wide-span (RBW: resolution bandwidth).

conditions for Kerr lensing; however, pure Kerr-lens mode locking was not achievable in our laser without the SA, which obviously played a crucial role for the self-starting and self-stabilization of the mode-locked regime.

The stability of the mode-locked regime was also assessed by measurements with a radio-frequency (RF) spectrum analyzer. The RF spectrum shows an extinction ratio above the noise level of \sim 78 dB for the fundamental beat note, measured with a resolution bandwidth of 300 Hz [see Fig. 6(a)]. The uniform harmonic beat notes recorded in Fig. 6(b) indicate stable mode locking without any Q switching or other instabilities. The pulse repetition rate of \sim 89.9 MHz corresponds to a laser cavity length of \sim 1.66 m. The results were similar for the different OCs.

As the output coupler transmission T was decreased to 0.5% and 0.2%, stimulated Raman scattering sidebands were observed in the laser spectrum. They can be seen in the spectra shown in Figs. 5(d) and 5(f). The first Stokes component was clearly observed at 2134.5 nm (with T=0.5%) and 2144.1 nm (with T=0.2%), at similar separation from the corresponding fundamental central laser wavelengths. Simultaneously, the first anti-Stokes component was also observed (see the same figures), although its intensity was much lower. These sidebands can be well explained by the Raman spectrum of Tm:CNNGG [20]. The first Stokes and anti-Stokes components are generated in the CNNGG crystal by impulsive stimulated Raman scattering with the mode centered at \sim 278.4 cm⁻¹.

3. CONCLUSION

In conclusion, CW, tunable, and passively mode-locked laser operation were studied for the first time to our knowledge with

a disordered Tm:CNNGG garnet crystal. A maximum slope efficiency of 43.5% and wavelength tuning range of 207 nm were achieved in the CW regime. Introducing an SWCNT-SA, the Tm:CNNGG laser generated nearly Fourier-limited pulses as short as 84 fs with an output power of 22 mW at a repetition rate of \sim 90 MHz. The stable and self-starting mode-locked laser performance is evidenced by the remarkably high extinction ratio of 78 dB above the carrier of the fundamental beat note in the RF spectrum. To the best of our knowledge, these results represent the first sub-100 fs mode-locked laser operation achieved with a garnet crystal around 2 μ m, as well as the shortest pulses ever achieved with Tm-doped crystalline materials.

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