## **RESEARCH ARTICLE**

# Two-micron all-fiberized passively mode-locked fiber lasers with high-energy nanosecond pulse

Meng Wang<sup>1,2,3</sup>, Yijian Huang<sup>4</sup>, Zongpeng Song<sup>1,2,3</sup>, Jincheng Wei<sup>3</sup>, Jihong Pei<sup>1</sup>, and Shuangchen Ruan<sup>1,2,3</sup>

<sup>1</sup>Key Laboratory of Advanced Optical Precision Manufacturing Technology of Guangdong Higher Education Institutes, Shenzhen Technology University, Shenzhen 518118, China

<sup>2</sup>College of Electronic Information Engineering, Shenzhen University, Shenzhen 518060, China

<sup>3</sup>Shenzhen Key Laboratory of Laser Engineering, Shenzhen University, Shenzhen 518060, China

<sup>4</sup>Key Laboratory of Optoelectronic Devices and System of Ministry of Education, College of Optoelectronic Engineering, Shenzhen University, Shenzhen 518060, China

(Received 23 December 2019; revised 25 January 2020; accepted 18 February 2020)

#### Abstract

We report on mode-locked thulium-doped fiber lasers with high-energy nanosecond pulses, relying on the transmission in a semiconductor saturable absorber (SESA) and a carbon nanotube (CNTs-PVA) film separately. A section of an SMF–MMF–SMF structure multimode interferometer with a transmission peak wavelength of ~2003 nm was used as a wavelength selector to fix the laser wavelength. When the SESA acted as a saturable absorber (SA), the mode-locked fiber laser had a maximum output power of ~461 mW with a pulse energy of ~0.14  $\mu$ J and a pulse duration of ~9.14 ns. In a CNT-film-based mode-locked fiber laser, stable mode-locked pulses with the maximum output power of ~46 mW, pulse energy of ~26.8 nJ and pulse duration of ~9.3 ns were obtained. To the best of our knowledge, our experiments demonstrated the first 2  $\mu$ m region 'real' SA-based dissipative soliton resonance with the highest mode-locked pulse energy from a 'real' SA-based all-fiberized resonator.

Keywords: high pulse energy; mode-locking; nanosecond pulse; Tm-doped fiber laser

#### 1. Introduction

Mode-locked thulium-doped fiber lasers (TDFLs), which operate in the 2  $\mu$ m spectral region with pulse duration ranging from nanoseconds to femtoseconds, have attracted a lot of interest due to their applications in mid-infrared supercontinuum generation, remote sensing, laser processing, medicine and free-space communication<sup>[1–10]</sup>. In the 2  $\mu$ m regime, research efforts associated with mode-locked fiber lasers so far have mainly aimed at ultra-short pulse (~fs or ~ps pulse) generation. The nanosecond-scale mode-locked pulses have received less attention. In effect, nanosecond pulses with high pulse energy have played an important role in many potential applications, such as laser cutting and welding, laser ablation for surface cleaning and as a pump source for mid-infrared pulse laser generation<sup>[11]</sup>.

In general, *Q*-switching and gain switching are two wellknown techniques for nanosecond pulse generation. Usually, it is difficult for passively *Q*-switched fiber lasers to produce pulses with several nanoseconds duration owing to the limitation of cavity length and modulation depth of the saturable absorber  $(SA)^{[12]}$ . As for actively *Q*-switched fiber lasers, stable nanosecond pulses can be obtained. Nevertheless, a complex electronics controlling system is needed, which highly increases the complexity of the fiber laser. Gainswitched fiber laser is another alternative, but a high-energy pulse pump source is a challenge for 2  $\mu$ m gain-switched pulse generation<sup>[13,14]</sup>. Furthermore, compared with the mode-locked pulse, the *Q*-switched pulse and gain-switched pulse have poorer stability and cannot be reshaped since the phase of the pulse of *Q*-switching or gain switching is random.

In fact, passively mode-locked fiber lasers also can generate nanosecond pulses with high energy<sup>[15]</sup> in which dissipative soliton resonance (DSR) has been proved as an effective method. Chang *et al.* theoretically predicted that the DSR can be formed in the anomalous dispersion regime<sup>[16]</sup>. In the following year, the existence of

Correspondence to: S. Ruan, 3002 Lantian Road, Pingshan District, Shenzhen 518118, China. Email: scruan@sztu.edu.cn

<sup>©</sup> The Author(s) 2020. This is an Open Access article, distributed under the terms of the Creative Commons Attribution licence (http://creativecommons.org/ licenses/by/4.0/), which permits unrestricted re-use, distribution, and reproduction in any medium, provided the original work is properly cited.

high-energy nanosecond DSR pulses in long cavity passively mode-locked erbium-doped fiber lasers with anomalous dispersion was experimentally reported by Li et al.<sup>[17]</sup>. Relying on the nonlinear polarization rotation (NPR) technique, they obtained the output pulse energy of  $\sim$ 79.5 nJ at a pulse repetition rate of  $\sim$ 278 kHz, and the pulse duration could be tuned from  $\sim 12.8$  ns to  $\sim 155.4$  ns. After that, several works about DSR pulse generation in anomalous dispersion have been reported. Nevertheless, most of the works are focused on the 1.5  $\mu$ m region<sup>[18–21]</sup>. As for the 2 µm region, there are a few reports about the nanosecond passively mode-locked fiber laser. In 2013, Fu et al. realized a passively mode-locked Tm, Ho-codoped fiber laser with a pulse duration of  $\sim$ 122 ns and pulse energy of  $\sim$ 35.2 nJ by utilizing graphene as an SA and lengthening the cavity to  $\sim$ 213 m<sup>[22]</sup>. Based on NPR and  $\sim$ 460 m long cavity length, fundamental mode-locked pulses with a pulse duration of  $\sim$ 304 ns were obtained by Wang *et al.*<sup>[23]</sup>. By adding a section of  $\sim 40$  m long single mode fiber (SMF) in a simple linear cavity TDFL, Liu et al. obtained self-mode-locked pulses with a pulse duration of  $\sim 40$  ns and pulse energy of  $\sim 32.7 \text{ nJ}^{[24]}$ . Recently, relying on a graphene SA and a narrow bandwidth fiber Bragg grating (FBG), Wang et al. realized a nanosecond-scale mode-locked TDFL with  $\sim$ 54 m long cavity length, but the laser has a relatively low pulse energy of  $\sim 1.5 \text{ nJ}^{[25]}$ . In our previous work, by utilizing a nonlinear optical loop mirror (NOLM) and an F-P fiber interferometer with a  $\sim 1000$  m long cavity, stable pulses with pulse duration ranging from  $\sim$ 2.4 ns to  $\sim$ 21.2 ns and a maximum pulse energy of  $\sim 1.5 \,\mu J$  were obtained<sup>[26]</sup>. For DSR in the 2 µm region, Xu et al. reported a DSR TDFL with net normal dispersion<sup>[27]</sup>. Zhao et al. first reported a nanosecond DSR pulse in TDFL with anomalous dispersion based on nonlinear amplification loop mirror (NALM). With a figure-of-eight configuration resonator, the direct output pulse energy is  $\sim 40.5 \text{ nJ}^{[28]}$ . With a figure-9 cavity configuration, Kharitonov et al. demonstrated a DSR mode-locked TDFL with an average output power of 670 mW and pulse energy of 400 nJ<sup>[29]</sup>. The highest average output power of DSR from the TDFL oscillator is ~1.4 W with pulse energy of ~353 nJ, which was obtained by Du *et al.* with an NOLM-based mode-locked  $\sigma$ -shaped cavity TDFL<sup>[30]</sup>. Nevertheless, we noted that the highenergy DSR nanosecond pulses in the 2 µm region were almost based on the 'artificial' SAs, such as NOLM and NALM. There are a few experimental reports of 'real' SAbased DSR TDFL.

In this paper, we will present high pulse energy nanosecondscale passively mode-locked TDFLs with semiconductor saturable absorbers (SESAs) and carbon nanotubes (CNTs). A stable rectangular-shape DSR pulse with a maximum pulse energy up to ~0.14  $\mu$ J was first obtained in 'real' SA-based mode-locked TDFL. Moreover, our experimental results verified that the SMF-MMF-SMF (SMS, SMF: SMF-28, Corning; MMF: multimode fiber, AFS105/125Y) fiber structure multimode interferometer (MMI)<sup>[31, 32]</sup> can be used as an effective wavelength filter for mode-locked TDFL.

#### 2. Experimental setup and result

In our experiment, the SMS fiber structure MMI filter was constructed by fusion splicing two pieces of SMFs at both ends of a section of ~191 mm long MMF. The structure configuration and transmission characteristics of the fiber MMI are shown in Figures 1(a) and 1(b), respectively. The main transmission peak is at ~2003 nm with a transmission loss of ~0.5 dB, as shown in Figure 1(b).

### 2.1. SESA-based nanosecond mode-locked fiber laser

The experimental setup for an SESA-based nanosecond mode-locked TDFL is schematically shown in Figure 2. The gain fiber is a section of  $\sim$ 3 m length single mode



Figure 1. SMS fiber structure MMI: (a) structure configuration; (b) transmission characteristics<sup>[31]</sup>.



Figure 2. Schematic configuration of SESA-based nanosecond TDFL. LD, laser diode; PC, polarization controller; CPS, cladding power stripper; SMF, single mode fiber.

double-clad Tm-doped fiber (SM-TDF-10P/130-HE), which is pumped by a fiber-pigtailed 793 nm multimode laser diode (LD) with a maximum output power of  $\sim 12$  W via a  $(2 + 1) \times 1$  pump/signal combiner. The inner cladding diameter of the TDF is  $\sim$ 130 µm with absorption of  $\sim$ 3 dB/m at 793 nm. The unabsorbed pump light will effectively strip by a self-made cladding power stripper (CPS) behind the TDF. To ensure the unidirectional propagation of the laser in the cavity, a polarization independent optical isolator with insertion loss of  $\sim$ 2.6 dB was used in the cavity. After the isolator, an optical coupler with an output ratio of  $\sim$ 55% was used for laser outputting. The SESA and SMS fiber structure MMI were placed between two polarization controllers (PCs), which act as a mode locker and a wavelength selector, respectively. The SESA (SA-2000-25-10ps-x, BATOP) has the modulation depth, transmittance and saturation fluence of  $\sim$ 15%,  $\sim$ 74% and  $\sim$ 2 mJ/cm<sup>2</sup>, respectively, at  $\sim$ 2000 nm. The PCs are utilized to adjust the polarization state of the light and initiate the mode-locking. Moreover, a section of a  $\sim$ 50 m long SMF was added into the cavity to manage the dispersion and implement the nanosecond mode-locked pulse. Thus, a simple ring cavity laser was formed with a total cavity length and a dispersion of  $\sim 63.5$  m and -5.8 ps<sup>2</sup>, respectively.

In the beginning, to verify that the SMS fiber structure MMI can function as an available wavelength selector, the SESA was not inserted into the cavity. The output properties of the fiber laser with and without the SMS fiber structure MMI are shown in Figure 3. When the fiber MMI was inserted into the cavity, the lasing wavelength was  $\sim 2003$  nm and the output efficiency of the fiber laser only had a slight decrease (from  $\sim 24.7\%$  to  $\sim 22.6\%$ ), which means the fiber MMI can be used as an effective wavelength selector. Moreover, no pulses were observed by increasing the pump power and fully manipulating the PCs, which verifies the fiber MMI is not responsible for mode-locking<sup>[33]</sup>.

Then, the SESA was sandwiched between the FC/PC ferrules, and the threshold of the mode-locked operation was about  $\sim 2.06$  W. The typical mode-locked output properties at pump power of  $\sim 2.83$  W are exhibited in Figure 4. Stable mode-locked pulses are shown in Figures 4(a) and 4(b), which were measured by a photodetector (ET-5000,  $\sim 28$  ps



Figure 3. Output properties of the fiber laser without SA: (a) typical output spectra with and without the SMS fiber structure MMI; (b) output efficiencies of the fiber laser with and without the SMS fiber structure MMI.



Figure 4. Typical mode-locked properties of the fiber laser at pump power of  $\sim$ 2.83 W: (a), (b) mode-locked pulses; (c) output spectrum; (d) output RF spectra.

rise/fall time) together with a 20 GS/s high-speed oscilloscope with 1-GHz bandwidth (Tektronix, DPO 7104C digital phosphor oscilloscope). The mode-locked pulses have a pulse duration of  $\sim$ 2.4 ns with a rectangular shape, as shown in Figure 4(a). Since the response threshold of our detecting system is less than 1 ns, we confirm that the pulse durations are truly detected. In addition, we also measured the autocorrelation traces of the nanosecond pulses with a time span of 50 ps by a commercial autocorrelator (APE Pulsecheck). But no sub-pulse was observed, which further proves that the nanosecond mode-locked pulses do not have complicated ultra-short sub-pulses. Figure 4(c)shows the output spectrum of the mode-locked fiber laser. The central wavelength is  $\sim$ 2002.9 nm with 3 dB bandwidth of  $\sim 1.7$  nm, corresponding to the transmission spectrum of the SMS fiber structure MMI. Thus the time-band product is  $\sim$ 305, which indicates a serious chirp. With the spectrum analyzer (Agilent, E4407B), the radio-frequency (RF) spectra were measured and plotted in Figure 4(d). The repetition frequency of the mode-locked pulse is  $\sim$ 3.254 MHz with a signal-to-noise ratio of  $\sim$ 65 dB, which verifies that the laser operated at a stable fundamental frequency mode-locked state. The harmonic RF spectra with 300 MHz bandwidth are shown in the inset of Figure 4(d).

In the experiment, we noted that the mode-locked pulse duration was increased with the pump powers, and the details are shown in Figure 5. Figure 5(a) shows the evolution of the pulse envelopes. With the increasing pump power from  $\sim$ 2.06 W to  $\sim$ 4.40 W, the pulse duration was broadened from  $\sim$ 949 ps to  $\sim$ 9.2 ns, as given in Figure 5(b). Moreover, the average output powers of the mode-locked pulses were measured, and the corresponding pulse energies and peak powers were calculated, as depicted in Figure 6. The average output power increased linearly as the enhancement of the pump power with a slope efficiency of  $\sim 18.8\%$ . When the pump power reached  $\sim$ 4.40 W, the maximum output power of ~461 mW and pulse energy of ~0.14  $\mu$ J were obtained. The corresponding peak powers of the pulses were calculated and shown in Figure 6(b). We noted that the pulse peak powers were approximately maintained at a stable value of  $\sim 21$  W, when the pump power was in the range of  $\sim 2.57$  W to  $\sim 3.74$  W. Similar results were reported in Refs. [20, 28, 34], and the phenomenon was caused by the peak power clamping (PPC) effect. The decrease in the peak power at higher pump power was considered as the result of performance degrading of the SESA, since the high pulse energy will lead to a large heat accumulation at the SESA. The above characteristics of the



Figure 5. (a) Mode-locked pulse envelopes at different pump powers; (b) pulse duration versus pump power.



Figure 6. (a) Average output power and pulse energy variance with the pump power; (b) peak power variance with the pump power.

mode-locked pulses, such as pulse duration increase with the pump power, nanosecond rectangular-shaped pulse and PPC effect, indicated that the mode-locked fiber laser operated in the DSR regime<sup>[17–21, 28]</sup>.

# 2.2. CNTs-PVA film based nanosecond mode-locked fiber laser

Similarly to the fiber laser configuration in Figure 2, the experimental setup of CNTs-based nanosecond mode-locked TDFL is shown in Figure 7. The laser output ratio is 45%, and a piece of ~110 m long SMF was used for managing the dispersion and enhancing the nonlinearity. Thus, the total cavity length is ~120 m and the dispersion is ~10.75 ps<sup>2</sup>. Through a typical balanced twin-detector method, the saturable absorber property of the CNTs-PVA film was measured and shown in Figure 8. The ultrapulse light source is a mode-locked fiber laser with central

wavelength, pulse duration and repetition rate of  $\sim 1958$  nm,  $\sim 725$  fs and  $\sim 55$  MHz, respectively. The CNTs-PVA film SA has a modulation depth of  $\sim 23.7\%$  and nonsaturable loss of  $\sim 29.4\%$ .

When the CNTs-PVA film was inserted into the cavity, with proper manipulation of the PCs, the initial mode-locked operation also could be obtained. Here, we characterized the mode-locking performance at the maximum pump power of  $\sim 2.12$  W, as shown in Figure 9. The mode-locked fiber laser centered at  $\sim 2003.1$  nm with 3 dB bandwidth of  $\sim 1.3$  nm, which depends on the transmission spectrum of the SMS fiber structure MMI, is shown in Figure 9(a). Figure 9(b) depicts the RF spectra. The repetition rate of the mode-locked pulse is  $\sim 1.718$  MHz and the signal-to-noise ratio is larger than 63 dB, which indicate a stable fundamental mode-locked pulses. The single pulse duration of  $\sim 9.3$  ns with an approximate rectangular shape was measured. Together with the spectral width, the time-band product is

			Maximum output	Maximum pulse	Pulse	
	SA	Wavelength (nm)	power (mW)	energy (nJ)	duration (ns)	Reference
'Real' SA	Graphene	1897.7-1930.2	~33.9	35.2	~122-143	[22]
	Graphene	2010.15	$\sim 5.5$	1.4	~3.8–94.3	[25]
	CNTs	2003.1	$\sim 46$	26.8	~9.3	Our work
	SESA	2002.9	~461	141.7	~1.6–9.1	Our work
'Artificial' SA	NPR	1960	~95	/	~304	[23]
	Self-mode-locking	1985.5	$\sim\!66.8$	~32.7	$\sim 40$	[24]
	NALM	1975.56	~43.1	$\sim \!\! 40.5$	~3.74-72.19	[28]
	NALM	1940.2-1969.2	~60.73	~19.51	$\sim 0.48 - 6.19$	[27]
	NALM	1985	$\sim 670$	$\sim 400$	$\sim 50$	[29]
	NOLM	1948.13	$\sim 295$	$\sim 1.5 \ \mu J$	$\sim 2.4 - 21.2$	[30]
	NOLM	2005.9	$\sim 1.4 \text{ W}$	$\sim 353$	~1.9-13.7	[31]
	Intermode beating	1983	$\sim 1.03 \text{ W}$	$\sim \! 107$	45	[35]

Table 1. Comparison of output properties of 2 µm nanosecond mode-locked fiber lasers



Figure 7. Schematic configuration of CNTs-PVA-based nanosecond TDFL.



Figure 8. Saturable absorber properties of the CNTs-PVA.

~904. At a pump power of ~2.12 W, the average output power of the fiber laser is ~46 mW, corresponding to the pulse energy of ~26.8 nJ.

In addition, Table 1 summarizes the representative output performances of all-fiberized nanosecond mode-locked TDFLs. In our experiment, the maximum average output power is  $\sim$ 461 mW corresponding to the pulse energy of  $\sim$ 141.7 nJ, which is higher than the previously reported results of 'real' SA-based nanosecond passively mode-locked TDFLs and comparable to the results of 'artificial' SA-based ones. However, owing to the relatively low damage threshold of the 'real' SA, the 'real' SA-based nanosecond mode-locked TDFLs still have the lower output power and pulse energy.

### 3. Conclusion

In summary, we have experimentally presented long cavity nanosecond-scale mode-locked fiber lasers with wavelength above  $2 \,\mu$ m based on real SAs (SESA and CNTs) and SMS fiber structure MMI. The DSR pulse with output pulse



Figure 9. Typical mode-locking performance of the CNT-based fiber laser at pump power of  $\sim$ 2.12 W: (a) output spectrum; (b) output RF spectrum; (c), (d) mode-locked pulses.

energy up to ~0.14  $\mu$ J was reported for the first time in real SA-based TDFL. This nanosecond mode-locked pulse with high pulse energy can be widely used in industrial processing and as a pump source for mid-infrared (~3.5  $\mu$ m) pulse laser generation after further amplification. Moreover, our experimental results verified the feasibility of SMS fiber structure MMI functioning as an effective wavelength filter in mode-locked TDFL.

### Acknowledgement

This work was supported by the National Natural Science Foundation of China (NSFC) (No. 61575129).

#### References

- N. Leindecker, A. Marandi, R. L. Byer, K. L. Vodopyanov, J. Jiang, I. Hartl, M. Fermann, and P. G. Schunemann, Opt. Express 20, 7046 (2012).
- 2. J. Świderski, M. Michalska, and G. Maze, Opt. Express **21**, 7851 (2013).
- C. W. Rudy, A. Marandi, K. L. Vodopyanov, and R. L. Byer, Opt. Lett. 38, 2865 (2013).

- Z. Zheng, D. Ouyang, J. Zhao, M. Liu, S. Ruan, P. Yan, and J. Wang, Photon. Res. 4, 135 (2016).
- M. Tao, T. Yu, Z. Wang, H. Chen, Y. Shen, G. Feng, and X. Ye, Sci. Rep. 6, 23759 (2016).
- K. Yin, B. Zhang, L. Yang, and J. Hou, Photon. Res. 6, 123 (2018).
- G. Ehret, C. Kiemle, M. Wirth, A. Amediek, A. Fix, and S. Houweling, Appl. Phys. B 90, 593 (2008).
- P. Kadwani, J. Chia, F. Altal, R. A. Sims, C. Willis, L. Shah, D. Killinger, and M. C. Richardson, Proc. SPIE **7924**, 79240L (2011).
- 9. N. M. Fried and K. E. Murray, J. Endourol. 19, 25 (2005).
- C. W. Rudy, M. J. F. Digonnet, and R. L. Byer, Opt. Fiber Technol. 20, 642 (2014).
- S. D. Chowdhury, A. Pal, D. Pal, S. Chatterjee, M. C. Paul, R. J. Sen, and M. Pal, Opt. Lett. 42, 2471 (2017).
- 12. T. Hakulinen and O. G. Okhotnikov, Opt. Lett. **32**, 2677 (2007).
- 13. M. Jiang and P. Tayebati, Opt. Lett. 32, 1797 (2007).
- B. C. Dickinson, S. D. Jackson, and T. A. King, Opt. Commun. 182, 199 (2000).
- H. Wang, Y. Wang, W. Zhao, W. Zhang, T. Zhang, X. Hu, Z. Yang, H. Liu, K. Duan, X. Liu, C. Li, D. Shen, Z. Sui, and B. Liu, Opt. Express 18, 7263 (2010).
- W. Chang, J. M. Soto-Crespo, A. Ankiewicz, and N. Akhmediev, Phys. Rev. A 79, 033840 (2009).
- X. Li, X. Liu, X. Hu, L. Wang, H. Lu, Y. Wang, and W. Zhao, Opt. Lett. 35, 3249 (2010).

- L. Duan, X. Liu, D. Mao, L. Wang, and G. Wang, Opt. Express 20, 265 (2012).
- Z. C. Luo, W. J. Cao, Z. B. Lin, Z. R. Cai, A. P. Luo, and W. C. Xu, Opt. Lett. 37, 4777 (2012).
- 20. K. Krzempek, Opt. Express 23, 30651 (2015).
- 21. J. Lee, J. Koo, and J. H. Lee, Opt. Eng. 55, 081309 (2016).
- B. Fu, L. Gui, X. Li, X. Xiao, H. Zhu, and C. Yang, IEEE Photonics Technol. Lett. 25, 1447 (2013).
- 23. X. Wang, P. Zhou, X. Wang, H. Xiao, and Z. Liu, Opt. Express 22, 6147 (2014).
- 24. C. Liu, Z. Luo, Y. Huang, B. Qu, H. Cheng, Y. Wang, D. Wu, H. Xu, and Z. Cai, Appl. Opt. 53, 892 (2014).
- 25. X. Wang, J. Zhang, Z. Gao, G. Xia, and Z. Wu, Acta Phys. Sin. 66, 114209 (2017).
- 26. M. Wang, Y. J. Huang, J. W. Yang, Y. Zhang, and S. C. Ruan, Laser Phys. Lett. **15**, 085110 (2018).
- 27. Y. Xu, Y. Song, G. Du, P. Yan, C. Guo, G. Zheng, and S. Ruan, IEEE Photonics J. 7, 1502007 (2015).

- J. Zhao, D. Ouyang, Z. Zheng, M. Liu, X. Ren, C. Li, S. Ruan, and W. Xie, Opt. Express 24, 12072 (2016).
- 29. S. Kharitonov and C. Brès, in 2017 European Conference on Lasers and Electro-Optics and European Quantum Electronics Conference (Optical Society of America, 2017), paper CJ-13-4.
- T. Du, W. Li, Q. Ruan, K. Wang, N. Chen, and Z. Luo, Appl. Phys. Express 11, 052701 (2018).
- M. Wang, J. Zhao, Y. Huang, S. Lin, and S. Ruan, IEEE Photonics J. 9, 1506408 (2017).
- 32. L. Zhang, J. Hu, J. Wang, and Y. Feng, Opt. Lett. **37**, 3828 (2012).
- T. Chen, Q. Zhang, Y. Zhang, X. Li, H. Zhang, and W. Xia, Photon. Res. 6, 1033 (2018).
- 34. L. Mei, G. Chen, L. Xu, X. Zhang, C. Gu, B. Sun, and A. Wang, Opt. Lett. 39, 3235 (2014).
- J. Zhang, D. Wu, R. Zhao, R. Wang, and S. Dai, High Power Laser Sci. Eng. 7, e65 (2019).