

# Effect of SiO<sub>2</sub> on the thermal stability and spectroscopic properties of Er<sup>3+</sup>-doped tellurite glasses\*

ZHENG Shi-chao (郑世超)\*\* and ZHOU Ya-xun (周亚训)

College of Information Science and Engineering, Ningbo University, Ningbo 315211, China

(Received 9 February 2014)

©Tianjin University of Technology and Springer-Verlag Berlin Heidelberg 2014

Er<sup>3+</sup>-doped tellurite glass (TeO<sub>2</sub>-ZnO-Na<sub>2</sub>O) prepared using the conventional melt-quenching method is modified by introducing the SiO<sub>2</sub>, and its effects on the thermal stability of glass host and the 1.53 μm band spectroscopic properties of Er<sup>3+</sup> are investigated by measuring the absorption spectra, 1.53 μm band fluorescence spectra, Raman spectra and differential scanning calorimeter (DSC) curves. It is found that for Er<sup>3+</sup>-doped tellurite glass, besides improving its thermal stability, introducing SiO<sub>2</sub> is helpful for the further improvement of the fluorescence full width at half maximum (FWHM) and bandwidth quality factor. The results indicate that the prepared Er<sup>3+</sup>-doped tellurite glass containing an appropriate amount of SiO<sub>2</sub> has good prospect as a candidate of gain medium applied for 1.53 μm broadband amplifier.

**Document code:** A **Article ID:** 1673-1905(2014)03-0209-4

**DOI** 10.1007/s11801-014-4016-8

With the increasing need for telecom bandwidth, the development of novel optical amplifiers offering wider spectrum and higher gain than current technologies is an urgent research topic. Most current amplifying technologies rely on Er<sup>3+</sup>-doped fiber amplifier (EDFA) made from silica glass fiber. However, the typical gain bandwidth is only about 35 nm in the 1.53 μm band, which cannot meet the future needs for broadband transmission<sup>[1]</sup>. Er<sup>3+</sup>-doped tellurite glass exhibits a large stimulated emission cross-section and a broad fluorescence full width at half maximum (FWHM) at 1.53 μm band<sup>[2-4]</sup>. In addition, tellurite glass also possesses high linear and nonlinear refractive indices, high rare-earth solubility and high gain coefficient. However, its poor glass thermal stability makes it difficult to be used in practice.

It is well known that SiO<sub>2</sub> is a good glass former and adding SiO<sub>2</sub> into glass can enhance the thermal stability and improve the resistance for crystallization<sup>[5]</sup>. However, samples containing TeO<sub>2</sub> and SiO<sub>2</sub> have been investigated very little so far. In this paper, SiO<sub>2</sub> is introduced into the Er<sup>3+</sup>-doped TeO<sub>2</sub>-ZnO-Na<sub>2</sub>O glasses. Especially, the effects of SiO<sub>2</sub> on the thermal stability, Judd-Ofeldt parameters, FWHM and stimulated emission cross-section of Er<sup>3+</sup>-doped tellurite glass are analyzed and discussed.

The Er<sup>3+</sup>-doped tellurite glasses with (75-x)TeO<sub>2</sub>-15ZnO-9.5Na<sub>2</sub>O-xSiO<sub>2</sub>-0.5Er<sub>2</sub>O<sub>3</sub> (x=0, 3%, 6%, 9% in

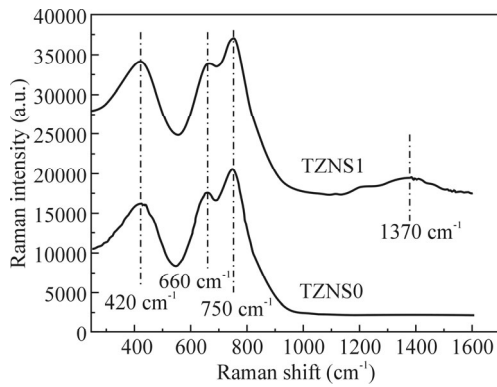
mole ratio) were prepared in alumina crucibles at 1000-1100 °C for about 30 min. The obtained glass melts were annealed at about 330 °C for 120 min, and next were cut and well polished for optical measurements. The prepared samples were denoted as TZNS0-3 for short, in which component Na<sub>2</sub>O was introduced in Na<sub>2</sub>CO<sub>3</sub> form.

The absorption spectra were measured with a Lambda 950 UV/VIS/NIR spectrophotometer in the wavelength range of 400-1700 nm. The fluorescence spectra were measured with a Jobin Yvon Triax550 upon excitation of a 975 nm laser diode (LD). The Raman spectra were measured by a Renishaw Micro-Raman instrument and the thermal stability of glass sample was determined using a differential scanning calorimeter (DSC) of TA Instrument Q2000 at a heating rate of 10 K/min.

Fig.1 displays the measured Raman spectra of TZNS0 and TZNS1 glass samples with and without SiO<sub>2</sub>. The Raman scattering band centered around 420 cm<sup>-1</sup> is assigned to the stretching vibrations of Te-O-Te or O-Te-O linkages between two [TeO<sub>4</sub>] four-coordinate atoms, and the bands centered around 660 cm<sup>-1</sup> and 750 cm<sup>-1</sup> originate from the stretching vibrations of Te-O bond in the continuous networks composed of [TeO<sub>4</sub>] trigonal bipyramid and [TeO<sub>3+6</sub>] trigonal pyramid structural units, respectively. A new Raman scattering peak centered around 1370 cm<sup>-1</sup> is observed in the SiO<sub>2</sub> contained glass sample, which is attributed to the characteristic stretching vibration modes of Si-O or Si=O bonds<sup>[6]</sup>.

\* This work has been supported by the National Natural Science Foundation of China (No.61178063), the Graduate Innovative Scientific Research Project of Zhejiang Province (No.YK2010048), the Scientific Research Foundation of Graduate School of Ningbo University (No.G13035), and the K. C. Wong Magna Fund and Hu Lan Outstanding Doctoral Fund in Ningbo University.

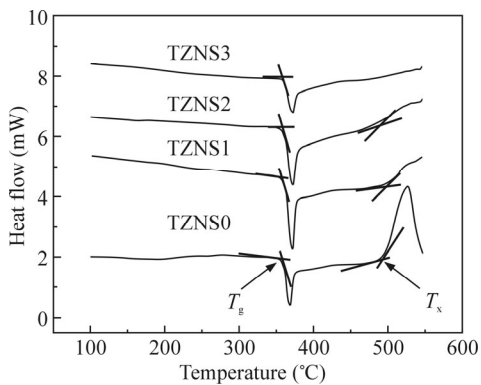
\*\* E-mail:526522025@qq.com



**Fig.1 Raman spectra of TZNS0 and TZNS1 glass samples**

Fig.2 displays the measured DSC curves of TZNS0-3 glass samples, and the obtained three characteristic temperatures, i.e., glass transition temperature ( $T_g$ ), crystallization onset temperature ( $T_x$ ) and the difference ( $\Delta T=T_x-T_g$ ) between them, are recorded in Tab.1.

In general, more than 100 °C for  $\Delta T$  can satisfy the requirement of conventional fiber drawing process. It can be seen that on the whole, the values of  $T_g$  and  $\Delta T$  for the prepared TZNS1-3 glass samples increase with the substitution of SiO<sub>2</sub> for TeO<sub>2</sub> component. Especially, no obvious crystallization onset peak is found in the glass sample TZNS3. The results indicate that the current investigated TZNS1-3 glasses are more desirable for fiber drawing in view of the thermal stability.



**Fig.2 DSC curves of TZNS0-3 glass samples**

Fig.3 reveals the measured absorption spectra of the prepared TZNS0-3 glass samples in the wavelength range of 400-1700 nm. All the inhomogeneously broadened absorption bands are attributed to 4f-4f transition of Er<sup>3+</sup> from the <sup>4</sup>I<sub>15/2</sub> ground state to the different excited states, and their baycenters are found to be almost identical except for some differences in the band intensities. It suggests that the introduction of SiO<sub>2</sub> has no obvious effect on the position of the lowest Stark sub-levels of Er<sup>3+</sup> manifolds.

The Judd-Ofelt theory<sup>[7,8]</sup> is a useful tool for analyzing

the spectroscopic properties of rare-earth doped glass or crystal, and is often used to evaluate the spectroscopic parameters, such as the three intensity parameters  $\Omega_t$  ( $t=2, 4, 6$ ). The parameters can be acquired from the measured absorption spectrum and the refractive index of glass host by a least-square fitting approach according to the procedure provided in Ref.[9]. The obtained three intensity parameters of the prepared glass samples together with the other glasses are listed in Tab.2.

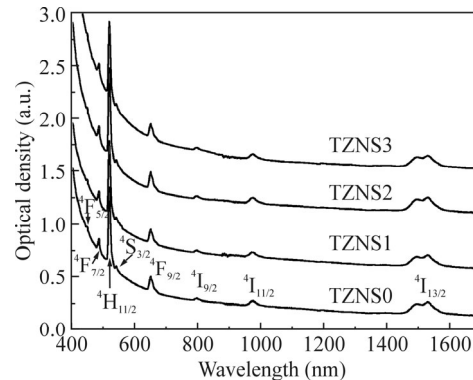
**Tab.1 Glass transition temperature ( $T_g$ ), crystallization temperature ( $T_x$ ), and the difference  $\Delta T$  between them for prepared TZNS0-3 glasses**

Sample	$T_g$ (°C)	$T_x$ (°C)	$\Delta T$ (°C)
TZNS0	357	470	113
TZNS1	357	498	141
TZNS2	358	491	133
TZNS3	359	-	$\infty$

According to the Tanabe studies<sup>[10]</sup>, the fluorescence emission, produced when Er<sup>3+</sup> transits from the various excited states to the ground state <sup>4</sup>I<sub>15/2</sub>, mainly comes from the contributions of electric-dipole transition. In order to achieve the broad and flat fluorescence emission spectrum, increasing the spectral line strength of electric-dipole transition is an effective way. Here, the spectral line strength of electric-dipole transition can be obtained based on the following linear combination of three intensity parameters<sup>[11]</sup>

$$S_{ed}[^4I_{13/2} \rightarrow ^4I_{15/2}] = 0.0188\Omega_2 + 0.1176\Omega_4 + 1.4617\Omega_6. \quad (1)$$

From the above equation, it is easily concluded that the spectral line strength  $S_{ed}$  mainly depends on the parameter  $\Omega_6$ , as it has a largest weighted coefficient of 1.4617. Apparently, increasing the parameter  $\Omega_6$  can effectively improve the value of  $S_{ed}$ . The calculated intensity parameter  $\Omega_6$  in this paper has a larger value compared with other glasses, such as silicate, aluminate and germanate<sup>[12]</sup>, indicating that Er<sup>3+</sup> ions in our prepared tellurite glass are more possible to realize broad emission.

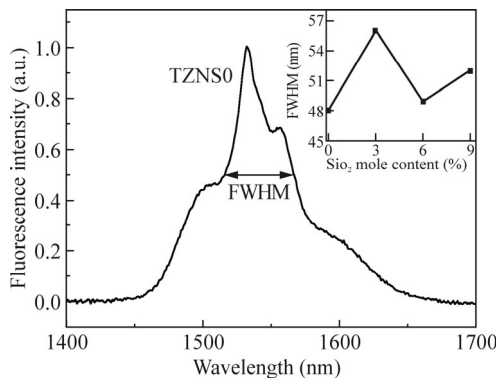


**Fig.3 Absorption spectra of TZNS0-3 glass samples**

**Tab.2 The compositional dependence of  $\Omega_t$  ( $t=2, 4, 6$ ) of  $\text{Er}^{3+}$  in TZNS0—3 samples**

Composition	$\Omega_2$ ( $10^{-20}$ cm)	$\Omega_4$ ( $10^{-20}$ cm)	$\Omega_6$ ( $10^{-20}$ cm)
TZNS0	5.39	1.41	0.93
TZNS1	5.82	1.66	0.94
TZNS2	5.87	1.71	0.91
TZNS3	5.75	1.54	1.00
Silicate <sup>[13]</sup>	4.23	1.04	0.61
Aluminate <sup>[13]</sup>	5.60	1.60	0.61
Germanate <sup>[13]</sup>	5.81	0.85	0.28

Fig.4 displays the measured 1.53  $\mu\text{m}$  band fluorescence spectrum corresponding to  $\text{Er}^{3+}:^4\text{I}_{13/2} \rightarrow ^4\text{I}_{15/2}$  transition of TZNS0 glass sample in the wavelength range of 1400-1700 nm. Because of the differences of the fluorescence spectra in different glass hosts, FWHM is often used as a semi-quantitative indication of the bandwidth. From the inset of Fig.4, it can be seen that the FWHM in the glass samples containing  $\text{SiO}_2$  has a large value.



**Fig.4 Fluorescence spectrum of TZNS0 glass sample (The inset is FWHM.)**

Amplification characteristics of the EDFA gain medium are described mainly by the gain quality factor and bandwidth quality factor. The bandwidth quality factor is generally measured by the product of peak stimulated emission cross-section  $\sigma_e$  and FWHM<sup>[13,14]</sup>. The larger the product, the better the amplification bandwidth characteristics of the gain medium<sup>[15,16]</sup>. Tab.3 lists the relevant spectral parameters and bandwidth quality factors of  $\text{Er}^{3+}$  ions in the TZNS0-3 glasses. The FWHM is obtained directly from the measured fluorescence emission spectra, and the stimulated emission cross-section is obtained from the absorption cross-section  $\sigma_a$  of  $\text{Er}^{3+}:^4\text{I}_{15/2} \rightarrow ^4\text{I}_{13/2}$  transition according to the McCumber theory<sup>[17]</sup>:

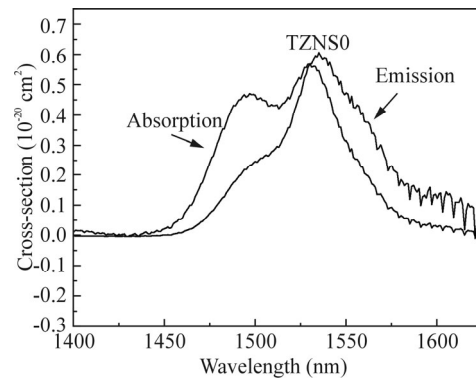
$$\sigma_e(\lambda) = \sigma_a(\lambda) \exp[(\epsilon - h\nu)/kT], \quad (2)$$

where  $h$  is the Planck constant,  $k$  is the Boltzmann constant,  $\epsilon$  is the net free energy required to excite one  $\text{Er}^{3+}$  from  $^4\text{I}_{15/2}$  to  $^4\text{I}_{13/2}$  level at temperature  $T$  and can be calculated using the method provided in Ref.[18].

Absorption cross-section  $\sigma_a$  is determined from the absorption spectra shown in Fig.3, which is related by:

$$\sigma_a(\lambda) = \frac{2.303}{NL} OD(\lambda), \quad (3)$$

where  $OD(\lambda)$  is the optical density,  $N$  is the  $\text{Er}^{3+}$  ion doping concentration, and  $L$  is the thickness of glass sample. From Tab.3, it can be observed that the bandwidth quality factor in the glass samples containing  $\text{SiO}_2$  has a larger value compared with that of glass sample without  $\text{SiO}_2$ . In addition, the bandwidth quality factor of glass samples developed in this paper is larger than that of silicate glass ( $22.0 \times 10^{-20} \text{ cm}^2 \text{ nm}$ )<sup>[19]</sup>, germanate glass ( $23.9 \times 10^{-20} \text{ cm}^2 \text{ nm}$ )<sup>[20]</sup> and phosphate glass ( $23.7 \times 10^{-20} \text{ cm}^2 \text{ nm}$ )<sup>[21]</sup>, which is attributed to the large stimulated emission cross-section and FWHM. Especially, the TZNS1 glass sample can be used as a potential host material for developing broadband EDFA in WDM system.



**Fig.5 Stimulated absorption and emission cross-sections of TZNS0 glass sample**

**Tab.3 The stimulated absorption ( $\sigma_a$ ) and emission ( $\sigma_e$ ) cross-sections, FWHMs, and bandwidth quality factors of  $\text{Er}^{3+}$  in TZNS0—3 glass samples**

Sample	$\sigma_a$ ( $10^{-21} \text{ cm}^2$ )	$\sigma_e$ ( $10^{-21} \text{ cm}^2$ )	FWHM (nm)	$\sigma_e \times \text{FWHM}$ ( $10^{-20} \text{ cm}^2 \text{ nm}$ )
TZNS0	5.71	6.08	48	29.18
TZNS1	6.02	6.63	56	37.13
TZNS2	6.04	6.66	49	32.63
TZNS3	6.27	6.58	52	34.22

In this paper, the effects of  $\text{SiO}_2$  on the spectroscopic properties and thermal stability of  $\text{Er}^{3+}$ -doped tellurite glasses are investigated. It is found that the introduction of an appropriate amount of  $\text{SiO}_2$  can enhance the thermal stability of glass hosts. Meanwhile, the introduction of  $\text{SiO}_2$  can also increase the FWHM and bandwidth quality factor of  $\text{Er}^{3+}$ . Therefore, as a glass host applied for the 1.53  $\mu\text{m}$  broadband amplification, an appropriate amount of  $\text{SiO}_2$  can be selectively introduced into the  $\text{Er}^{3+}$ -doped tellurite glasses.

## References

- [1] R. El-Mallawany, A. Patra, C. S. Friend, R. Kapoor and P. N. Prasad, *Opt. Mater.* **26**, 267 (2004).
- [2] S. Wang, Y. X. Zhou, S. X. Dai, X. S. Wang, X. Shen, Y. Wu and X. C. Xu, *Journal of Optoelectronics·Laser* **22**, 12 (2011). (in Chinese)
- [3] T. Sasikala, M. L. Rama, K. Pavani and T. Chengaiah, *J. Alloy Compd.* **542**, 271 (2012).
- [4] Y. X. Zhou, X. C. Xu, F. Chen, J. H. Lin and G. B. Yang, *Optoelectronics Letters* **8**, 273 (2012).
- [5] T. F. Xu, X. D. Zhang, S. X. Dai, Q. H. Nie, X. Shen and X. H. Zhang, *Physica B* **389**, 242 (2007).
- [6] M. Okuno, B. Reynard, Y. Shimada, Y. Syono and C. Willaime, *Phys. Chem. Minerals* **26**, 304 (1999).
- [7] B. R. Judd, *Phys. Rev.* **127**, 750 (1962).
- [8] G. S. Ofelt, *J. Chem. Phys.* **37**, 511 (1962).
- [9] S. Tanabe, T. Ohyagi, N. Soga and T. Hanada, *Phys. Rev. B* **46**, 3305 (1992).
- [10] S. Tanabe, *J. Non-Cryst. Solids* **259**, 1 (1999).
- [11] M. J. Weber, *Phys. Rev.* **157**, 262 (1967).
- [12] J. Qiu, Y. Shimizugawa, Y. Iwabuchi and K. Hirao, *Appl. Phys. Lett.* **71**, 43 (1997).
- [13] S. C. Zheng, Y. W. Qi, S. X. Peng, D. D. Yin, Y. X. Zhou and S. X. Dai, *Optoelectronics Letters* **9**, 0461 (2013).
- [14] J. S. Wang, E. M. Vogel and E. Snitzer, *Opt. Mater.* **3**, 187 (1994).
- [15] M. Naftaly, S. Shen and A. Jha, *Appl. Opt.* **39**, 4979 (2000).
- [16] S. X. Dai, T. F. Xu, Q. H. Nie, X. Shen, J. J. Zhang and L. L. Hu, *Acta Phys. Sin.* **55**, 1479 (2006).
- [17] D. E. McCumber, *Phys. Rev.* **A299**, 134 (1964).
- [18] W. J. Miniscalco and R. S. Quimby, *Opt. Lett.* **16**, 258 (1991).
- [19] B. C. Hwang, S. Jiang, T. Luo, K. Seneschal, G. Sorbello, M. Morrell, F. Smektala, S. Honkanen, J. Lucas and N. Peyghambarian, *IEEE Photon. Technol. Lett.* **13**, 657 (2001).
- [20] H. Lin, E. Y. B. Pun, S. Q. Man and X. R. Liu, *J. Opt. Soc. Am. B* **18**, 602 (2001).
- [21] G. C. Righini, S. Pelli, M. Fossi, M. Brenci, A. A. Lipovskii, E. V. Kolobkova, A. Speghini and M. Bettinelli, *Proc. SPIE* **210**, 4282 (2001).