

# Room temperature ferromagnetism of nonmagnetic element Ca-doped LiNbO<sub>3</sub> films\*

LIU Hua-rui (刘华蕊), LI Min (李敏), SUN Peng (孙鹏), WANG Shi-qi (王诗琪), JIN Xin (金鑫), SUN Xian-ke (孙现科), AN Yu-kai (安玉凯)\*\*, and LIU Ji-wen (刘技文)\*\*

Tianjin Key Laboratory for Photoelectric Materials and Devices, School of Material Science and Engineering, Tianjin University of Technology, Tianjin 300384, China

(Received 5 December 2013)

©Tianjin University of Technology and Springer-Verlag Berlin Heidelberg 2014

The nonmagnetic element Ca-doped LiNbO<sub>3</sub> films were prepared on Si (111) substrates by radio frequency (RF) magnetron sputtering technique. X-ray diffraction (XRD) and X-ray photoelectron spectroscopy (XPS) indicate that the Ca atoms enter the LiNbO<sub>3</sub> lattice in the form of Ca<sup>2+</sup> ions. Superconducting quantum interference device (SQUID) results show that the Ca-doped LiNbO<sub>3</sub> films have room-temperature ferromagnetism and a maximum saturation magnetization of 4800 A/m at the 3% of Ca atom doping concentration. The room temperature ferromagnetism of the Ca-doped LiNbO<sub>3</sub> films can be attributed to the occurrence of vacancies due to Ca doping and is the intrinsic property.

**Document code:** A **Article ID:** 1673-1905(2014)02-0115-4

**DOI** 10.1007/s11801-014-3225-5

In recent years, the multiferroic materials with simultaneous ferromagnetism and ferroelectricity have attracted much attention because of their interesting physical properties and a wide range of potential applications in some devices, such as magnetic data storage, non-volatile memories, and sensors<sup>[1-3]</sup>. However, the single-phase multiferroic materials are relatively rare<sup>[4]</sup>, due to two reasons: the ferromagnetism and ferroelectricity turn out to be mutually exclusive<sup>[5-7]</sup> and ferromagnetism tends to be metallic while ferroelectricity coexists with insulating state<sup>[8]</sup>. Therefore, in order to resolve these problems, many investigations are being extensively carried out.

Some researches suggested that the ABO<sub>3</sub>-type ferroelectric oxide can be easily doped by the 3d transition metal to produce ferromagnetism. For this reason, as a representative of the ABO<sub>3</sub>-type ferroelectric oxide, LiNbO<sub>3</sub> is worth investigating for its unexceptionable ferroelectric and crystal structural properties, such as a strong room temperature spontaneous polarization of 70 μC/cm<sup>2</sup> and a high ferroelectric transition temperature of 1483 K<sup>[9,10]</sup>. Recently, some investigations showed that the ferromagnetism of LiNbO<sub>3</sub> can be obtained by the introduction of transition metal (TM) ions into LiNbO<sub>3</sub>. Song et al<sup>[11,12]</sup> reported that the implantation of cobalt ions into the LiNbO<sub>3</sub> wafers induces the ferromagnetic behavior of LiNbO<sub>3</sub> at room temperature. Chen et al<sup>[13]</sup> also reported that the LiNbO<sub>3</sub> wafers implanted with manganese ions exhibit ferromagnetism. Lately, Sheng et

al<sup>[14]</sup> reported that vanadium can induce the room temperature ferromagnetic behavior of LiNbO<sub>3</sub> by doping it into LiNbO<sub>3</sub>. These researches are beneficial to the exploration of new ABO<sub>3</sub>-type multiferroic materials.

However, the TM-doped LiNbO<sub>3</sub> easily produces the TM clusters and magnetic secondary phases, which obscures the intrinsic magnetic property of the TM-doped LiNbO<sub>3</sub>. But the nonmagnetic element doped LiNbO<sub>3</sub> can effectively avoid the above disadvantages if the nonmagnetic element doped LiNbO<sub>3</sub> has magnetic property, because unlike TM elements, the nonmagnetic element itself and related compounds are not magnetic. Therefore, the nonmagnetic element doped LiNbO<sub>3</sub> is an ideal system to investigate the magnetic mechanism of the doped LiNbO<sub>3</sub>. As far as we know, there is no report on the room temperature ferromagnetism of nonmagnetic element doped LiNbO<sub>3</sub>.

In this work, we choose Ca as doping element of LiNbO<sub>3</sub> films to investigate the magnetic properties of LiNbO<sub>3</sub> and the results exhibit that the Ca-doped LiNbO<sub>3</sub> films present room temperature ferromagnetism. Because the Ca clusters and its secondary phases are not magnetic, the magnetic property of the Ca-doped LiNbO<sub>3</sub> films should be an intrinsic property. The origin of the ferromagnetic behavior is also discussed in this paper.

The Ca-doped LiNbO<sub>3</sub> films were deposited on Si (111) substrates with a SiO<sub>2</sub> buffer layer by radio frequency (RF) magnetron sputtering technique. The base pressure

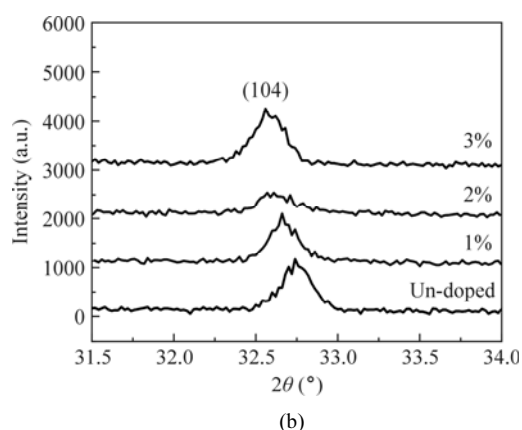
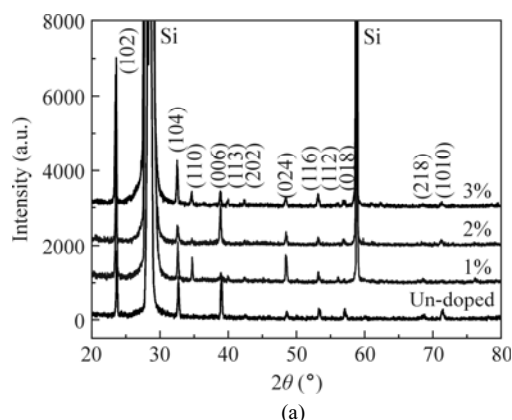
\* This work has been supported by the National Natural Science Foundation of China (Nos.11174217 and 10904110).

\*\* E-mails: jwliu@tjut.edu.cn; ayk\_bj@126.com

of the sputtering system was about  $9 \times 10^{-5}$  Pa and the films were deposited at the sputtering pressure of 1.6 Pa for 2 h. A mixture of  $O_2/Ar$  with a flow ratio of 5:10 was introduced into the deposition chamber as the sputtering gas. The deposition temperature and RF power were controlled at 575 °C and 85 W, respectively. The  $LiNbO_3$  target on top of some Ca chips (square in  $3 \text{ mm} \times 3 \text{ mm}$ , 99.99% purity) symmetrically attached as the co-sputtering target was used to fabricate the Ca-doped  $LiNbO_3$  films. By adjusting the number of Ca chips, the  $LiNbO_3$  films with different Ca-doped concentrations were obtained. In order to improve the crystalline quality of the films, the deposited-films were annealed by a tube furnace at 1000 °C for 1 h in air.

X-ray diffraction (XRD) analysis with  $Cu \text{ K}\alpha$  radiation ( $\lambda=0.15406 \text{ nm}$ ) was used to determine the crystal structure and phases in the films. The thickness of the films was estimated using a scanning electron microscopy (SEM). The composition and the element valence of the films were characterized by X-ray photoelectron spectroscopy (XPS) spectra. The 1%, 2% and 3% Ca atom doping concentrations in  $LiNbO_3$  films were determined by the elemental sensitivity factor method based on the XPS measurement. The ferromagnetism of the films was recorded using a superconducting quantum interference device (SQUID) magnetometer at 300 K.

Fig.1(a) shows the XRD patterns of the undoped, 1%, 2% and 3% Ca-doped  $LiNbO_3$  films annealed at 1000 °C. From Fig.1(a), it can be seen that the diffraction peaks of the films are indexed as a typical R3c structure of  $LiNbO_3$ <sup>[15]</sup> and no diffraction peaks connect with the metallic calcium or its secondary phases, suggesting that the  $LiNbO_3$  film is formed and there are not the calcium clusters or secondary phases in the Ca-doped  $LiNbO_3$  films. Fig.1(b) shows the enlarged view of (104) diffraction peaks of the Ca-doped  $LiNbO_3$  films. The clear shift of the (104) peak to lower  $2\theta$  angles is observed with the increase of Ca doping concentration, because the Ca doping results in the change of lattice parameter of  $LiNbO_3$ , indicating that the Ca atoms are incorporated into the lattice of  $LiNbO_3$ .

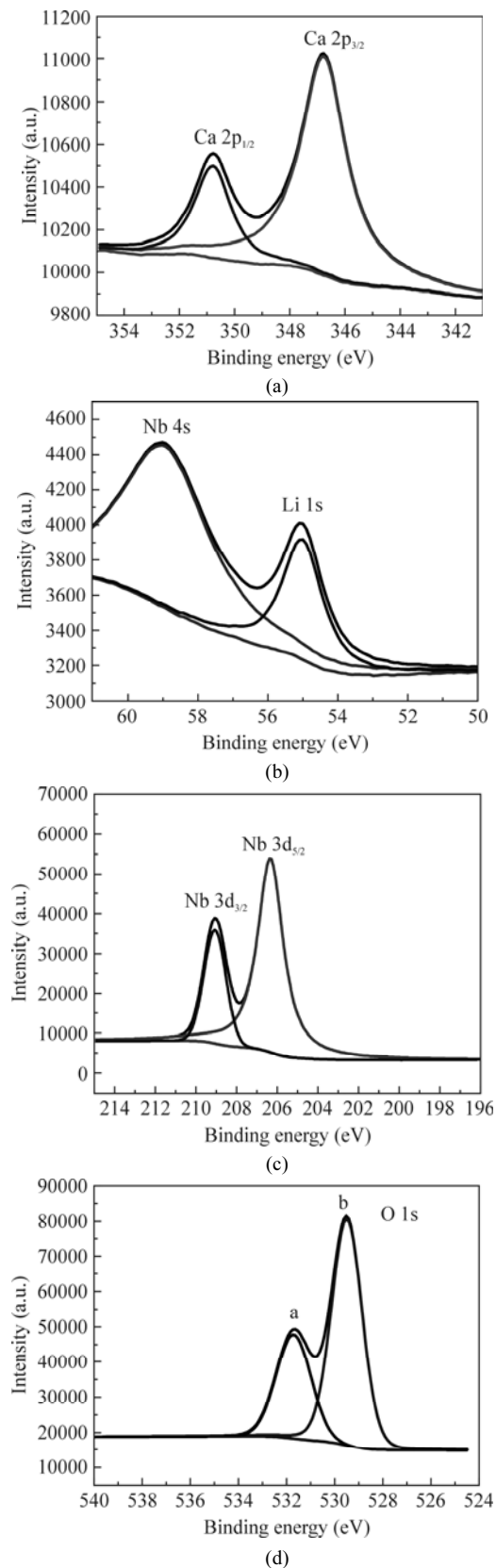


**Fig.1 (a) X-ray diffraction patterns of the un-doped, 1%, 2% and 3% Ca-doped  $LiNbO_3$  films; (b) Enlarged view of (104) diffraction peaks of the Ca-doped  $LiNbO_3$  films**

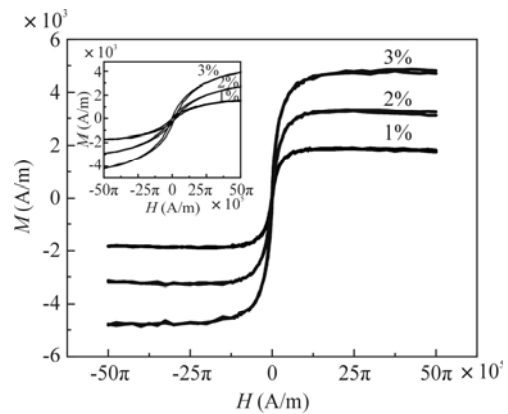
Fig.2 shows the Ca 2p, Li 1s (with the nearby Nb 4s level), Nb 3d and O 1s XPS spectra of the 2% Ca-doped  $LiNbO_3$  films. The charge-shifted spectra were corrected using the adventitious C 1s photoelectron signal at 284.6 eV. The Ca 2p XPS spectrum is depicted as shown in Fig.2(a). The peaks located at 346.78 eV and 350.78 eV can be ascribed to Ca  $2p_{3/2}$  and Ca  $2p_{1/2}$  binding energy, respectively. The Ca  $2p_{3/2}$  binding energy of the Ca-doped  $LiNbO_3$  film located at 346.78 eV is clearly different from that of metallic calcium (345.9 eV) and is very close to that of CaO (346.65 eV)<sup>[16]</sup>. This suggests that Ca atoms in the form of  $Ca^{2+}$  ions are incorporated into the  $LiNbO_3$  lattice. Fig.2(b) shows the Li 1s XPS spectrum. Li 1s binding energy is located at 55.08 eV, indicating that the Li remains in a typical  $Li^+$  valence. Fig.2(c) shows the Nb 3d XPS spectrum. The Nb  $3d_{5/2}$  and Nb  $3d_{3/2}$  peaks are located at 206.38 eV and 209.08 eV, respectively, indicating that Nb is  $Nb^{5+}$  valence<sup>[17]</sup>. Fig.2(d) shows the O 1s XPS spectrum. The peaks located at 529.56 eV and 531.75 eV respectively, can be ascribed to the  $LiNbO_3$  lattice oxygen and the oxygen deficient regions<sup>[18,19]</sup>.

Fig.3 shows the room temperature M-H curves of the Ca-doped  $LiNbO_3$  films with different Ca concentrations. From Fig.3, it can be seen that the Ca-doped  $LiNbO_3$  films exhibit the room temperature ferromagnetism and the clear hysteretic behavior at lower fields (the inset of Fig.3), and the saturation magnetization ( $M_s$ ) increases from 1870 A/m to 4800 A/m with Ca atom doping concentration from 1% to 3%. This proves that the Ca doping is effective for the room temperature ferromagnetism of  $LiNbO_3$  films. From above XRD and XPS results, we know that the Ca atoms enter the  $LiNbO_3$  lattice in the form of  $Ca^{2+}$  ions. Because the valence of Ca ions is different from that of Nb and Li ions, the O vacancies and Li vacancies are induced to compensate the charge non-equilibrium<sup>[20]</sup>. The Li vacancies could induce the spin-

polarization of O 2p electrons<sup>[21]</sup> and O vacancies would result in the exchange interactions between located electron



**Fig.2 (a) Ca 2p, (b) Li 1s, (c) Nb 3d and (d) O 1s XPS spectra of the 2% Ca-doped LiNbO<sub>3</sub> film**



**Fig.3 Magnetization curves of the 1%, 2% and 3% Ca-doped LiNbO<sub>3</sub> films at room temperature (Inset shows an enlarged view of M-H curves.)**

spin moments<sup>[22]</sup>, so the vacancies should be responsible for the ferromagnetism of the Ca-doped LiNbO<sub>3</sub> films. Recent first principle calculation indicated that the magnetic properties of the A (A = Ca, Sr, Ba) substitution for Li in LiNbO<sub>3</sub> result from the s orbital of Nb atoms around the doping element<sup>[23]</sup>. Anyways, the room-temperature ferromagnetism induced by the nonmagnetic element Ca doping in LiNbO<sub>3</sub> film is the instinct property and deserves careful investigation.

In summary, the undoped, 1%, 2% and 3% Ca-doped LiNbO<sub>3</sub> films were deposited on Si (111) substrates by radio frequency magnetron sputtering technique. The Ca-doped LiNbO<sub>3</sub> films show room-temperature ferromagnetism. The saturation magnetization increases from 1870 A/m to 4800 A/m with Ca atom doping concentration from 1% to 3%. The Ca atoms in the form of Ca<sup>2+</sup> ions enter the LiNbO<sub>3</sub> lattice and the Ca<sup>2+</sup> incorporation can induce the formation of the Li vacancies or O vacancies for the charge compensation. The vacancies are responsible for the intrinsic ferromagnetism of the Ca-doped LiNbO<sub>3</sub> films.

## References

- [1] M. Bibes and A. Barthelemy, *Nat. Mater.* **7**, 425 (2008).
- [2] N. Nagao and K. Iijima, *Vacuum* **83**, 1132 (2009).
- [3] G. Catalan and J. F. Scott, *Adv. Mater.* **21**, 2463 (2009).
- [4] N. A. Hill and A. Filippetti, *J. Magn. Magn. Mater.* **242**, 976 (2002).
- [5] R. Ramesh and N. A. Spaldin, *Nat. Mater.* **6**, 21 (2007).
- [6] J. F. Scott, *Science* **315**, 954 (2007).
- [7] N. A. Hill, *J. Phys. Chem. B.* **104**, 6694 (2000).
- [8] M. Fiebig, *J. Phys. D: Appl. Phys.* **38**, 123 (2005).
- [9] Ji Renlong, Zuo Changyun and Liu Jiwen, *Journal of Optoelectronics-Laser* **21**, 1337 (2010). (in Chinese)
- [10] Xiao Qing, An Yukai, Li Xiang, Xu Lili, Wu Yichen, Duan Lingshen and Liu Jiwen, *Journal of Optoelectronics-Laser* **22**, 1667 (2011). (in Chinese)

- [11] C. Song, C. Z. Wang, Y. C. Yang, X. J. Liu, F. Zeng and F. Pan, *Appl. Phys. Lett.* **92**, 262901 (2008).
- [12] C. Song, C. Wang, X. Liu, F. Zeng and F. Pan, *Cryst. Growth Des.* **9**, 1235 (2009).
- [13] C. Chen, F. Zeng, J. H. Li, P. Sheng, J. T. Luo and Y. C. Yang, *Thin Solid Films* **520**, 764 (2011).
- [14] P. Sheng, F. Zeng, G. S. Tang, F. Pan, W. S. Yan and F. C. Hu, *J. Appl. Phys.* **112**, 033913 (2012).
- [15] R. S. Weis and T. K. Gaylord, *Applied Physics A* **37**, 191 (1985).
- [16] Yue Liu, Tingting Gu, Yan Wang, Xiaole Weng and Zhongbiao Wu, *Catalysis Communications* **18**, 106 (2012).
- [17] Maëlen Aufray, Stéphane Menuel, Yves Fort, Julien Eschbach, Didier Rouxel and Brice Vincent, *J. Nanosci. Nanotechnol.* **9**, 4780 (2009).
- [18] Shiming Yan, Shihui Ge, Yalu Zuo, Wen Qiao and Li Zhang, *Scripta Mater.* **61**, 387 (2009).
- [19] S. M. Yan, S. H. Ge, W. Qiao, Y. L. Zuo, F. Xu and L. Xi, *J. Magn. Magn. Mat.* **323**, 264 (2011).
- [20] Zhongqiang Hu, Meiya Li, Benfang Yu, Ling Pei, Jun Liu, Jing Wang and Xingzhong Zhao, *J. Phys. D: Appl. Phys.* **42**, 185010 (2009).
- [21] Ensi Cao, Yongjia Zhang, Hongwei Qin, Ling Zhang and Jifan Hu, *Physica B* **410**, 68 (2013).
- [22] Carlos Díaz-Moreno, Rurik Farias, Abel Hurtado-Macias, Jose Elizalde-Galindo and Juan Hernandez-Paz, *J. Appl. Phys.* **111**, 07D907 (2012).
- [23] Yongjia Zhang, Jifan Hu, Hua Liu and Hongwei Qin, *Transactions on Magnetics* **47**, 2916 (2011).