# Optimal design and characteristics analysis of a polymer electro－optic switch with seven vertical－turning serial－cou－ pled microrings＊ 

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#### Abstract

A novel $1 \times 2$ polymer electro－optic（EO）switch based on seven vertical－turning serial－coupled microrings is proposed for dropping crosstalk and obtaining flat boxlike spectrum．The device structure，theory and formulation are presented， and the microring resonance order and coupling gaps are optimized．The switching voltage of the device for obtaining crosstalk lower than -30 dB under through state is decided to be about 1.86 V ．Under the operation voltages of 0 V （drop state）and 1.86 V （through state），the switching performance is characterized，and the output spectrum is ana－ lyzed．The calculation results show that the crosstalk at through state and that at drop state are -30.2 dB and -53.2 dB ， respectively，while the insertion losses are 0.86 dB and 3.18 dB ，respectively．Owning to the seven serial－coupled mi－ crorings resonance structure，the proposed switch reveals the favorable boxlike spectrum compared with the simple device with only one microring，and thus the crosstalk under drop state is improved from -26.8 dB to -53.2 dB ．Due to the low crosstalk，this device can be used in optical networks－on－chip for signal switching and routing．


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Electro－optic（EO）switches ${ }^{[1-6]}$ play a crucial role in the optical communication systems．Poled polymers ${ }^{[7-10]}$ are widely used in the fabrication of optical switches and modulators．In 2009，a polymer EO switch composed of a single microring coupled with two parallel channel waveguides was reported by our group ${ }^{[11]}$ ．In this paper， we propose a novel polymer EO switching configuration by using seven vertical－turning serial－coupled microrings， which is an improvement to the structure reported in Ref．［11］．

The structure model of the microring resonator（MRR） EO switch is shown in Fig．1（a），which consists of a horizontal channel，a vertical channel and seven vertical－ turning serial－coupled microrings．The cross section of MRR waveguide is shown in Fig．1（b）．The electrodes are only deposited on microrings，so the EO effect only oc－ curs in the waveguide cores of microrings．The waveguide layers are upper electrode，upper buffer layer，core，under buffer layer and lower electrode．By applying different op－ eration voltages on the microrings，the switching func－ tions can be realized between the horizontal channel and
the vertical channel．The device length and width are both $L_{1}+L_{2}+L_{3}$ ．During design，we select $L_{1}=100 \mu \mathrm{~m}$ ， $L_{3}=100+w_{\mathrm{c}} / 2$ ，and $L_{2}=3\left(2 R+w_{\mathrm{r}}\right)+R+w_{\mathrm{r}} / 2+3 h_{\mathrm{RR}}+h_{\mathrm{CR}}+w_{\mathrm{c}} / 2$ ， where $h_{\mathrm{RR}}$ is the coupling gap between two adjacent mi－ crorings，$h_{\mathrm{CR}}$ is that between the channel waveguide and MRR waveguide，and $w_{\mathrm{r}}$ and $w_{\mathrm{c}}$ are core widths of the MRR waveguide and channel waveguide，respectively．

（a）

[^0]

Fig. 1 (a) Structure model of the MRR EO switch based on seven vertical-turning serial-coupled microrings; (b) Cross-section of the MRR waveguide

Under 1550 nm , the related material parameters are taken as follows. The refractive index of the EO polymer material (AJL8/APC) is $n_{10}=1.59$, its bulk amplitude attenuation coefficient is $\alpha_{10}=0.25 \mathrm{~dB} / \mathrm{cm}$, and its EO coefficient is $\gamma_{33}=68 \mathrm{pm} / \mathrm{V}^{[12]}$. The refractive index of the upper/under buffer layer material [P(PFS-GMA)] is $n_{20}=1.461$, and its bulk amplitude attenuation coefficient is $\alpha_{20}=0.25 \mathrm{~dB} / \mathrm{cm}^{[13]}$. The electrode is made of aurum, its refractive index is $n_{30}=0.19$, and its bulk extinction coefficient is $\kappa_{30}=6.1^{[14]}$. Air is adopted as the left/right cladding layer, i.e., $n_{40}=1.0$ and $\alpha_{40}=0$.

The thickness of upper/under buffer layer is taken as $t_{2}=2.5 \mu \mathrm{~m}$, and the thickness of electrode is taken as $t_{3}=0.2 \mu \mathrm{~m}$. For assuring single-mode propagation in microring rectangular waveguide, its width and thickness are both taken as $w_{\mathrm{r}}=t_{1}=1.7 \mu \mathrm{~m}$. We take the ring bending radius as $R=19.49 \mu \mathrm{~m}$, and the calculated mode effective refractive index is about 1.5189 under the case of applying no voltage. We take the core thickness of the channel equal to that of the microring as $t_{1}=1.7 \mu \mathrm{~m}$, and select the width of the channel as $w_{\mathrm{c}}=1.85 \mu \mathrm{~m}$, and in this case, the channel and microring possess the identical mode effective refractive index. Under the parameters above, the mode loss in MRR waveguide is almost equal to that of the channel waveguide, which is $\alpha_{\mathrm{R}}=\alpha_{\mathrm{C}}=0.256 \mathrm{~dB} / \mathrm{cm}$.

For the seven vertical-turning serial-coupled microrings, define $\kappa_{\mathrm{RR}}$ and $t_{\mathrm{RR}}$ as the amplitude coupling ratio and amplitude transmission ratio between two adjacent microrings, respectively. Then the relation between the amplitudes $a_{1}, b_{1}$ and $a_{7}^{\prime}, b_{7}^{\prime}$ can be expressed by

$$
\left[\begin{array}{l}
a_{7}^{\prime}  \tag{1}\\
b_{7}^{\prime}
\end{array}\right]=\boldsymbol{P}_{\mathrm{R}} \boldsymbol{M P}_{\mathrm{R}}\left[\begin{array}{l}
a_{1} \\
b_{1}
\end{array}\right]=\boldsymbol{M}_{\mathrm{t}}\left[\begin{array}{l}
a_{1} \\
b_{1}
\end{array}\right],
$$

where

$$
\boldsymbol{P}_{\mathrm{R}}=\left[\begin{array}{cc}
0 & \exp \left(-\mathrm{j} \phi_{\mathrm{R}}\right) \\
\exp \left(\mathrm{j} \phi_{\mathrm{R}}\right) & 0
\end{array}\right]\left(\frac{1}{\mathrm{j} \kappa_{\mathrm{RR}}}\right)^{2} \times
$$

$$
\begin{align*}
& {\left[\begin{array}{cc}
-\exp \left(\mathrm{j} \phi_{\mathrm{R}}\right) & t_{\mathrm{RR}} \exp \left(-\mathrm{j} \phi_{\mathrm{R}}\right) \\
-t_{\mathrm{RR}} \exp \left(\mathrm{j} \phi_{\mathrm{R}}\right) & \exp \left(-\mathrm{j} \phi_{\mathrm{R}}\right)
\end{array}\right]^{2}, }  \tag{2}\\
\boldsymbol{M}= & \frac{1}{\mathrm{j} \kappa_{\mathrm{RR}}}\left[\begin{array}{cc}
t_{\mathrm{RR}} & -1 \\
1 & -t_{\mathrm{RR}}
\end{array}\right]\left[\begin{array}{cc}
0 & \exp \left(-\mathrm{j} \phi_{1}\right) \\
\exp \left(\mathrm{j} \phi_{2}\right) & 0
\end{array}\right] \times \\
& \frac{1}{\mathrm{j} \kappa_{\mathrm{RR}}}\left[\begin{array}{cc}
t_{\mathrm{RR}} & -1 \\
1 & -t_{\mathrm{RR}}
\end{array}\right], \tag{3}
\end{align*}
$$

where $\phi_{\mathrm{R}}=\pi R\left(\beta_{\mathrm{R}}-\mathrm{j} \alpha_{\mathrm{R}}\right), \phi_{1}=\frac{3}{2} \pi R\left(\beta_{\mathrm{R}}-\mathrm{j} \alpha_{\mathrm{R}}\right)$ and $\phi_{2}=$ $\frac{1}{2} \pi R\left(\beta_{\mathrm{R}}-\mathrm{j} \alpha_{\mathrm{R}}\right)$.

According to EO modulation theory, we can obtain the variation of refractive index $\Delta n_{10}$ of the EO core material resulting from the operation voltage $U$ as

$$
\begin{equation*}
\Delta n_{10}(U)=\frac{1}{2} n_{10}^{3} \gamma_{33} E_{1}=\frac{n_{10}^{3} n_{20}^{2} \gamma_{33} U}{2\left(2 n_{10}^{2} t_{2}+n_{20}^{2} t_{1}\right)}, \tag{4}
\end{equation*}
$$

where $U$ can be equal to zero or not. The refractive index of the EO core material $n_{10}$ will be changed to $n_{10}+\Delta n_{10}$. Then the mode propagation constant in each microring can be regarded as a function of $U$, denoted by $\beta_{\mathrm{R}}(U)$. Since $\phi_{1}, \phi_{2}, \phi_{\mathrm{R}}, \kappa_{\mathrm{RR}}$ and $t_{\mathrm{RR}}$ are all related to mode propagation constant, $\boldsymbol{M}_{\mathrm{t}}$ will be a function of $U$, that is $\boldsymbol{M}_{\mathrm{t}}=\boldsymbol{M}_{\mathrm{t}}(U)$.

The relation between $a_{0}, b_{0}$ and $a_{1}, b_{1}$ is

$$
\left[\begin{array}{l}
a_{1}  \tag{5}\\
b_{1}
\end{array}\right]=\frac{1}{\mathrm{j} \kappa_{\mathrm{CR}}}\left[\begin{array}{cc}
t_{\mathrm{CR}} & -1 \\
1 & -t_{\mathrm{CR}}
\end{array}\right]\left[\begin{array}{l}
a_{0} \\
b_{0}
\end{array}\right],
$$

and that between $a_{7}^{\prime}, b_{7}^{\prime}$ and $a_{8}, b_{8}$ is

$$
\left[\begin{array}{l}
a_{8}  \tag{6}\\
b_{8}
\end{array}\right]=\frac{1}{\mathrm{j} \kappa_{\mathrm{CR}}}\left[\begin{array}{cc}
t_{\mathrm{CR}} & -1 \\
1 & -t_{\mathrm{CR}}
\end{array}\right]\left[\begin{array}{l}
a_{7}^{\prime} \\
b_{7}^{\prime}
\end{array}\right],
$$

where $\kappa_{\mathrm{CR}}$ and $t_{\mathrm{CR}}$ are the amplitude coupling ratio and amplitude transmission ratio between MRR waveguide and channel waveguide, respectively. So the total transfer matrix of the switch can be written as

$$
\left[\begin{array}{c}
a_{8}  \tag{7}\\
b_{8}
\end{array}\right]=\frac{1}{\mathrm{j} \kappa_{\mathrm{CR}}}\left[\begin{array}{cc}
t_{\mathrm{CR}} & -1 \\
1 & -t_{\mathrm{CR}}
\end{array}\right] \boldsymbol{M}_{\mathrm{t}} \frac{1}{\mathrm{j} \kappa_{\mathrm{CR}}}\left[\begin{array}{cc}
t_{\mathrm{CR}} & -1 \\
1 & -t_{\mathrm{CR}}
\end{array}\right]=\boldsymbol{P}_{\mathrm{t}}\left[\begin{array}{l}
a_{0} \\
b_{0}
\end{array}\right],
$$

where $\boldsymbol{P}_{\mathrm{t}}=\frac{1}{\mathrm{j} \kappa_{\mathrm{CR}}}\left[\begin{array}{cc}t_{\mathrm{CR}} & -1 \\ 1 & -t_{\mathrm{CR}}\end{array}\right] \boldsymbol{M}_{\mathrm{t}} \frac{1}{\mathrm{j} \kappa_{\mathrm{CR}}}\left[\begin{array}{cc}t_{\mathrm{CR}} & -1 \\ 1 & -t_{\mathrm{CR}}\end{array}\right]$ is a $2 \times 2$ matrix, and its invert matrix is defined as $\boldsymbol{P}_{\mathrm{t}}^{-1}$. Eq.(7) can be further written as

$$
\left[\begin{array}{l}
a_{0}  \tag{8}\\
b_{0}
\end{array}\right]=\boldsymbol{P}_{\mathrm{t}}^{-1}\left[\begin{array}{l}
a_{8} \\
b_{8}
\end{array}\right] .
$$

Considering that $a_{8}=0, a_{0}=\left[\boldsymbol{P}_{\mathrm{t}}^{-1}(1,2)\right]^{-1} b_{8}$ and $b_{0}=$ $\boldsymbol{P}_{\mathrm{t}}^{-1}(2,2) b_{8}$, we obtain from Eq.(8) that

$$
\begin{equation*}
\frac{b_{8}^{\prime}}{a_{0}}=\left[\boldsymbol{P}_{\mathrm{t}}^{-1}(1,2)\right]^{-1}, \tag{9}
\end{equation*}
$$

$$
\begin{equation*}
\frac{b_{0}}{a_{0}}=\frac{\boldsymbol{P}_{\mathrm{t}}^{-1}(2,2)}{\boldsymbol{P}_{\mathrm{t}}^{-1}(1,2)} . \tag{10}
\end{equation*}
$$

Combing the following relations of $a_{0}=a_{0}^{\prime} \exp \left(-\mathrm{j} \psi_{1}\right)$, $b_{0}^{\prime}=b_{0} \exp \left[-\mathrm{j}\left(\psi_{1}+\psi_{2}\right)\right]$ and $b_{8}^{\prime}=b_{8} \exp \left(-\mathrm{j} \psi_{1}\right)$, the output light amplitudes from the drop port and through port can be finally expressed as

$$
\begin{align*}
\frac{b_{8}^{\prime}}{a_{0}^{\prime}}= & \frac{b_{8} \exp \left(-\mathrm{j} \psi_{1}\right)}{a_{0} \exp \left(\mathrm{j} \psi_{1}\right)}=\left[\boldsymbol{P}_{\mathrm{t}}^{-1}(1,2)\right]^{-1} \exp \left(-\mathrm{j} 2 \psi_{1}\right),  \tag{11}\\
\frac{b_{0}^{\prime}}{a_{0}^{\prime}}= & \frac{b_{0} \exp \left[-\mathrm{j}\left(\psi_{1}+\psi_{2}\right)\right]}{a_{0} \exp \left(\mathrm{j} \psi_{1}\right)}= \\
& \frac{\boldsymbol{P}_{\mathrm{t}}^{-1}(2,2)}{\boldsymbol{P}_{\mathrm{t}}^{-1}(1,2)} \exp \left[-\mathrm{j}\left(2 \psi_{1}+\psi_{2}\right)\right] . \tag{12}
\end{align*}
$$

where $\psi_{1}=L_{1}\left(\beta_{\mathrm{C}}-\mathrm{j} \alpha_{\mathrm{C}}\right)$ and $\psi_{2}=\left(L_{2}+L_{3}\right)\left(\beta_{\mathrm{C}}-\mathrm{j} \alpha_{\mathrm{C}}\right)$.
The output power from the two ports in dB form will be

$$
\begin{align*}
& P_{\mathrm{D}}=20 \lg \left|\frac{b_{8}^{\prime}}{a_{0}^{\prime}}\right|,  \tag{13}\\
& P_{\mathrm{B}}=20 \lg \left\lvert\, \frac{\left|\frac{b_{0}^{\prime}}{a_{0}^{\prime}}\right| .}{} .\right. \tag{14}
\end{align*}
$$

The switching characteristics are described below.
(1) Under $U=0 \mathrm{~V}$, the switch is operated at drop state, the drop port is named as ON-port, and the through port is named as OFF-port. The insertion loss and crosstalk under this state are characterized as

$$
\begin{align*}
& I L_{\text {drop }}=\left.P_{\mathrm{D}}\right|_{U=0},  \tag{15}\\
& C T_{\text {drop }}=\left.\left(P_{\mathrm{B}}-P_{\mathrm{D}}\right)\right|_{U=0} . \tag{16}
\end{align*}
$$

(2) Under the switching voltage, i.e., $U=U_{\mathrm{s}}$, the switch is operated at through state, the vertical output port is named as OFF-port, and the horizontal output port is named as ON-port. The insertion loss and crosstalk under this state are defined as

$$
\begin{align*}
& I L_{\mathrm{thr}}=\left.P_{\mathrm{B}}\right|_{U=U},  \tag{17}\\
& C T_{\mathrm{thr}}=\left.\left(P_{\mathrm{D}}-P_{\mathrm{B}}\right)\right|_{U=U} . \tag{18}
\end{align*}
$$

The resonance order of microring is a key parameter that affects the performance of the MRR EO switch, which should be firstly optimized for obtaining low crosstalk. By using Eqs.(13) and (14), Fig.2(a) shows the output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ versus operation wavelength $\lambda$ under different resonance orders, where $h_{\mathrm{RR}}=0.68 \mu \mathrm{~m}$, $h_{\mathrm{CR}}=0.14 \mu \mathrm{~m}$ and $U=0 \mathrm{~V}$. It can be found from Fig.2(a) that under different resonance orders, the crosstalks of $P_{\mathrm{B}}-P_{\mathrm{D}}$ under wavelength of 1550 nm are also different. The curves of crosstalk and insertion loss versus resonance order are further plotted in Fig.2(b). We can find that when $m=120$, the crosstalk becomes the smallest, and the insertion loss is nearly a constant. So in practical design, we choose $m=120$, and the corresponding ring radius $R$ is determined to be $\sim 19.49 \mu \mathrm{~m}$.


Fig. 2 (a) Curves of output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ versus operation wavelength $\lambda$ under different resonance orders; (b) Curves of crosstalk and insertion loss under drop state versus resonance order, where $h_{R R}=0.68 \mu \mathrm{~m}, h_{C R}=0.14 \mu \mathrm{~m}$ and $U=0 \mathrm{~V}$

The central coupling gap $h_{\mathrm{CR}}$ between channel waveguide and MRR waveguide and the central coupling gap $h_{\text {RR }}$ between two adjacent microrings also need to be selected properly. By means of Eqs.(13) and (14), Fig.3(a) shows the output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ versus operation wavelength $\lambda$ under different $h_{\mathrm{RR}}$, where $h_{\mathrm{CR}}=0.14 \mu \mathrm{~m}$ and $h_{\mathrm{RR}}=0.38$ $\mu \mathrm{m}, 0.48 \mu \mathrm{~m}, 0.58 \mu \mathrm{~m}, 0.68 \mu \mathrm{~m}$ and $0.78 \mu \mathrm{~m}$. It can be found that $h_{\mathrm{RR}}$ affects the output spectral characteristics, such as flat property, insertion loss and crosstalk. When $h_{\mathrm{RR}}=0.38 \mu \mathrm{~m}$ and $0.48 \mu \mathrm{~m}$, the spectrum of the drop port is not flat, and the crosstalk is larger than -30 dB . When $h_{\mathrm{RR}}=0.58 \mu \mathrm{~m}$ or $0.78 \mu \mathrm{~m}$, though the spectrum of the drop port becomes flat, the crosstalk is still larger than -30 dB . Fig.3(b) shows the curves of the crosstalk and insertion loss versus $h_{\text {RR }}$. When $h_{R R}=0.68 \mu \mathrm{~m}$, the crosstalk becomes the smallest, and under this state, the spectrum is also flat, as seen from Fig.3(a).

Fig.3(c) shows the output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ versus operation wavelength $\lambda$ under different $h_{\mathrm{CR}}$, where $h_{\mathrm{RR}}=0.68 \mu \mathrm{~m}$, $h_{\mathrm{CR}}=0.04 \mu \mathrm{~m}, 0.14 \mu \mathrm{~m}, 0.24 \mu \mathrm{~m}$ and $0.34 \mu \mathrm{~m}$. It can be found that $h_{\mathrm{CR}}$ slightly affects the output spectrum, especially the flat property, but it obviously influences the crosstalk. Fig.3(d) shows the curves of crosstalk and insertion loss versus $h_{\mathrm{CR}}$. When $h_{\mathrm{CR}}=0.14 \mu \mathrm{~m}$, the crosstalk becomes the smallest, and the insertion loss nearly be-
comes the smallest. Therefore, we select $h_{\mathrm{CR}}=0.14 \mu \mathrm{~m}$ and $h_{\mathrm{RR}}=0.68 \mu \mathrm{~m}$.


Fig. 3 Curves of output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ versus operation wavelength $\lambda$ under different central coupling gaps: (a) $h_{\mathrm{CR}}=0.14 \mu \mathrm{~m}, \lambda=1550 \mathrm{~nm}, h_{\mathrm{RR}}=0.38 \mu \mathrm{~m}, 0.48$ $\mu \mathrm{m}, 0.58 \mu \mathrm{~m}, 0.68 \mu \mathrm{~m}$ and $0.78 \mu \mathrm{~m}$, and (c) $\boldsymbol{h}_{\mathrm{RR}}=0.68$ $\mu \mathrm{m}, \lambda=1550 \mathrm{~nm}, h_{\mathrm{CR}}=0.04 \mu \mathrm{~m}, 0.14 \mu \mathrm{~m}, 0.24 \mu \mathrm{~m}$ and $0.34 \mu \mathrm{~m}$; Curves of $C T_{\text {dorp }}$ and $I L_{\text {dorp }}$ versus (b) $h_{R R}$ and (d) $h_{\mathrm{CR}}$, where $\lambda=1550 \mathrm{~nm}$

Next, with the optimized parameters above, we investigate the performance of the switch. Fig.4(a) shows the curves of output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ versus the operation voltage $U$. We can see that when $U=0 \mathrm{~V}, P_{\mathrm{D}}$ is the maximum, and the device is operated at drop state. When $U$ increases, $P_{\mathrm{B}}$ increases, $P_{\mathrm{D}}$ decreases, and the through state can be realized. Fig.4(b) shows the result of $P_{\mathrm{D}}-P_{\mathrm{B}}$ versus the applied voltage $U$, and the relation between $P_{\mathrm{D}}-P_{\mathrm{B}}$ and the device's crosstalk of $C T_{\text {drop }}$ and $C T_{\text {thr }}$ can be described as

$$
\begin{align*}
& P_{\mathrm{D}}-P_{\mathrm{B}}=-C T_{\text {drop }}, P_{\mathrm{D}}-P_{\mathrm{B}}>0,  \tag{19}\\
& P_{\mathrm{D}}-P_{\mathrm{B}}=C T_{\text {thr }}, P_{\mathrm{D}}-P_{\mathrm{B}}<0 . \tag{20}
\end{align*}
$$

When the operation voltage is equal to the switching voltage, the output power from horizontal channel will be the largest, while that output from the vertical channel will be the smallest, thus the switching function is realized. We define the switching voltage for the MRR EO switch as follows for obtaining $C T_{\mathrm{thr}}=-30 \mathrm{~dB}$, i.e.,

$$
\begin{equation*}
U_{\mathrm{s}}=\left.U_{\mathrm{s}}\right|_{C T_{\mathrm{tut}}=-30 \mathrm{~dB}} . \tag{21}
\end{equation*}
$$

Then the switching voltage for the device is about 1.86 V .
Fig. 5 shows the output spectra of the switch under two operation states ( $U=0 \mathrm{~V}$ for drop state and $U=1.86 \mathrm{~V}$ for through state). The simulation results indicate that under the applied voltages of 0 V and 1.86 V , the switch presents good switching functions at 1550 nm , and it reveals the boxlike spectral responses for the drop port and the extremely low crosstalk for both states. The crosstalk at through state is -30.2 dB , while that at drop state is
about -53.2 dB . The insertion losses at through and drop states are 0.86 dB and 3.18 dB , respectively.


Fig. 4 (a) Curves of output power $P_{\mathrm{B}}$ and $P_{\mathrm{D}}$ from the drop and through ports versus the applied voltage $U$ under 1550 nm ; (b) Curve of $P_{\mathrm{D}}-P_{\mathrm{B}}$ versus the applied voltage $U$


Fig. 5 Spectral responses of the MRR switch under drop and through states

In order to confirm the superiority of the MRR switch using seven microrings, we make a comparison on the output spectra of the two switches consisting of seven microrings and one microring, respectively, as shown in Fig.6. The inset of Fig. 6 shows the structure of the MRR switch with only one microring cross-coupled with the two channel waveguides. It can be seen that compared with the simple device with only one microring, the
switch with seven vertical-turning serial-coupled microrings reveals favorable flat boxlike spectrum, and the crosstalk is improved greatly from -26.8 dB to -53.2 dB .


Fig. 6 Comparison on the spectral responses of the two MRR switches (drop state) with seven microrings and only one microring, respectively (The inset shows the structure of the MRR switch with only one microring.)

As a conclusion, optimum design and characteristic analysis of a novel MRR EO switch based on seven ver-tical-tuning serial-coupled microrings are carried out under the operation wavelength of 1550 nm . By selecting proper values of resonance order and central coupling gaps, the crosstalk can be decreased as small as possible, and the output spectrum can be made as flat as possible. Also, compared with the switch with only one microring, the intensity of the non-resonance light can be reduced as weak as possible, and flat boxlike spectrum and low crosstalk can be achieved. Simulation results show that the driving voltage of the device is 1.86 V ; the crosstalk at through state is -30.2 dB , while that at drop state is about -53.2 dB ; the insertion loss at through state is
0.86 dB , while that at drop state is about 3.18 dB .

## References

[1] C. Chen, F. Zhang, H. Wang, X. Sun, F. Wang and D. Zhang, IEEE J. Quantum Electron. 47, 959 (2011).
[2] A. M. Yan, Y. N. Zhi, J. F. Sun and R. Liu, Appl. Phys. B 107, 421 (2012).
[3] Y. Enami, D. Mathin, C. T. DeRose, R. A. Norwood, J. Luo, A. K. Y. Jen and N. Peyghambarian, Appl. Phys. Lett. 94, 213513 (2009).
[4] A. Lorenz and H. S. Kitzerow, Appl. Phys. Lett. 98, 241106 (2011).
[5] J. Van Campenhout, W. M. Green, S. Assefa and Y. A. Vlasov, Opt. Express 19, 11568 (2011).
[6] J. Yang, W. Chen, W. Wang, M. Wang and J. Yang, J. Optoelectron. Laser 24, 16 (2013). (in Chinese)
[7] S. Huang, J. D. Luo, H. L. Yip, A. Ayazi, X. H. Zhou, M. Gould, A. T. Chen, T. Baehr-Jones, M. Hochberg and A. K. Y. Alex, Adv. Mater. 24, OP1 (2012).
[8] Y. Tsuboi, K. Tsuboi and T. Michinobu, J. Photopolymer Science Technol. 24, 305 (2011).
[9] C. Cabanetos, E. Blart, Y. Pellegrin, V. Montembault, L. Fontaine, F. Adamietz, V. Rodriguez and F. Odobel, Polymer 52, 2286 (2011).
[10] X. Yan, C. S. Ma, C. T. Zheng, X. Y. Wang and D. M. Zhang, Opt. Laser Technol. 42, 526 (2010).
[11] Z. Fan, B. Yun, G. Hu, Y. Yan and Y. Cui, J. Optoelectron. Laser 23, 1727 (2012). (in Chinese)
[12] G. Y. Xu, Z. F. Liu, J. Ma, B. Y. Liu, S. T. Ho, L. Wang, P. W. Zhu, T. J. Marks, J. D. Luo and A. K. Y. Jen, Opt. Exp. 13, 7380 (2005).
[13] C. Pitois, S. Vukmirovic, A. Hult, D. Wiesmann and M. Robertsson, Macromolecules 32, 2903 (1999).
[14] W. G. Driscoll and W. Vaughan, Handbook of Optics, New York: McGraw-Hill, 7 (1978).


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