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一种用于传感的多Fano通道高灵敏度 MIM 波导

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摘 要:提出了一种金属-绝缘体-金属波导结构,该结构由带有中央矩形空气路径的方形环谐振腔和 带有挡板的总线波导组成。利用有限元法研究了该结构的磁场分布、透射特性和传感性能。仿真结果 表明,谐振腔中的中央矩形空气路径可以改变表面等离极化激元在谐振腔中的传播路径,提供更多的 等离子体共振模式。在所提出的结构中可以激发四重Fano共振,透射谱中形成一个滤波带。改变结构 参数可以调节Fano共振的个数,最多可以获得6个Fano共振和两个滤波带,还可以对Fano共振的位 置、强度以及滤波带的宽度进行灵活方便的调节。该结构的最大灵敏度和品质因数分别为3028 nm/RIU 和157.14,可用于制作多波段带阻滤波器和检测葡萄糖等液体浓度的传感器。

关键词:表面等离激元;金属-绝缘体-金属波导;多重Fano共振;高灵敏度;滤波器
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0 引言

随着科技的发展和进步,人们对器件的微型化和集成化的要求越来越高。由于光学衍射极限的存在, 传统光子器件的光学性能受到极大的限制。表面等离极化激元(Surface Plasmon Polaritons, SPPs)是入射 光波和金属相互作用激发的电磁波^[1-2],它沿金属-绝缘体表面传播,并在垂直于界面的两侧呈指数衰减^[3-4]。 SPPs可以克服传统的光学衍射极限并将光限制在亚波长尺寸范围内,因此它在高密度光子集成电路中的传 输、处理和控制方面具有非常重要的应用^[5-6]。

近年来,研究人员已经提出了许多基于 SPPs 的微纳结构,例如金属-绝缘体-金属(Metal-insulatormetal, MIM)波导、纳米狭缝、混合布拉格波导、金属纳米粒子波导和等离子体纳米团簇等。MIM 波导具有 欧姆损耗低、传播距离长和易于制作等优点,可用于纳米传感器^[7-11]、光开关^[12-13]、滤波器^[14]和解复用器^[15]等 光学器件,它已成为最有前途的波导结构之一^[16]。

1961年,UGO Fano首次提出了Fano共振,认为原子系统中宽的连续态和窄的离散态相互干涉导致了Fano共振的形成^[17-19]。近年来,研究人员发现,在一些等离激元纳米结构中也可以产生Fano共振。与洛伦兹线型不同,Fano共振具有尖锐且不对称的线型,对周围环境折射率也非常敏感^[20-21]。在MIM 波导中,波导结构对Fano共振的传输特性有很大的影响。例如,ZAFAR R等^[22]研究了与一对椭圆环形谐振腔耦合的MIM 波导中的Fano共振,其最大灵敏度为1100 nm/RIU;RAHMATIYAR M等^[23]研究了一个圆锥缺陷环形谐振腔与一个具有锥形缺陷的 MIM 等离子体波导耦合的折射率传感器,其最大灵敏度为1295 nm/RIU;ZHU J^[24]等研究了由 MIM 波导和半圆谐振腔与键形谐振腔耦合组成的波导结构的折射率传感器,灵敏度最大可达到1261.67 nm/RIU。与单个Fano系统相比,多重Fano共振可以实现滤波器的功能,但一个结构同时具有多波段滤波功能和优秀传感特性的情况还比较少见。

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基于 SPPs 的 MIM 波导系统中使用的谐振腔通常采用方形,因为它的外观简单且易于制造。本文提出了 一种由带有中央矩形空气路径的方形环谐振腔和带有挡板的总线波导组成的 MIM 波导,利用基于有限元法 (Finite Element Method, FEM)的 COMSOL Multiphysics 5.4 软件对该结构的光学特性进行了研究,并根据透射 光谱和磁场分布探索了 Fano共振的形成机制。与单个方形腔和方形环腔不同,带有中央矩形空气路径的方形 环谐振腔可以增加谐振腔的有效腔长并改变 SPPs 在谐振腔中的传播路径,提供了更多的等离子体共振模式。 仿真结果表明,所提出的结构中可激发四重 Fano共振,透射光谱中出现了一个滤波带。改变方形环的边长,最 多可以获得六重 Fano共振和两个滤波带,改变结构参数可以调节 Fano共振的位置、强度和滤波带的宽度,该结 构还具有良好的灵敏度和 FOM 值,可用于制作多波段带阻滤波器和检测葡萄糖等液体浓度的传感器。

1 结构及仿真方法

图 1 为所提 MIM 波导的二维(2D)结构示意图,该结构由带有挡板的总线波导和一个带有中央矩形空 气路径的方形环谐振腔组成。方形环的边长为L,中央矩形空气路径的高度为h,谐振腔与带有挡板的总线 波导间的耦合距离为g,总线波导中挡板宽度为d,带有中央矩形空气路径的方形环谐振腔和总线波导的宽 度都为w。文中所有结果均是使用 COMSOL Multiphysics 5.4 模拟得到,输入端口(Input)和输出端口 (Ouput)分别位于 MIM 波导的左侧和右侧。该结构在y方向的厚度足够大(远大于光波长),3D 模型的仿真 结果与 2D 模型的仿真结果基本相同,为了减少计算量,均采用 2D 结构来进行模拟计算^[26-27]。



图1 所提出MIM波导二维(2D)结构示意

Fig. 1 Two-dimensional (2D) structure schematic of the proposed MIM waveguide

图 1 中,蓝色部分和白色部分分别代表银和空气。由于银功耗较低,因此选择它作为 MIM 波导的金属 材料。Ag的相对介电常数(ε_m)可以通过 Drude 模型来描述^[28,29]

$$\boldsymbol{\varepsilon}_{\mathrm{m}}(\boldsymbol{\omega}) = \boldsymbol{\varepsilon}_{\infty} - \frac{\boldsymbol{\omega}_{\mathrm{p}}^{2}}{\boldsymbol{\omega}^{2} + \mathrm{i}\boldsymbol{\omega}\boldsymbol{\gamma}}$$
(1)

式中, ω 是入射光频率, ω_p =9.10 eV, ε_{∞} =3.7 和 γ =0.018 eV分别是等离子体振荡频率、入射光频率无穷大时对应的介电常数和电子碰撞频率。

方形环腔可视为法布里-珀罗腔,当SPPs耦合到环形谐振腔中并满足谐振条件时,就会发生谐振,在谐振腔中形成稳定的驻波场。由驻波理论可得^[30]

$$\lambda_{\rm res} = \frac{2\text{Re}(n_{\rm eff})L_{\rm eff}}{N - \frac{\varphi}{\pi}} N = 1, 2, 3, \cdots$$
(2)

式中,N表示谐振的阶数, L_{eff} 表示有效腔长, φ 为SPPs在介质金属界面发生反射时引起的相位变化,Re(n_{eff}) 是有效折射率的实部。Re(n_{eff})可表示为

$$\operatorname{Re}\left(n_{\text{eff}}\right) = \left(\varepsilon_{\text{m}} + \left(\frac{k}{k_{0}}\right)^{2}\right)^{\frac{1}{2}}$$
(3)

式中, $k = \frac{2\pi}{\lambda} \pi k_0$ 分别是波导和自由空间中的波矢量, ϵ_m 是银的相对介电常数。

2 结果与讨论

2.1 Fano共振和滤波带的产生

图 2 为四种结构的透射光谱, Structure Ⅰ为带有挡板的总线波导, Structure Ⅱ, Structure Ⅲ和 The whole system分别由方形谐振腔、方形环谐振腔和带有中央空气路径的方形环谐振腔与带有挡板的总线波导组成。为了确保结构中仅传播横向磁模式,带有中央矩形空气路径的方形环谐振腔和总线波导的宽度都固定为 w=50 nm, 其它初始结构参数分别为 h=50 nm,g=10 nm,L=400 nm和 d=50 nm。



图 2 不同结构的透射光谱 Fig. 2 Transmission spectra of different structures

图 2(a)是带有挡板的总线波导产生的宽的连续态。从图 2 可以看出,后三种结构都产生了 Fano 共振,且随着方形谐振腔、方形环谐振腔和带有中央矩形空气路径的方形环谐振腔的引入,产生的 Fano 共振数量逐渐增加。图 2(b)中只有一个位于 950 nm 的 Fano 共振,是由方形腔产生的离散态和带有挡板的总线波导产生的连续态耦合而成。图 2(c)中共有三个 Fano 共振,分别位于 1 910 nm,995 nm 和 935 nm,是由方形环腔产生的离散态与带有挡板的总线波导产生的连续态耦合产生。图 2(d)中产生了四个 Fano 共振,分别位于 1 885 nm, 1 495 nm,1 215 nm 和 990 nm,将称其为 FR1、FR2、FR3 和 FR4。全系统中产生了更多的 Fano 共振,这是因为全系统在内部纳米盘中央引入了一条矩形空气路径,改变了 SPPs 在谐振腔中的传播路径,提供了更多的等离子体共振模式。Fano 共振谱线的显著不对称性在纳米光子器件中有很多应用,如图 2(d)中,由于 FR2 和 FR3 的方向相反,透射光谱在 1 200~1 500 nm 之间呈 U型,因此所设计的结构可以实现带阻滤波功能。带阻宽度定义为透射率小于 1%的波长范围,由图 2(d)可知,阻带左侧 λ_c 为关闭位置,阻带右侧 λ_0 为打开位置,命名其为阻带 1。则阻带宽度可以定义为 $\Delta l = \lambda_0 - \lambda_c$ 。对本结构, $\lambda_c = 1$ 260 nm, $\lambda_0 = 1$ 455 nm,可以算出阻带1中心波长为1 357.5 nm,阻带 1宽度 Δl 为 195 nm,因此,该结构可用于制作带阻滤波器。

为了更深入地了解所提出结构的Fano共振的形成机制,图3给出了全系统中产生的FR1、FR2、FR3和 FR4峰值位置的磁场分布。如图3(a)~(d)所示,在各个Fano峰值波长处,带有空气路径的方形环谐振腔中 均出现了驻波,电磁波能量主要集中在谐振腔中。具体而言,FR1处能量主要集中在谐振腔上下两侧,FR2 处能量主要集中在谐振腔的左右两侧,FR3处能量集中在谐振腔水平部分,FR4处的能量主要集中在谐振 腔的四角处。全系统中FR1、FR2、FR3和FR4的磁场中的节点数分别为2、2、3和4,不同的节点数是由于波 长不同时,总线波导和谐振腔中波数不同而引起的。



图3 全系统中不同峰值位置的磁场分布

2.2 Fano共振和滤波带的调节

结构参数对Fano共振的强度和位置有着强烈的影响。因此研究了全系统结构参数对其透射光谱的影响,未提及参数均保持不变。

首先,研究了全系统中耦合距离g和空气路径宽度h对透射光谱的影响。图4(a)给出了耦合距离g以5 nm 的步长从0增大到20 nm 时全系统的透射光谱。g=0 nm 时,耦合效应最强,透射率显示出剧烈的振荡。耦合距离g直接影响离散态和连续态之间的耦合效果,随着g的增加,耦合效应减弱,因此各Fano共振透射率和半高全宽(Full Width at Half Maximum, FWHM)减小,各Fano共振发生轻微蓝移。由图4(a)还可以得知,随着g从5 nm 增大到20 nm,FR2和FR3之间的阻带1宽度由170 nm 增大到230 nm,阻带1中心波长由1210 nm 移动到1345 nm。

图 4(b)给出的是空气路径宽度 h以 10 nm 的步长从 40 nm 增大到 80 nm 时全系统的透射光谱。随着 h的 增大,各 Fano共振发生不同程度的蓝移。空气路径可以改变 SPPs 在谐振腔中的传播路径,提供更多的等离 子体共振模式。由图 4(b)可以得知,随着 h的增加,FR1和 FR3 几乎不移动,FR4 只发生轻微的蓝移,而 FR2 发生明显的蓝移。改变 h不仅可以实现对 FR2的独立调控,同时还调节了滤波宽度和中心波长。由图 4(b)可 以得知,随着 h的降低,阻带 1的中心波长逐渐向低频方向移动,阻带 1宽度最大可以增加到 230 nm。

接着,研究了全系统中方形环边长L和总线波导中挡板的宽度d的透射光谱。图5(a)给出了全系统中 方形环边长L以50 nm的步长从300 nm增大到600 nm时的透射光谱,观察到随着L的增加透射光谱具有更 多的Fano共振,且各Fano共振发生显著的红移。根据式(2),红移现象可归因于谐振腔有效长度L_{eff}的增

Fig. 3 Magnetic field distribution at different peak positions in the whole system



图4 耦合距离g和中央空气路径宽度h对全系统透射光谱的影响 Fig. 4 Influence of coupling distance g and central air path width h on transmission spectrum of the whole system



图5 方形环边长L和总线波导中挡板宽度d对全系统透射光谱的影响

Fig. 5 Influence of side length *L* of square ring and baffle width *d* in bus waveguide on transmission spectrum of the whole system 加。因此,可以通过改变 *L*将所需的 Fano 共振峰移动到特定波长。对于 *L*=300,350,400,450,500,550和 600 nm, Fano 共振的数量分别为 2,3,4,4,5,6 和 6。各 Fano 峰对 *L* 比较敏感,从图中可以看出当 *L* 从 300 nm 增大到 600 nm 时, FR1 从 1 320 nm 移动到 3 040 nm, FR2 从 1 060 nm 移动到 2 400 nm。当 *L*=

500 nm时,在945 nm处出现一个新的Fano,形成了一个新的阻带,命名为阻带2,从而为实现多波段滤波提供了可能。当L增加时,两个阻带中心波长都向着低频方向移动,阻带1宽度逐渐增加到230 nm,阻带2宽度最大可达到275 nm。因此,改变边长L可以有效调节阻带中心波长和阻带宽度。

图 5(b)给出的是总线波导中挡板的宽度 d 以 10 nm 的步长从 10 nm 增大到 70 nm 时全系统的透射光谱。 d的大小影响总线波导和带有中央空气路径的方形环谐振腔之间的耦合效果,从图 5(b)可以看出,随着 d 的增加,FR1~FR4的共振峰位置几乎不发生移动,其中 FR1~FR3 的透射率随着 d 的增加变化较小,FR4强度有较大增强。综合考虑透射率、FWHM 以及阻带宽度,各个参数分别选取 d=50 nm,g=10 nm 和 h=50 nm。

3 Fano共振的灵敏度和FOM 值

等离子体折射率传感器的原理是基于共振波长会随着周围材料折射率的变化而变化。在实际应用中, 折射率还会随着被测材料的温度、湿度和浓度等的变化而变化。灵敏度和品质因数(Figure of Merit, FOM) 是评价折射率传感器性能的两个重要的指标。灵敏度S的定义为^[31]

$$S = \frac{\Delta \lambda}{\Delta n} \tag{4}$$

式中, Δn 表示折射率的变化, $\Delta \lambda$ 表示折射率变化前后Fano共振峰对应的波长变化,S的单位为nm/RIU。 FOM定义为^[32]

$$FOM = \frac{S}{FWHM}$$
(5)

FWHM表示共振峰的半高全宽。FOM是一个无量纲参数,由式(5)可知,可以通过提高灵敏度和降低 FWHM来优化传感性能。

为分析整个结构的传感特性,设置波导和谐振腔中的折射率n在1.00~1.08内变化,间隔为0.02。从图6(a)、6(b)可以看出,当填充介质的折射率从1.00变为1.08时,图中的透射光谱显著红移。这与式(2)



(c) The relationship between the resonant wavelength of the whole system and the refractive index n

图 6 不同折射率 n 对全系统的透射光谱的影响 Fig. 6 Effect of different refractive index n on transmission spectrum of the whole system

 $\lambda_{res} = 2 \operatorname{Re}(n_{eff}) L_{eff} / (N - \frac{\varphi}{\pi})$ 得到的理论结果是一致的,随着有效折射率的增加, λ_{res} 也在增加。图 6(c)进一 步表明FR1~FR4的共振波长与折射率的变化基本呈线性关系。当L=400 nm,根据式(4)和式(5)计算得 到的FR1~FR4的折射率灵敏度S分别为1900,1500,1200和900 nm/RIU,而品质因数FOM分别为15.6, 57.7,40和69.2。

改变参数L时,可以获得更多的Fano共振,为了探究参数L对结构性能的影响,分别计算了全系统中方 形环边长L以50 nm的步长从300 nm 增大到600 nm 时各 Fano 共振的共振波长λ_{res},灵敏度 S 和品质因数 FOM 如表1。

| Table 1 | Resonance wavelength A | res, sensitivity 5 | and FOM of Fal | to resonance of | the whole syste | em at differen | ΠL |
|----------------------|---|--------------------|----------------|-----------------|-----------------|----------------|----------|
| | | $FR1_4$ | $FR2_6$ | FR3 | FR4 | FR5 | FR6 |
| L=300 nm | $\lambda_{ m res}/ m nm$ | 1 344 | 1 053 | | | | |
| | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 1 350 | 1 050 | | | | |
| | $\rm FOM/RIU^{-1}$ | 22.7 | 26.6 | | | | |
| | $\lambda_{ m res}/ m nm$ | 1 611 | 1 275 | 1 028 | | | |
| L = 350 nm | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 1 600 | 1 250 | 1 000 | | | |
| | FOM/RIU^{-1} | 21.9 | 50 | 32.3 | | | |
| | $\lambda_{ m res}/ m nm$ | 1 884 | 1 494 | 1 218 | 988 | | |
| $L{=}400 \text{ nm}$ | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 1 900 | 1 500 | 1 200 | 900 | | |
| | $\rm FOM/RIU^{-1}$ | 69.2 | 40 | 57.7 | 15.6 | | |
| | $\lambda_{ m res}/ m nm$ | 2 174 | 1 716 | 1 402 | 1 124 | | |
| $L{=}450 \text{ nm}$ | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 2 200 | 1 700 | 1 400 | 1 100 | | |
| | $\rm FOM/RIU^{-1}$ | 8.0 | 54 | 47.5 | 157.1 | | |
| | $\lambda_{ m res}/ m nm$ | 2 472 | 1 942 | 1 586 | 1 260 | 947 | |
| $L{=}500 \text{ nm}$ | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 2 450 | 1 950 | 1 550 | 1 250 | 950 | |
| | $\mathrm{FOM}/\mathrm{RIU}^{-1}$ | 12.9 | 36 | 47 | 73.5 | 21 | |
| | $\lambda_{ m res}/ m nm$ | 2 766 | 2 171 | 1 771 | 1 398 | 1 051 | 932 |
| $L{=}550 \text{ nm}$ | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 2 750 | 2 200 | 1 750 | 1 350 | 1 050 | 900 |
| | $\mathrm{FOM}/\mathrm{RIU}^{-1}$ | 15.1 | 32.3 | 38 | 75 | 15.2 | 16.7 |
| | $\lambda_{ m res}/ m nm$ | 3 050 | 2 400 | 1 955 | 1 535 | 1 170 | $1\ 015$ |
| L=600 nm | $S/(\mathrm{nm}\cdot\mathrm{RIU}^{-1})$ | 3 000 | 2 500 | 2 000 | 1 500 | 1 000 | $1\ 000$ |
| | $\rm FOM/RIU^{-1}$ | 21.4 | 50 | 36.4 | 75 | 13.3 | 13.5 |

表1 不同L时全系统的Fano共振的共振波长 λ_{res} ,灵敏度S和品质因数FOM

citivity S and FOM of F 41 3 0.41 . .

在全系统中,可以获得多个具有高灵敏度的Fano共振。如果L很小,制造上较为困难;如果L很大,波导 的欧姆损耗、Fano共振的FWHM和结构尺寸变大,会导致透射率较低。在模拟中,L的范围选择为300 nm< L<600 nm,当L=450 nm时,可以得到最大FOM值为157.1。其性能优于表2中的大多数结构^[33-39]。

表2 本文提出结构与近年文献中波导结构的传感性能比较

| Table 2 | Comparison of sensing performance between the proposed structure and the waveguide structure in recent literature | | | | | | | |
|---------|---|-------------------------|--------------------|-----------|--|--|--|--|
| Year | Structure | $S/(nm \cdot RIU^{-1})$ | $\rm FOM/RIU^{-1}$ | Ref. | | | | |
| 2020 | Two unequal vertical rectangular cavities | 2 625.87 | 26.04 | [33] | | | | |
| 2021 | A connected-concentricdouble rings resonator | 2 260 | 56.5 | [34] | | | | |
| 2021 | A racetrack ring resonator | 1 774 | 61 | [35] | | | | |
| 2022 | A U-shaped ring resonator | 2 020 | 53.16 | [36] | | | | |
| 2022 | A semi-circular resonant cavity | 579 | 12.46 | [37] | | | | |
| 2022 | A rectangular root and a double-ring with a rectangular cavity | 2 280 | 76.7 | [38] | | | | |
| 2022 | A toroidal cavity with abuilt-in elliptical rin | 2 220 | 58.7 | [39] | | | | |
| | A square ring resonator with a central rectangular air path | 3 028 | 157.14 | This work | | | | |

4 全系统应用与葡萄糖浓度传感器

基于 SPPs 的 MIM 波导可以应用于折射率传感器,用来检测液体的折射率范围,在生物和化学领域也 有广泛的应用。在生物医学中,常常通过检测葡萄糖浓度来检验血糖的高低,如糖尿病就是一种典型由血 糖过高引起的疾病,该疾病会对身体造成很多不良影响,很容易引发各种并发症。因此,检测人体中葡萄糖 浓度极为重要,本文研究了所提出结构在葡萄糖浓度检测中的应用,其中葡萄糖溶液填充在波导和谐振腔 中。已有实验数据表明^[40],葡萄糖浓度与其折射率之间存在特定的线性关系,葡萄糖浓度的折射率表示为 $n = 0.000 \, 118 \, 9C + 1.332 \, 305 \, 45$ (6)

式中,C是葡萄糖溶液的浓度,单位为g/L,n是葡萄糖溶液的折射率。式(6)阐明了葡萄糖浓度和折射率之 间的线性关系。图7(a)、(b)显示了当葡萄糖浓度C从0g/L增大到200g/L时结构的透射光谱。根据 式(6),葡萄糖溶液的折射率从1.33230545增加到1.35608345。结构的其他参数为L=600 nm,h=50 nm、 d=50 nm和g=10 nm。从图7(c)、(d)可以得出,随着葡萄糖浓度的增加,共振波λres发生红移,C和λres之间 具有近似线性的关系,这与式(6)可以很好地吻合。为了更好地评估葡萄糖浓度传感器的性能,用式(4)表 示其灵敏度。由式(4)可以算出,该传感器的灵敏度高达3028 nm/RIU。本文研究为制造折射率传感器提 供了一种新思路,该传感器简单、方便、性能良好,为生物光学传感器提供了一种新途径。



图7 葡萄糖浓度对结构的影响 Fig.7 Effect of glucose concentration on structure

5 结论

本文提出了一种由带有中央矩形空气路径的方形环谐振腔和带有挡板的总线波导组成的MIM波导结构。该结构的透射光谱中出现四个典型的非对称线型的Fano共振和一个滤波带。利用有限元法研究了该结构的磁场分布、透射特性和传感性能。通过优化结构最多可以获得6个Fano共振和两个滤波带。改变结构参数可以调节Fano共振的位置、强度和滤波带的宽度。波导对结构内填充介质的折射率很敏感,该结构的最大灵敏度和品质因数(FOM)分别为3028 nm/RIU和157.14。该结构可用于集成纳米光学和生物医学

领域,如多波段带阻滤波器和检测葡萄糖等液体浓度的传感器。

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A Multi-Fano Channel High Sensitivity MIM Waveguide for Sensing

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Abstract: Due to the existence of diffraction limit, the optical performance of traditional photonic devices is limited greatly. Surface Plasmon Polaritons (SPPs) can overcome the traditional optical diffraction limit and localize light in sub-wavelength range, so it has very important applications in the transmission, processing and control of optical waves in high-density photonic integrated circuits. For the advantages of low ohmic loss, long propagation distance and easy fabrication, MIM waveguide has become one of the most promising waveguides. Square resonators are often used in SPPs-based MIM waveguide systems because of simple structure and easy fabrication. Fano resonance has a sharp and asymmetric line shape and is very sensitive to the refractive index of surrounding environment. Compared with a single Fano system, multiple Fano resonances can realize multi-channel sensing and have the ability of parallel processing, which has attracted great attention. In this paper, a metal-insulator-metal waveguide structure consisting of a square ring resonator with a central rectangular air path and a bus waveguide with a baffle is proposed, which is studied by using COMSOL Multiphysics 5.4 based on the Finite Element Method (FEM). Because of its low power dissipation, silver is chosen as the metal material of MIM waveguide. The thickness of the structure is large enough (much larger than the optical wavelength), and the simulation results of 3D model and 2D model are basically the same, so the 2D model is adopted to reduce calculation quantity. The optical properties of the structure are studied, and the formation mechanism of Fano resonance is discussed according to the transmission spectrum and magnetic field distribution. In addition,

the influence of changing structural parameters is discussed. The application of the proposed structure in sensing is summarized. Firstly, the transmission spectra of a single square cavity, a square ring cavity and the proposed structure are given. The square ring resonator with a central rectangular air path can increase the effective cavity length and change the propagation path of SPPs in the resonator by providing more plasma resonance modes. Secondly, in order to understand the formation mechanism of Fano resonance of the proposed structure, the magnetic field distributions at the peak positions of Fano resonance are given. The simulation results show that the proposed structure can excite quadruple Fano resonance, and a filter band is formed in the transmission spectrum. The position and intensity of Fano resonance and the width of the filter band can be adjusted conveniently by changing the structure parameters. Thirdly, the number of Fano resonances can be adjusted by changing the side length of the square ring, and six Fano resonances and two filter bands can be obtained at most. At the same time, the filter width and the center wavelength can be regulated. The bandwidth of the proposed structure is defined as the wavelength range in which the transmittance is less than 1%, and the maximum bandwidth of the proposed structure is 275 nm. Therefore, the structure can be used for making band-stop filter. Finally, the application of MIM waveguide based on SPPs in refractive index sensor is studied. In biomedicine, it is very important to measure glucose concentration in the body, because it is often used to check the level of blood glucose. The proposed structure is very sensitive to the refractive index of the filled medium. The application of the proposed structure in the detection of glucose concentration is studied by investigating the relationship between the resonant wavelength of each Fano and glucose concentration. In order to evaluate the performance of glucose concentration sensor, the maximum sensitivity and Figure of Merit (FOM) of the proposed structure are calculated, which are 3 028 nm/RIU and 157.14 nm/RIU, respectively. The sensitivity and FOM value of the proposed structure are compared with that of MIM waveguide structure proposed in recent years.

Key words: Surface plasmon polaritons; Metal-insulator-metal waveguide; Multiple Fano resonance; High sensitivity; Filter

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