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近红外通讯波段的石墨烯/锗肖特基结光电探测器

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摘要: 制作了一种近红外通讯波段的石墨烯/锗肖特基结光电探测器。采用化学气相沉积法制备高质量石墨烯, 通过湿法转移法转移到 n 型锗表面, 获得高性能石墨烯/锗的肖特基结器件。仿真与实验结果表明, 石墨烯透明电极与锗衬底形成良好的肖特基接触, 大大提升了器件的光生载流子收集效率。该器件在无光照条件下, 整流比为 5.3×10^2 , 在光强为 0.3 mW/cm^2 的 1550 nm 近红外光的照射下, 开关比为 10^2 , 响应度和探测率分别可达 635.7 mA/W 、 $9.8 \times 10^{10} \text{ Jones}$ 。同时, 器件具备较快的响应速度, 在 3 dB 带宽处上升和下降时间分别为 $40 \mu\text{s}$ 和 $35 \mu\text{s}$ 。研究结果展示了高性能石墨烯/锗光电探测器在近红外光电系统中具有潜在的应用前景。

关键词: 光电探测器; 石墨烯; 能带分析; 锗; 肖特基结

中图分类号: TN215

文献标识码: A

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0 引言

石墨烯是碳原子经 sp^2 杂化形成的蜂窝状单层二维片层材料^[1], 它的厚度只有原子尺度 (0.35 nm), 是当前世界上已知最薄的二维材料^[2], 由于石墨烯具有良好的电学和光电性能, 它在电源电池^[3,4]、电容导体^[5,6]、传感器件^[7]、吸波材料^[8]、光电器件^[9,10]中显示出广阔的应用前景。同时, 硅材料成本低, 在地壳中含量丰富, 为光探测提供了一个通用的平台。通过石墨烯和硅半导体材料结合, 构成异质结光探测器^[9-14], 入射光可以很容易地穿透石墨烯薄膜, 到达结区, 激发的电子-空穴对被内置电场隔开, 形成光电流, 具有优异的光探测性能。

尽管硅基光电探测器已经取得了巨大的进展, 但硅的带隙是 1.12 eV , 导致硅基近红外光探测器的截止波长一般在 $1.1 \mu\text{m}$ 左右, 其探测波长范围相对较窄。为解决这个问题, 研究人员采用带隙比硅小得多的锗 (Ge) 材料, 其带隙为 0.67 eV , 截止波长一般在 $1.8 \mu\text{m}$ 左右, 利用 Ge 可以探测到更长波段的近红外光, 且 Ge 具有在近红外波段吸收系数大、成本低以及与硅工艺并行处理的良好兼容性等特点^[15]。到目前为止, 已经开发了许多不同器件结构的 Ge 基红外探测器, 包括金属-半导体-金属^[16]、p-i-n 型^[17]、锗/硅结^[18], 但高昂的制造成本或复杂的器件结构等因素阻碍了 Ge 基红外探测器的进一步发展。一种有效的解决方法是锗与二维 (Two Dimensionality, 2D) 材料结合组装形成肖特基结或异质结^[19-20]。在这些结中, 光吸收主要发生在窄带隙半导体中, 2D 材料通常作为一种活性层, 用于光载流子的分离和传输, 且具有较大透光率。这种器件结构简单, 易于制造。LUO L 等^[19]提出了一种由锗纳米锥 (Germanium Nanocones, GeNCs) 阵列和 PdSe_2 多层膜组成的敏感红外光电探测器, 该探测器通过直接硒化方法获得。组装好的 $\text{PdSe}_2/\text{GeNCs}$ 杂化异质结在 1550 nm 的光照下表现出显著的光伏性能, 使得近红外光探测器成为一种不需要外部电源的自驱动器件。

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进一步的器件分析表明,基于PdSe₂/GeNCs的近红外光探测器在1 550 nm光照下响应度达530.2 mA/W。WANG L等^[20]在研究中,将多层PtSe₂薄膜以简单的方式转移到Ge晶圆上形成垂直化异质结光电探测器,在1 550 nm辐照下,近红外光电探测器的响应率为602 mA/W。但是,PdSe₂等二维材料的合成需要价格昂贵的贵金属,且材料厚度一般较大,光的吸收率在15%左右^[21],会降低衬底锗的吸收从而影响异质结的响应。此时,石墨烯又成为了更优质的候选2D材料。

本文报道了一种基于石墨烯/锗肖特基结近红外光探测器,可应用于通讯波段,该探测器通过简单的化学气相沉积法在铜箔上沉积出高质量高性能石墨烯,转移到锗材料上形成结接触,通过光电性能测试与能带分析验证其性能。

1 制备

1.1 石墨烯的合成与转移

工艺中采用化学气相沉积法(Cheical Vapor Deposition, CVD)^[22],以25 μm厚的铜箔为催化剂,气体为CH₄(40 sccm)和H₂(20 sccm),在1 000°C下制备了单层石墨烯薄膜。生长后,将表面有单层石墨烯膜的铜箔上表面在3 000 r/min转速下旋转涂覆实验室自制的浓度为5%的聚甲基丙烯酸甲酯(Polymethyl Methacrylate, PMMA),70°C退火7 min,然后将铜箔置于CuSO₄溶液中。试剂溶液比例为CuSO₄:HCl:H₂O=10 g:50 mL:50 mL,通过HCl刻蚀铜箔,稳定剂为CuSO₄溶液。刻蚀结束后,将石墨烯膜附于载玻片上,转移至去离子水中反复清洗,直至去除残余的刻蚀液。

1.2 器件的制备

器件制备采用化学气象沉积法、湿法刻蚀法等,用到的仪器有紫外光刻机、等离子体清洗机、电子束蒸发系统。如图1所示,具体流程步骤为

1)对Ge衬底进行清洗,将覆盖300 nm厚SiO₂绝缘层的Ge衬底分别用丙酮、无水乙醇和去离子水超声清洗15 min。

2)在Ge衬底表面旋涂正光刻胶,在600 r/min转速下匀胶10 s,然后提速至3 000 r/min匀胶30 s,前烘5 min使光刻胶凝固,接着用紫外光刻机在覆盖绝缘层的Ge衬底上制备窗口(0.2 cm×0.2 cm),曝光显影后,进行后烘7 min稳定光刻胶,然后利用缓冲氧化物刻蚀液(Buffered Oxide Etch, BOE)刻蚀,刻蚀速率与刻蚀液浓度有关,刻蚀结束后用丙酮、无水乙醇、去离子水清洗干净。

3)转移刻蚀好的石墨烯薄膜至Ge衬底上形成异质结,室温下空置一夜,最后将器件放入丙酮中反复浸泡去除表面PMMA。

4)进行二次光刻,刻蚀器件的上电极,使用尺寸为0.25~0.3 cm的环形掩模版,之后将器件贴在电子束

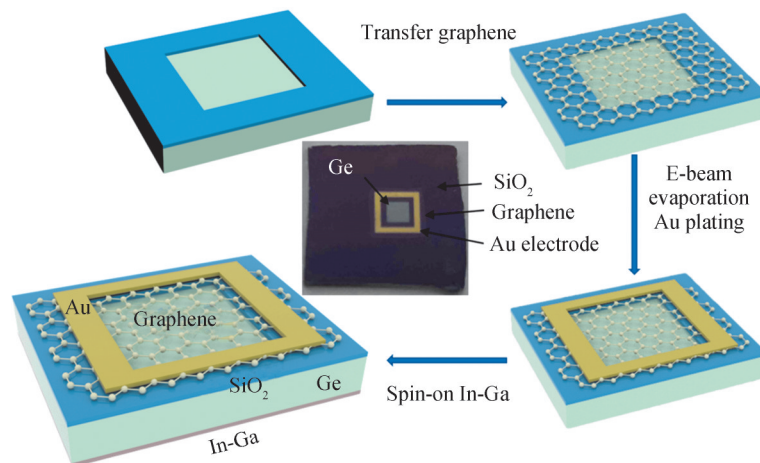


图1 制备锗/石墨烯肖特基结光探测器的步骤示意图和器件实物

Fig. 1 Schematic diagram of steps for constructing germanium/graphene Schottky junction light detector and photograph of the detector

蒸镀样品托盘上,预沉积 50 nm 的 Au 薄膜。

5)在器件底部涂抹均匀的 In-Ga 合金作为底部电极。

通过光刻、刻蚀、转移石墨烯、电子束蒸镀最终得到器件,从上到下依次是金电极、石墨烯、二氧化硅、锗、In-Ga 合金。器件实物如图 1 中心插图所示。

2 实验与分析

石墨烯/锗光电探测器的电学性能通过半导体表征系统(Keithley2400SP 2150,普林斯顿公司)进行测试。使用的测试光源为 1 550 nm 近红外激光二极管(M15500LP1,索雷博公司)。光谱响应在单色仪(LE-SP-M300)上进行。采用不同波长(265、365、450、530、660、730、810、970 和 1 050、1 200、1 300、1 400、1 500、1 550、1 600 nm)的激光二极管作为光源,研究光响应。所有光源的功率强度都经过功率计(Thorlabs GmbH, PM100D)仔细校准。所有测试均在室温、相对湿度 40%~60% 环境下进行。

图 2 描述了该光电探测器的工作机理。首先,对 Ge 的吸收进行仿真。不同波长的入射光以 90° 的入射角被引入锗表面的正面,采用 Synopsys Sentaurus TCAD 仿真 n-Ge 衬底(Sb 掺杂,载流子浓度为 10^{15} cm^{-3})、尺寸分别为 $20 \mu\text{m} \times 200 \mu\text{m}$ 的二维结构中的光子吸收速率,如图 2(a)所示。可以看出光子吸收速率的独特分布,当入射光的波长较短(如 265~365 nm)时,入射光的穿透深度(光子吸收速率值较高)很浅(小于 10 nm),表明光子几乎被吸收在异质结的表面,由于表面缺陷和/或悬空键的存在,该区域存在严重的载流子重组,从而降低了光响应。但随着入射光波长的增加,穿透深度将逐渐增加,在 1 600 nm 处达到最强吸收。通过文献中锗的折射率(n)、消光系数(k)^[19]计算出锗的吸收率与波长之间的关系,如图 2(b)所示。可以看

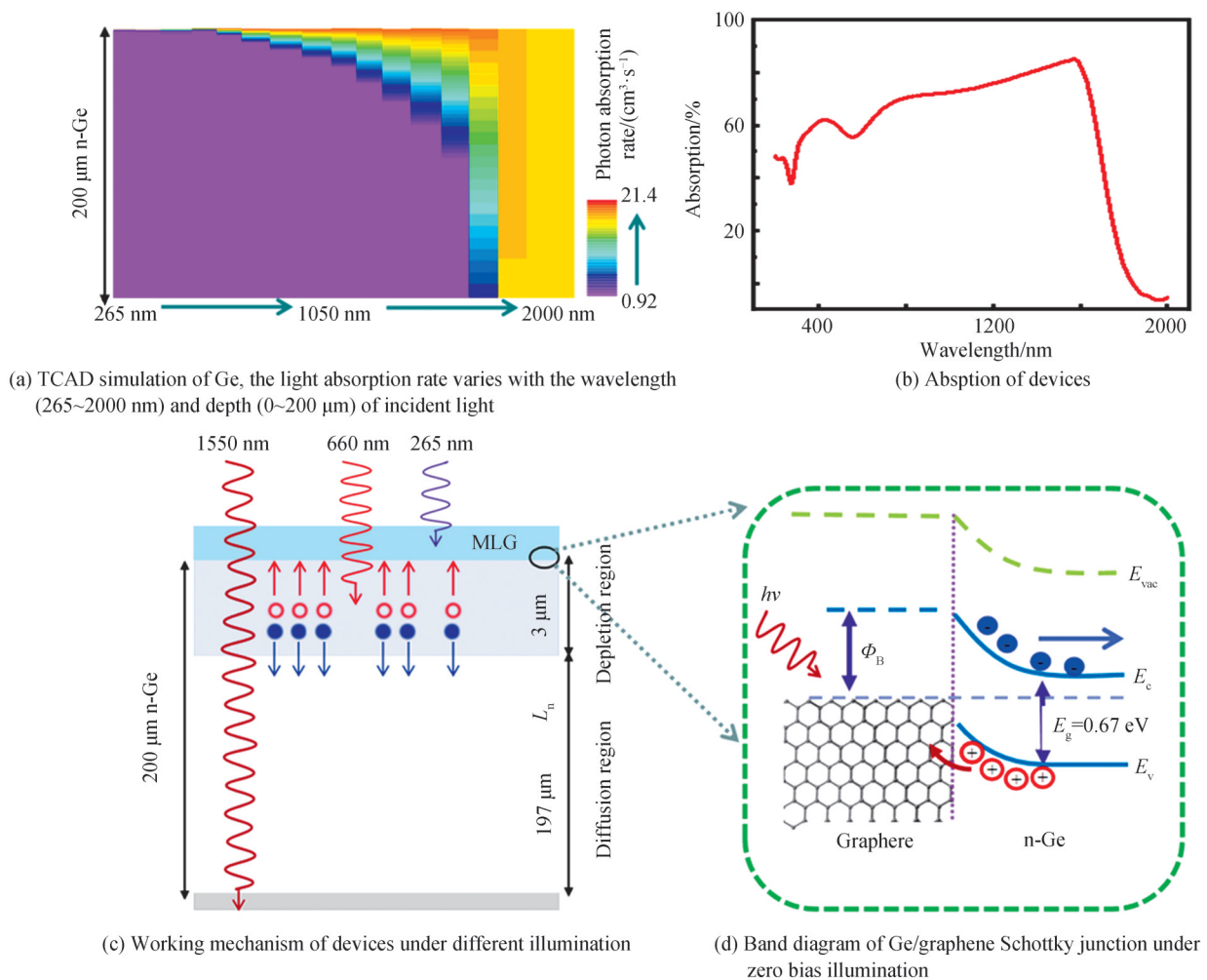


图 2 器件仿真与工作原理分析
Fig. 2 Device simulation and working principle analysis

出随着波长的增加,吸收先增加后减小,在1 550 nm处达到峰值。图2(c)显示了在短、中、长波长(如265、660、1 550 nm)照射下的电子空穴对。在光照条件下,探测器的结可分为耗尽区和扩散区,这两个区域都有助于光生载流子的扩散。根据计算,器件的耗尽区宽度在3 μm 左右,而与n型Ge(迁移率为1 800 cm^2/Vs)相对应的扩散长度(L_h)估计为700 μm 。当电子空穴对在该区域产生时,这两个区域都有助于光电流的形成。然而,随着入射波长从265 nm增加到1 550 nm,电子空穴对的产生将逐渐从单层石墨烯/n-Ge结的耗尽区扩展至扩散区,由于锗的厚度与扩散长度相比较小,大多数光诱导电子空穴对将被内置电场隔开,因此造成该光电探测器在不同波段光响应不同。当器件被短波长光照明时,相对较高的吸收系数会在异质结表面附近产生较强的吸收。相反,随着波长的增加,相应的相对较低的吸收系数将导致更深的穿透深度。图2(d)分析了器件零偏压下的能带。探测器的工作机理可以利用结的能带图和载流子输运过程来分析解释。石墨烯的费米能级(E_F)为4.7 eV,具有1~10 Ωcm^{-1} 电阻率的n-Ge的功函数大致为4.37 eV,在这两种材料形成肖特基结后,由于功函数的差异,电子从锗扩散到石墨烯,与此同时,空穴在锗的耗尽区形成。这种电荷转移打破了原来各自的能带平衡,锗表面附近能级向上弯曲,并出现内置电场。当受到能量超过锗(0.67 eV)禁带宽度的光照射时,锗的耗尽区会产生电子-空穴对。耗尽区附近产生的载流子会扩散到耗尽区。随后,电子和空穴在内置电场的作用下迅速分离,电子在底部的In-Ga电极被收集,而空穴则通过石墨烯转移并最终被上表面的金电极收集。上述过程在零偏置下产生了光电流。而在正向偏压下,结区宽度减小,且随着正向偏压逐步增加,正向电流将迅速增大;在反向偏压下,锗能带的弯曲更加显著,结区宽度增加,内建电场增强,反向电流变大。

图3(a)描绘的是器件在无光照条件下的 I - V 曲线,采用线性和半对数尺度对比,在 ± 1 V电压下,整流特性明显,器件的整流比约为 5.3×10^2 ,优于之前许多Ge基的异质结构,例如,PtSe₂/Ge(± 1.5 V时约430)^[23]、石墨烯/PtSe₂/Ge(在 ± 3 V时约25)^[20]、PdSe₂/Ge纳米锥阵列(在 ± 5 V时约5)异质结构^[19]。这里Au/石墨烯/Au和In-Ga/Ge/In-Ga都是欧姆接触,因此整流特性归因于石墨烯/锗肖特基结。图3(b)展示了该肖特基结光电探测器的稳定性,在无任何保护的条件下存储5个月后,用功率为0.3 mW/cm^2 的1 550 nm的光照射,仍能保持其优异的光开关特性;器件的开关比仍能达到 10^2 ,相比于存储之前,器件的光电流几乎没有衰减。图3(c)和(d)为器件在1 550 nm近红外光照射下,在不同光照强度下的 I - V 曲线和时间响应曲线。可以看出,光电压和光电流均与入射光功率呈正相关关系,分析其机理是因为随着光功率的增加,有更多的电子-空穴对产生且分离后形成更大的光电流。进一步地,从图3(d)中不同光功率下的光电流和暗电流计算出净光电流,绘制在图3(e)中。利用Origin软件进行线性拟合,可以用常用的幂律来描述,即

$$I_{\text{ph}} = aP^\theta \quad (1)$$

式中, I_{ph} 表示净光电流, a 表示与光波长有关的常数, P 表示入射光强, θ 是反映光载流子复合活性的经验值。最终得出净光电流的值与光强呈近似线性的关系, θ 的值为0.947。从图3(c)中可以得出光电压的值,绘制出光电压与光照强度的变化,如图3(f)。光电压首先在低光强区域(0.055~0.638 mW/cm^2)出现急剧增加,然后在高光强区域(0.638~1.4 mW/cm^2)出现缓慢增加。

计算了肖特基结光电探测器的光电流响应度 R_1 、噪声等效功率NEP和探测率 D^* 三个光电探测器关键性能指标。其中, R_1 可以解释为光电探测器在有源区输出光电流与入射光功率的比值;NEP又称为最小可测功率,定义为光电器件输出信号电压的有效值等于噪声均方根电压值时的入射光功率; D^* 表示光电探测器对微弱信号的探测能力。这三个参数计算公式分别为

$$R_1 = \frac{I_{\text{light}} - I_{\text{dark}}}{P_{\text{in}} A} \quad (2)$$

$$\text{NEP} = \frac{i_n^{1/2}}{R_1} \quad (3)$$

$$D^* = \frac{\sqrt{A} f}{\text{NEP}} \quad (4)$$

式中, I_{light} 表示光电流, I_{dark} 表示暗电流, A 表示光电探测器光敏区面积, P_{in} 表示入射光功率密度, i_n 表示噪声等效电流, f 表示频率带宽。

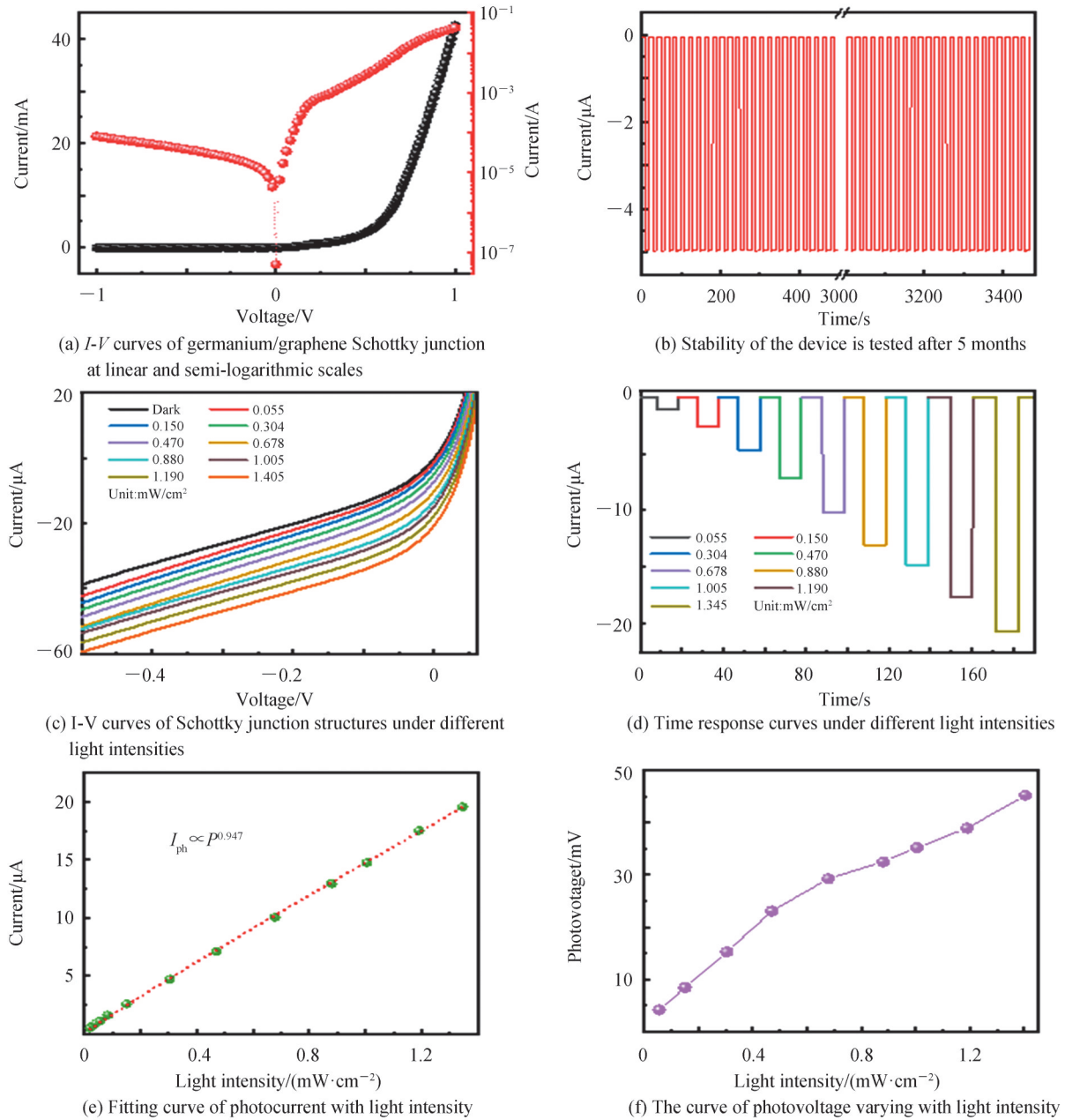


图3 锗/石墨烯肖特基结器件光电性能测试与分析

Fig. 3 Photoelectric performance test and analysis of Germanium/graphene Schottky junction devices

利用光谱响应在单色仪(LE-SP-M300)测量了同一功率不同波长下器件的光暗电流,带入式(2)计算得到归一化光谱响应,如图4(a)所示。器件的响应峰值在1 550 nm附近,与器件的吸收吻合。计算发现 R_1 均随光功率密度的增强而减小,如图4(b)所示。因此,依据实验数据计算,在零偏压,光照强度为0.055 mW/cm^2 的1 550 nm近红外光照射下,最大的 R_1 值为635.7 mA/W 。在1 550 nm波长下,相对较大的 R_1 值超过石墨烯/ PtSe_2/Ge (602 mA/W)^[20]、 PdSe_2/Ge 纳米锥阵列(530.2 mA/W)^[19]异质结构。为了验证该异质结构光电探测器记录快速变化光信号的能力,对其响应速度进行了研究。采用1 550 nm近红外激光二极管作为光源,由函数发生器驱动产生高频光信号,用数字示波器记录输出电压与时间的函数关系。如图4(c)所示,光电探测器显示了良好的开关特性,对不同频率(1,5,10 kHz)的脉冲光具有良好的可重复性,表明该器件在宽调制频率范围内具有良好的可操作性。此外,在10 kHz的光调制频率下,器件仍能保持光电压最大值的70%,如图4(d),表明该器件的3 dB带宽频率约为10 kHz(3 dB带宽定义为光响应下降到其峰值的

70.7%时所对应的光调制频率)。通过分析10 kHz时的单放大响应曲线(图4(c)),计算出上升和下降时间(t_r 和 t_f)分别为40 μ s和35 μ s,器件呈现出较快的响应速度,这与Graphene/Ge(t_r 和 t_f 分别为23 μ s和108 μ s)^[21]以及PdSe₂/Ge纳米锥阵列(t_r 和 t_f 分别为24.5 μ s和38.5 μ s)^[19]相当甚至更优。通过对器件在零偏压下的暗电流进行傅里叶变换可得出图4(e)的噪声等效频谱,如图4(e)插图所示,器件的暗电流约为 4.35×10^{-8} A,与石墨烯/PtSe₂/Ge(约 3.08×10^{-8} A)^[20]和PdSe₂/Ge纳米锥阵列(约 6.09×10^{-8} A)^[19]探测器的暗电流相当。提取1 Hz频率下 $i_n^{1/2}$ 的噪声值,约为 1.02×10^{-12} AHz^{-1/2},带入式(3)可以得出不同波长下的NEP大小,其中响应峰值波长1550 nm下NEP为 1.6×10^{-12} Hz^{-1/2}W⁻¹。对应的探测率 D^* 用式(4)计算,结果是 9.8×10^{10} Jones。

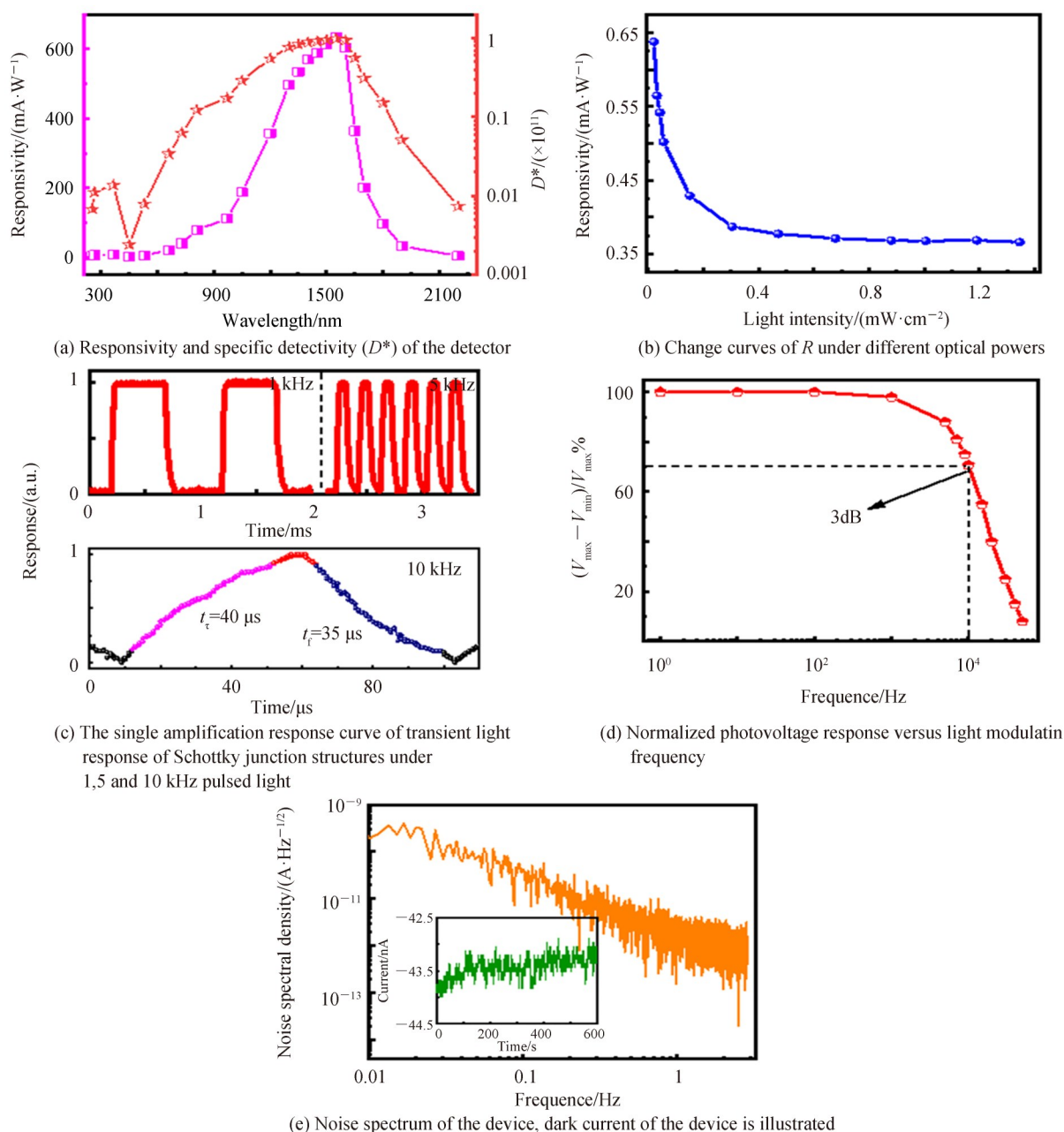


图4 锗/石墨烯肖特基结器件光电响应测试与器件噪声分析

Fig. 4 Photoelectric response test and devices noise analysis of Germanium/graphene Schottky junction devices

3 结论

本文制备了高性能近红外通讯波段的石墨烯/锗肖特基结光电探测器。仿真与实验结果表明,石墨烯

透明电极与锗衬底形成良好的肖特基接触,大大提升了器件的光生载流子收集效率。在无光照条件下,器件的整流比在 5.3×10^2 ,在光强为 0.3 mW/cm^2 的 1550 nm 近红外光的照射下,开关比达 10^2 ,光电流响应和探测率分别可达 635.7 mA/W 、 $9.8 \times 10^{10} \text{ Jones}$ 。器件响应速度较快,在 3 dB 带宽处上升和下降时间分别为 $40 \mu\text{s}$ 和 $35 \mu\text{s}$ 。此外,器件的制备工艺简单,在空气环境中5个月后,光电流几乎没有衰减,具有良好的再现性,展示了高性能石墨烯/锗光电探测器在近红外光电系统中具有潜在的应用前景。

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Graphene/germanium Schottky Junction Photodetectors in the Near-infrared Communication Band

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Abstract: Graphene has good electrical and optoelectronic properties, and the heterojunction photodetector formed by the combination of graphene and silicon semiconductor materials has excellent photodetection properties. The band gap of silicon is 1.12 eV, resulting in the cut-off wavelength of silicon-based near infrared ray photodetectors generally around 1.1 μm , and its detection wavelength range is relatively narrow, and Ge-based photodetectors can detect longer-wavelength near infrared ray light. In this study, a graphene/germanium Schottky junction photodetector in the near-infrared communication band was fabricated. First, high-quality graphene was prepared on the surface of copper foil by chemical vapor deposition method. Subsequently, the graphene was transferred to n-Ge surface by the wet transfer method, thus forming a Schottky contact. Finally, a high-performance graphene/germanium Schottky junction device is obtained by depositing gold electrode on the front side of the germanium substrate and spin-coating the In-Ga electrode on the back side. By simulating the photon absorption rate in the two-dimensional structure of the n-Ge substrate with Synopsys Sentaurus TCAD, the unique distribution of the photon absorption rate can be seen. When the wavelength of the incident light is short, the penetration depth of the incident light is very shallow (less than 10 nm), indicating that photons are almost absorbed in the surface of the heterojunction. Due to the existence of surface defects and/or dangling bonds, there is severe carrier recombination in this region, which reduces the photoresponse. But as the wavelength of the incident light increases, the penetration depth will gradually increase, reaching the strongest absorption at 1 600 nm. With the introduction of the graphene transparent electrode, it will form a good Schottky contact with the germanium substrate, which greatly improves the photo-generated carrier collection efficiency of the device. As the incident wavelength increases from 265 nm to 1 550 nm, the generation of electron-hole pairs will gradually expand from the depletion region of the monolayer graphene/n-Ge junction to the diffusion region, causing the photodetector has different responses with different wavelengths. Using the band diagram of the junction and the carrier transport process to analyze and explain the working mechanism of the detector, the Fermi level (E_F) of graphene is 4.7 eV. The resistivity of n-Ge is $1\sim 10\ \Omega\text{cm}^{-1}$, and its work function is about 4.37 eV. After the two materials form a Schottky junction, electrons diffuse from germanium to graphene due to the difference in work function, while holes are formed in the depletion region of germanium. This charge transfer breaks the original respective energy band balance, the energy levels near the germanium surface bend upward, and a built-in electric field appears. The depletion region of germanium produces electron-hole pairs when exposed to light with energies exceeding the forbidden band width of germanium (0.67 eV). Carriers generated near the depletion region diffuse into the depletion region. Subsequently, the electrons and holes are rapidly separated by the built-in electric field, the electrons are collected at the bottom In-Ga electrode, while the holes are transferred through the graphene and finally collected by the gold electrode at the upper surface. The same is true of the experimental results, the graphene/germanium Schottky junction shows obvious rectification characteristics, and the rectification ratio of the device is about 5.3×10^2 under the condition of $\pm 1\ \text{V}$ without illumination, which is better than many previous Ge-based heterogeneities structure. The detector shows good switching characteristics and good repeatability for pulsed light at different frequencies (1 kHz, 5 kHz, 10 kHz), indicating that the device has good operability over a wide modulation frequency range. After 5 months in the air environment, the photocurrent almost did not decay, and it still maintained its excellent optical switching characteristics, showing good stability. The device also exhibits obvious photoelectric response performance. Under the irradiation of 1 550 nm near-infrared light with a light intensity of $0.3\ \text{mW}/\text{cm}^2$, the responsivity and detection rate of the device can reach $635.7\ \text{mA}/\text{W}$ and 9.8×10^{10} Jones, respectively. At the same time, the device has a fast response speed, with rise and fall times of 40 μs and 35 μs , respectively, at a 3 dB

bandwidth. The potential applications of high-performance graphene/germanium photodetectors in near-infrared optoelectronic systems are demonstrated.

Key words: Photodetector; Graphene; Energy band analysis; Germanium; Schottky junction

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