doi:10.3788/gzxb20174611.1122004

折/衍混合大视场消热差红外双波段光学系统设计

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摘 要:根据双波段消热差理论设计了大视场消热差红外双波段光学系统.系统为4片式反远结构,包括一个衍射面和一个非球面,设计波长为 $3.5\sim4.8 \mu m/8\sim12 \mu m$,焦距为8 mm,全视场 80.2° , F/\ddagger 为2,系统出瞳与冷光阑严格匹配,满足100%冷光阑效率.根据消热差条件和波段间消色差条件,得出4片分离薄透镜光焦度分配的解,进一步建立了三维投影消热差图,据此合理选择光学材料.根据环境温度要求,利用光学被动式消热差的方法实现了系统在 $-40\sim60^\circ$ 的温度变化范围内的消热差设计.结果表明,系统在环境温度变化范围内成像质量良好,实现了光学系统无热化.

关键词:光学设计;双波段;消热差;折/衍混合;大视场

中图分类号:TN216 文献标识码:A

文章编号:1004-4213(2017)11-1122004-9

Designfor Cooled Dual-band Infrared Refractive-Diffractive Hybrid Optical System of Athermalization and Wide FOV

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Abstract: On the basis of the dual-band athermal theory, a cooled dual-band infrared refractive-diffractive hybrid optical system of athermalized and wide field of view was designed. The system consists of 4 lenses and introduces only 1 diffractive surface and 1 aspherical surface, the full field of view is 80.2° , the focal length is 8 mm, F number is 2, the operating wavelength range is $3.5 \sim 4.8 \mu m$ and $8 \sim 12 \mu m$, and the cold stop efficiency is 100%. The solutions of 4 separate thin lenses were obtained by analyzing dual-band achromatic and athermalized theory. Further, establishing the material model of athermalized dual-band optic system and choosing the optical materials properly. According to requirements, the dual-band athermal optical system with temperature of -40° C to 60° C was designed with the use of optical passive. The reasults prove that the optical system obtains a good image quality during -40° C to 60° C, and realizes athermalization.

Key words: Optical design; Dual-band; Athermalization; Refractive/diffractive hybrid; Wide field of view OCIS Codes: 040.3060; 080.2740; 080.3620; 080.6755; 090.1970

0 引言

大气中同一个目标物体在不同波段有不同的辐射特性,其中短波红外与可见光较为相似.在湿热的环境

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基金项目:国家自然科学基金(No.61520106015)资助

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中,波长为 3~5µm 的中波红外(Mid-Wavelength Infrared, MWIR)的侦查能力较为优越,主要用于观测高 温事件;波长为 8~12µm 的长波红外(Long-Wavelength Infrared, LWIR)在存在杂散辐射或靠近热源的情 况下具有较强的侦查能力,主要用于探测常温物体轮廓^[1-2].多波段光学系统的优势是具有全天候观察的能 力,分辨率高、采集光信息量大,可以应用于雾霾、沙暴等诸多复杂环境条件中,在许多自然灾害后的救援工 作中发挥着重要作用.

随着双色红外探测器的发展,双波段红外光学系统的研究也在不断深入,能在同一个光学系统中观察不同类型目标、不同环境条件、不同识别任务^[3-5].目前,对红外中波长波段同时敏感的探测器阵列主要有双波 段碲镉汞焦平面阵列和量子阱红外光电探测器阵列^[6].双波段共焦光学系统能同时采集两种波长,在同一像 面位置得到稳定的像质.温度变化会引起像差和热离焦,需要对设计进行无热化处理.

自 20 世纪 70 年代法国研制出双波段红外光学系统 VAMPIR 至今^[7],运用各种方法设计的双波段红外 消热差光学系统已经取得了丰富的成果.张欣婷等在光学系统中引入双层衍射元件,利用衍射元件负色散的 特性和光热膨胀系数小的特性,同时消色差和消热差^[7];常军等采用离轴反射结构设计光学无热化系统^[8]; 贾永丹等在双视场/双色红外消热差光学系统设计中使用投影消热差图合理选择光学材料^[9];杨新军等通过 计算消热差消色差方程,合理分配光焦度^[10],付强等提出了计算多个波长下平均离焦的方法^[11]实现红外双 波段无热化.

本文研究了光学被动式消热差条件、波段间消色差条件,建立了新三维消热差图(TPC 图).设计了 4 片 折/衍混合结构的大视场(Field of View,FOV)双波段共焦消热差光学系统,仅使用了一面非球面和一面衍 射面.系统采用了制冷型双波段探测器,实现了 100%冷光阑效率,满足消除背景杂光的要求.系统总长小于 84 mm,光阑距离像面 20 mm,后截距大于 12 mm,在-40~60℃实现了消热差,即温度导致的像面移动在 焦深范围内.

1 设计思路及理论分析

1.1 红外双波段波段间消色差原理

分别在中波长波波段应用物像关系公式有

$$\frac{1}{l_{1c}} - \frac{1}{l_{1c}} = (n_{1c} - 1)(\frac{1}{R_1} - \frac{1}{R_2}) = \varphi_{1c}$$
(1)

$$\frac{1}{l_{2c}} - \frac{1}{l_{2c}} = (n_{2c} - 1)(\frac{1}{R_1} - \frac{1}{R_2}) = \varphi_{2c}$$
(2)

式中, l_{1c} 、 l_{1c} 、 n_{1c} 、 φ_{1c} 分别表示长波波段中心波长的物距、像距、折射率、光焦度, l_{2c} 、 l_{2c} 、 n_{2c} 、 φ_{2c} 分别表示中波波段中心波长的物距、像距、折射率、光焦度, R_1 、 R_2 为元件曲率半径.

对式(1)、(2)作一阶近似得

$$\frac{l_{2c}^{'}-l_{1c}^{'}}{l^{'^{2}}}-\frac{l_{2c}^{'}-l_{1c}}{l^{2}}=(n_{1c}^{'}-n_{2c}^{'})(\frac{1}{R_{1}}-\frac{1}{R_{2}})=\frac{(n_{1c}^{'}-n_{2c}^{'})}{(n_{1c}^{'}-1)}\varphi_{1c}$$
(3)

类比阿贝数定义,取 $\frac{(n_{1c}-n_{2c})}{(n_{1c}-1)}=P$ 为波段间色差系数,波段间色差为 $L'_{pk}=l'_{2c}-l'_{1c}$,则^[12]

$$L_{pk}' = \frac{1}{u^2} \sum y_k^2 P_k \varphi_k = \frac{1}{(h_1 \varphi)^2} \sum y_k^2 P_k \varphi_k$$
(4)

式中,L'_{pk}为系统迭加波段间色差,y_k表示第k 面光线高度.

1.2 红外双波段消热差原理

薄透镜光焦度公式为

$$\varphi = (n - n_0)(c_1 - c_2) \tag{5}$$

微分可得

$$\frac{\mathrm{d}\varphi}{\mathrm{d}T} = \left(\frac{\mathrm{d}n}{\mathrm{d}T} - \frac{\mathrm{d}n_0}{\mathrm{d}T}\right) \left(c_1 - c_2\right) + \left(n - n_0\right) \left(\frac{\mathrm{d}c_1}{\mathrm{d}T} - \frac{\mathrm{d}c_2}{\mathrm{d}T}\right) \tag{6}$$

式中,T 表示温度,n。为介质折射率,n 为薄透镜折射率,c1、c2为透镜前表面和后表面的曲率.

由光学材料热膨胀系数为

$$\alpha_{g} = c_{1}^{-1} \frac{\mathrm{d}c_{1}}{\mathrm{d}T} = c_{2}^{-1} \frac{\mathrm{d}c_{2}}{\mathrm{d}T}$$
(7)

可得透射薄透镜的光热离焦系数为[13]

$$T_{f,r} = \alpha_{g} - \frac{1}{n - n_{0}} \left(\frac{\mathrm{d}n}{\mathrm{d}T} - \frac{\mathrm{d}n_{0}}{\mathrm{d}T} \right)$$
(8)

从式(8)可以看出,折射元件的光热离焦系数只由玻璃材料性质决定.

衍射光学元件光热离焦系数为[14]

$$T_{f,D} = \frac{1}{f} \frac{\mathrm{d}f}{\mathrm{d}T} = 2\alpha_{g} + \frac{1}{n_{o}} \frac{\mathrm{d}n_{o}}{\mathrm{d}T}$$

$$\tag{9}$$

式中,f为光焦度 φ 对应的焦距,从式(9)可以看出,衍射元件的光热离焦系数只与玻璃材料的光热膨胀系数 a_g 有关.由表1中中波红外材料的折射率温度变化梯度 $\frac{dn}{dT}$,和长波红外材料的折射率温度变化梯度 $\frac{dn}{dT}$ 可知薄 透镜的光热离焦系数小于零,衍射光学元件光热离焦系数大于零,机械材料的热膨胀系数大于零.故本系统 采用光学被动消热差的办法,即透射元件的光热离焦量、衍射元件的光热离焦的综合离焦与机械材料的热膨 胀离焦量相互补偿,使光学系统的像面离焦量在焦深范围内.双波段红外光学系统无热化,要满足五个方程:

系统光焦度方程为

$$\sum_{i=1}^{J} h_i \phi_i = \phi \tag{10}$$

中波波段内消色差为

$$\left(\frac{1}{h_{\perp}\phi}\right)^2 \sum_{i=1}^{j} h_i^2 C_{\mathrm{m}i} \phi_i = \Delta f_{\perp \mathrm{b}}$$

$$\tag{11}$$

长波波段内消色差为

$$\left(\frac{1}{h_{1}\phi}\right)^{2}\sum_{i=1}^{j}h_{i}^{2}C_{1i}\phi_{i} = \Delta f_{2b}$$
(12)

消热差条件为

$$\alpha_{\rm H}L = \left(\frac{1}{h_{\rm 1}\phi}\right)^2 \sum_{i=1}^{j} h_i^2 T_i \phi_i$$
(13)

波段间消色差为

$$L'_{p} = \frac{1}{(h_{1}\phi)^{2}} \sum h_{i}^{2} P_{i}\phi_{i}$$
(14)

其中 *q* 为系统的总光焦度,*h*_i 为第一近轴光线在第 *i* 个透镜上的入射高度,*q*_i 每个透镜的光焦度,*C*_{mi}是第 *i* 个透镜材料在中波红外波段的色散因子,*C*_{li}是第 *i* 个透镜材料在长波红外波段的色散因子,*T*_i 是第 *i* 个透镜的光热离焦系数.

1.3 红外材料特性及选取

表1列出中波红外、长波红外常用材料,及其对应的应用波段范围、光学性质、热学性质[15-16].

	表 1	光学材料的中波红外特性和长波红外特性	
Table 1	Material prop	erties of mid-wavelength infrared and long-wavelength infrar	ed

Character/ trademark	Refractive index at 4 μm	Refractive index at 10 μm	Abbe number 10 μm (8-12 μm)	Abbe number 4 μm (3-5 μm)	Inter band chromatic aberration/ $(\times 10^{-3})$	Thermal dispersive power/ (×10 ⁻⁶ /K)	dn/dt at $4\mu m/$ (×10 ⁻⁶ /K)	$\frac{dn/dt}{at \ 10 \mu m/} (\times 10^{-6}/K)$
Ge	4.024 3	4.032	783.394 3	116.315	6.836	12.61	396	396
ZnSe	2.433 2	2.406 4	57.468 92	177.988	18.646	37.7	60	61
ZnS	2.252 5	2.199 9	22.756 17	109.635	41.965	30.2	38.7	44
AMTIR-1	2.514 1	2.497 6	113.116 5	172.337 8	11.197 451	36.1	79	70.5
IG4	2.620 9	2.608 2	172.436	194.922 6	9.431 56	20.4	92	91
IG6	2.794 5	2.777 7	160.118	168.091 1	3.108 059	16	76	76

分析密接两组元系统,有

$$\begin{cases} \varphi_{1} + \varphi_{2} = \varphi \\ \frac{C_{m1}\varphi_{1} + C_{m2}\varphi_{2}}{\phi} = 0 \\ \frac{C_{11}\varphi_{1} + C_{12}\varphi_{2}}{\phi} = 0 \\ \frac{T_{1}\varphi_{1} + T_{2}\varphi_{2}}{\phi} = 0 \\ P_{1}\varphi_{1} + P_{2}\varphi_{2} = 0 \end{cases}$$
(15)

其解为

$$\begin{cases} \varphi_{1} = \frac{C_{m2}\varphi}{C_{m2} - C_{m1}} = -\frac{C_{12}\varphi}{C_{12} - C_{11}} = \frac{T_{2}\varphi}{T_{2} - T_{1}} = \frac{P_{2}\varphi}{P_{2} - P_{1}} \\ \varphi_{2} = -\frac{C_{m1}\varphi}{C_{m2} - C_{m1}} = -\frac{C_{11}\varphi}{C_{12} - C_{11}} = -\frac{T_{1}\varphi}{T_{2} - T_{1}} = \frac{-P_{2}\varphi}{P_{2} - P_{1}} \end{cases}$$
(16)

方程有解的条件是 $\frac{C_{m2}}{C_{m1}} = \frac{C_{12}}{C_{11}} = \frac{T_2}{T_1} = \frac{P_2}{P_1}$.

分析密接四组元系统,有

$$\begin{cases} \varphi_{1} + \varphi_{2} + \varphi_{3} + \varphi_{4} = \phi \\ \frac{C_{m1}\varphi_{1} + C_{m2}\varphi_{2} + C_{m3}\varphi_{3} + C_{m4}\varphi_{4}}{\phi} = 0 \\ \frac{C_{11}\varphi_{1} + C_{12}\varphi_{2} + C_{13}\varphi_{3} + C_{14}\varphi_{4}}{\phi} = 0 \\ \frac{T_{1}\varphi_{1} + T_{2}\varphi_{2} + T_{3}\varphi_{3} + T_{4}\varphi_{4}}{\phi} = 0 \\ P_{1}\varphi_{1} + P_{2}\varphi_{2} + P_{3}\varphi_{3} + P_{4}\varphi_{4} = 0 \end{cases}$$
(17)

其解为

$$\begin{cases} \varphi_{1} = \frac{P_{2}(C_{m3}T_{4} - C_{m4}T_{3}) + P_{3}(T_{2}C_{m4} - C_{m2}T_{4}) + P_{4}(C_{m2}T_{3} - T_{2}C_{m3})}{N_{1}} = \\ \frac{P_{2}(C_{13}T_{4} - C_{14}T_{3}) + P_{3}(T_{2}C_{14} - C_{12}T_{4}) + P_{4}(C_{12}T_{3} - T_{2}C_{13})}{N_{2}} \\ \varphi_{2} = \frac{P_{1}(C_{m3}T_{4} - C_{m4}T_{3}) + P_{3}(T_{1}C_{m4} - C_{m1}T_{4}) + P_{4}(C_{m1}T_{3} - T_{1}C_{m3})}{N_{1}} = \\ \frac{P_{1}(C_{13}T_{4} - C_{14}T_{3}) + P_{3}(T_{1}C_{14} - C_{11}T_{4}) + P_{4}(C_{11}T_{3} - T_{1}C_{13})}{N_{2}} \\ \varphi_{3} = \frac{P_{1}(C_{m2}T_{4} - C_{m4}T_{2}) + P_{2}(T_{1}C_{m4} - C_{m2}T_{4}) + P_{4}(C_{m1}T_{2} - T_{1}C_{m2})}{N_{1}} = \\ \frac{P_{1}(C_{12}T_{4} - C_{14}T_{2}) + P_{2}(T_{1}C_{14} - C_{12}T_{4}) + P_{4}(C_{11}T_{2} - T_{1}C_{12})}{N_{2}} \\ \varphi_{4} = \frac{P_{1}(C_{m2}T_{3} - C_{m3}T_{2}) + P_{2}(T_{1}C_{m3} - C_{m2}T_{3}) + P_{3}(C_{m1}T_{2} - T_{1}C_{m2})}{N_{1}} = \\ \frac{P_{1}(C_{12}T_{3} - C_{13}T_{2}) + P_{2}(T_{1}C_{13} - C_{12}T_{3}) + P_{3}(C_{11}T_{2} - T_{1}C_{12})}{N_{2}} \end{cases}$$

其中,

$$\begin{cases} N_{1} = P_{1} (C_{m4} T_{3} - C_{m3} T_{4} + C_{m2} T_{4} - C_{m4} T_{2} + C_{m2} T_{3} - C_{m3} T_{2}) + P_{2} (C_{m1} T_{4} - C_{m4} T_{1} + C_{m1} T_{3} - C_{m3} T_{1} + C_{m1} T_{4} - C_{m4} T_{1} + C_{m1} T_{4} - C_{m4} T_{1} + C_{m1} T_{2} - C_{m2} T_{1} + C_{m2} T_{4} - C_{m4} T_{2}) + P_{4} (C_{m1} T_{2} - C_{m2} T_{1} + C_{m1} T_{3} - C_{m3} T_{1} + C_{m2} T_{3} - C_{m3} T_{2}) \end{cases}$$

$$\begin{cases} N_{2} = P_{1}(C_{14}T_{3} - C_{13}T_{4} + C_{12}T_{4} - C_{14}T_{2} + C_{12}T_{3} - C_{13}T_{2}) + P_{2}(C_{11}T_{4} - C_{14}T_{1} + C_{11}T_{3} - C_{13}T_{1} + C_{12}T_{4} - C_{14}T_{1} + C_{11}T_{2} - C_{12}T_{1} + C_{12}T_{4} - C_{14}T_{2}) + (19) \\ P_{4}(C_{11}T_{2} - C_{12}T_{1} + C_{11}T_{3} - C_{13}T_{1} + C_{12}T_{3} - C_{13}T_{2}) \\ \vec{x}(10) \sim (14) \oplus, \mathfrak{P} h_{i}\varphi_{i} = \varphi_{i}^{'}, h_{i}T_{i} = T_{i}^{'}, h_{i}C_{i} = C_{i}^{'}, \beta \# \beta \\ \frac{\varphi_{1}^{'} + \varphi_{2}^{'} + \varphi_{3}^{'} + \varphi_{4}^{'} = \phi}{\varphi} \\ \frac{C_{11}^{'}\varphi_{1}^{'} + C_{12}^{'}\varphi_{2}^{'} + C_{13}^{'}\varphi_{3}^{'} + C_{14}^{'}\varphi_{4}^{'}}{\phi} = 0 \\ \frac{C_{11}^{'}\varphi_{1}^{'} + T_{2}^{'}\varphi_{2}^{'} + T_{3}^{'}\varphi_{3}^{'} + T_{4}^{'}\varphi_{4}^{'}}{\phi} = 0 \end{cases}$$

$$(20)$$

其解为

$$\begin{cases} \varphi_{1}^{'} = \frac{P_{2}^{'}(C_{m3}^{'}T_{4}^{'} - C_{m4}^{'}T_{3}^{'}) + P_{3}^{'}(T_{2}^{'}C_{m4}^{'} - C_{m2}^{'}T_{4}^{'}) + P_{4}^{'}(C_{m2}^{'}T_{3}^{'} - T_{2}^{'}C_{m3}^{'})}{N_{1}} = \\ \frac{P_{2}^{'}(C_{13}^{'}T_{4}^{'} - C_{14}^{'}T_{3}^{'}) + P_{3}^{'}(T_{2}^{'}C_{14}^{'} - C_{12}^{'}T_{4}^{'}) + P_{4}^{'}(C_{12}^{'}T_{3}^{'} - T_{2}^{'}C_{13}^{'})}{N_{2}} \\ \varphi_{2}^{'} = \frac{P_{1}^{'}(C_{m3}^{'}T_{4}^{'} - C_{m4}^{'}T_{3}^{'}) + P_{3}^{'}(T_{1}^{'}C_{m4}^{'} - C_{m1}^{'}T_{4}^{'}) + P_{4}^{'}(C_{m1}^{'}T_{3}^{'} - T_{1}^{'}C_{m3}^{'})}{N_{1}} \\ = \frac{P_{1}^{'}(C_{13}^{'}T_{4}^{'} - C_{14}^{'}T_{3}^{'}) + P_{3}^{'}(T_{1}^{'}C_{m4}^{'} - C_{m2}^{'}T_{4}^{'}) + P_{4}^{'}(C_{m1}^{'}T_{3}^{'} - T_{1}^{'}C_{m3}^{'})}{N_{2}} \\ \varphi_{3}^{'} = \frac{P_{1}^{'}(C_{m2}^{'}T_{4}^{'} - C_{m4}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{m4}^{'} - C_{m2}^{'}T_{4}^{'}) + P_{4}^{'}(C_{m1}^{'}T_{2}^{'} - T_{1}^{'}C_{m2}^{'})}{N_{1}} \\ \varphi_{3}^{'} = \frac{P_{1}^{'}(C_{12}^{'}T_{4}^{'} - C_{14}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{m3}^{'} - C_{m2}^{'}T_{3}^{'}) + P_{4}^{'}(C_{11}^{'}T_{2}^{'} - T_{1}^{'}C_{m2}^{'})}{N_{1}} \\ \varphi_{4}^{'} = \frac{P_{1}^{'}(C_{m2}^{'}T_{3}^{'} - C_{m3}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{m3}^{'} - C_{m2}^{'}T_{3}^{'}) + P_{3}^{'}(C_{m1}^{'}T_{2}^{'} - T_{1}^{'}C_{m2}^{'})}{N_{1}} \\ \frac{P_{1}^{'}(C_{12}^{'}T_{3}^{'} - C_{m3}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{m3}^{'} - C_{m2}^{'}T_{3}^{'}) + P_{3}^{'}(C_{m1}^{'}T_{2}^{'} - T_{1}^{'}C_{m2}^{'})}{N_{1}} \\ = \frac{P_{1}^{'}(C_{12}^{'}T_{3}^{'} - C_{13}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{13}^{'} - C_{12}^{'}T_{3}^{'}) + P_{3}^{'}(C_{11}^{'}T_{2}^{'} - T_{1}^{'}C_{m2}^{'})}{N_{1}} \\ = \frac{P_{1}^{'}(C_{12}^{'}T_{3}^{'} - C_{13}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{13}^{'} - C_{12}^{'}T_{3}^{'}) + P_{3}^{'}(C_{11}^{'}T_{2}^{'} - T_{1}^{'}C_{m2}^{'})}{N_{1}} \\ = \frac{P_{1}^{'}(C_{12}^{'}T_{3}^{'} - C_{13}^{'}T_{2}^{'}) + P_{2}^{'}(T_{1}^{'}C_{13}^{'} - C_{12}^{'}T_{3}^{'}) + P_{3}^{'}(C_{11}^{'}T_{2}^{'} - T_{1}^{'}C_{12}^{'})}{N_{1}}} \\ = \frac{P_{1}^{'}(C_{12}^{'}T_{3}^{'} - C_{13}^{'}$$

其中

$$\begin{cases} N_{1} = P'_{1} (C'_{m4} T'_{3} - C'_{m3} T'_{4} + C'_{m2} T'_{4} - C'_{m4} T'_{2} + C'_{m2} T'_{3} - C'_{m3} T'_{2}) + P'_{2} (C'_{m1} T'_{4} - C'_{m4} T'_{1} + C'_{m1} T'_{3} - C'_{m3} T'_{1} + C'_{m1} T'_{4} - C'_{m4} T'_{1} + C'_{m1} T'_{3} - C'_{m3} T'_{1} + C'_{m1} T'_{2} - C'_{m2} T'_{1} + C'_{m2} T'_{4} - C'_{m4} T'_{2}) + P'_{4} (C'_{m1} T'_{2} - C'_{m2} T'_{1} + C'_{m1} T'_{3} - C'_{m3} T'_{1} + C'_{m2} T'_{3} - C'_{m3} T'_{2}) \\ N_{2} = P'_{1} (C'_{14} T'_{3} - C'_{13} T'_{4} + C'_{12} T'_{4} - C'_{14} T'_{2} + C'_{12} T'_{3} - C'_{13} T'_{2}) + P'_{2} (C'_{11} T'_{4} - C'_{14} T'_{1} + C'_{11} T'_{3} - C'_{13} T'_{1} + C'_{12} T'_{4} - C'_{14} T'_{2} + C'_{12} T'_{3} - C'_{13} T'_{2}) + P'_{2} (C'_{11} T'_{4} - C'_{14} T'_{1} + C'_{11} T'_{3} - C'_{13} T'_{1} + C'_{12} T'_{4} - C'_{14} T'_{3}) + P'_{3} (C'_{11} T'_{4} - C'_{14} T'_{1} + C'_{11} T'_{2} - C'_{12} T'_{1} + C'_{12} T'_{4} - C'_{14} T'_{2}) + P'_{4} (C'_{11} T'_{2} - C'_{12} T'_{1} + C'_{11} T'_{3} - C'_{13} T'_{1} + C'_{12} T'_{3} - C'_{13} T'_{1} + C'_{12} T'_{3} - C'_{13} T'_{2}) \\ \end{pmatrix}$$

在初始解中为减小系统的高级像差,应尽量减小各组元的光焦度.由式(18)和(21)可以看出,分母越大, 各组元的光焦度越小,即 N_1 、 N_2 越小,系统各组元的光焦度越小.根据分析建立三维坐标系,X 轴为C,Y 轴 为T,Z 轴为P.在同一坐标系中标记出各个材料在 3~5 μ m 波段内各性能指标的散点图(\triangle)和 8~12 μ m 波段内各性能指标的散点图(O).根据表 1 在图 1 中标出各个材料,绘制出其中三种材料组成的新型消热差 TPC 图,如图 1.以长波下的三种材料所组成的三角形为基准平面,将中波下的三角形投影到基准平面上.其 重合面积越大,系统初始结构时分配到各个元件的光焦度越小,对像质越有利.



图 1 双波段消热差 TPC 图 Fig.1 TPC diagram of dual-band athermalization

2 设计实例

表 2 列出了制冷型双波段消热差光学系统和大视场长波红外消热差光学系统的设计方法与指标.基于 --款像面尺寸为 384×288,像元尺寸为 25 μm 的制冷型双波段红外探测器,设计参数见表 3.系统共采 2 片 GE,1 片 ZNSE,1 片 IG4 材料,系统 *F*/♯为 2,将光阑后置于光学系统后,探测器窗口前,出瞳大小和冷屏大 小配合,满足 100%的冷光阑效率.根据新三维投影热差图选择光学材料,合理分配光焦度,使用光学仿真软 件 ZEMAX 优化光学系统,最终得到光学系统结构如图 2.系统为反远结构,主面后移,后截距大于 12 mm.第 一面的相对孔径大,轴外光线高度高.第一片透镜为负光焦度大折射率的 GE 材料,有利于大视场下光线的 偏折,GE 材料的波段内色散和波段间色散都很小,不产生大的像差.第二片为正光焦度透镜,在第二片透镜 的后表面引入一面二元衍射面平衡色差,负色散的衍射元件有利于消除色差,减小镜片的曲率和所承担的光 焦度.第三片透镜为负光焦度的 ZnSe.第四片透镜的材料为 GE 材料,在四片透镜的前表面引入一面非球面, 校正大视场引起的轴外剩余像差,以及剩余色差.

Reference	Wavelength/ μm	FOV/ (°)	F/#	Effective focal length/mm	Detector	The total number of lens(PCS)	Total length/ mm	Major technique
No.1	3.5-4.8/ 9.2-10.8	74	2.5	9.5	384×288/ 25μm	3	202	Three mirrorunobscured system
No.2	3.5~4.8/ 8~12	68	2	10	384×288/ 25μm	5	113	Refractive system with 3 aspheric surfaces
No.3	$4.4 \sim 5.7$ 7.8 ~ 8.8	96	4	6	384×288/ 25μm	6	62	Double-layer harmonic diffraction element
No.4	3-5	82	2	8	384×288/ 25μm	3	65	Refractive-Diffractive Hybrid System
This design	3.5~4.8/ 8~12	80.2	2	8	384×288/ 25μm	4	84	Refractive-Diffractive Hybrid System

表 2 制冷型定焦红外消热差系统的指标和主要设计方法 Table 2 Design index and method of cooled infrared optical system of athermalized

表 3 系统设计参数					
Table 3 System design parameters					
Parameters	Value				
Wavelength/ μ m	3.5-4.8 8-12				
FOV/(°)	80.2				
Effective focal length/mm	8				
F/ $#$	2				
Cold shield efficiency	100%				
Operating temperature/°C	$-40\!\sim\!60$				





3 设计评价

系统长波星点在衍射极限内,中波星点接近衍射极限,如图 3.中波焦深 $\Delta z = \pm 2\lambda (F/\#) = 64 \mu m$,长波 焦深为 $\Delta z = \pm 2\lambda (F/\#) = 160 \mu m$,系统最大离焦量为 62.892 μm ,小于两个波段的焦深,能够实现共焦.系统 MTF 如图 4,在 $-40 \sim 60^{\circ}$ 的温度范围内变化极小,长波 MTF 在 20 lp/mm 大于 0.6,中波 MTF 在 20 lp/mm大于 0.34,成像质量良好.





图 4 制冷型红外双波段消热差系统 MTF 图 Fig.4 MTF of athermalized dual-band infrared optical system

4 结论

本文分析了红外双波段消热差和消色差条件,结合新三维消热差图合理选择材料,构建了红外双波段消 热差光学系统,实现了在-40~60℃光学系统无热化设计.运用新型三维消热差图选择材料更加直观,且利 于材料重新选择分配.该方法需要逐一对比不同的材料组合,由于红外材料种类少,方法优点明显,可广泛应 用于各红外波段的光学设计中.但是,在可见光的双波段消热差设计中,因为可选择的材料多,用新型三维消 热差图选择材料效率低.可以采用计算机编写材料分析程序,对大量光学材料作选择对比,结合理论进一步 扩展该方法在可见光领域的应用.

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Foundation item: National Natural Science Foundation of China(No. 61520106015)