

标量衍射理论的非傍轴近似及其有效性*

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摘 要 当光束束腰(或衍射孔孔径)可与波长相比拟或光束具有较大的发散角时,傍轴近似不再成立. 在标量瑞利-索末菲衍射积分的基础上,进一步研究了衍射场的非傍轴近似解,并详细分析了解的有效性. 以平面波圆孔衍射为例,对衍射场的精确解、非傍轴近似解以及菲涅耳近似解进行了详细的数值计算和比较研究. 结果表明,非傍轴近似对微小孔衍射非常精确、有效.

关键词 物理光学;标量衍射理论;非傍轴近似;圆孔衍射;有效性

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0 引言

对于大多数光学衍射与光束传输的实际问题,傍轴标量衍射理论都是非常精确、有效的. 然而,对强聚焦光束或二极管激光器发出的光束,当束腰宽度为波长量级、发散角很大时,傍轴近似不再成立,甚至标量近似都已不再成立^[1~7]. 近年来国内外关于非傍轴衍射的研究工作十分活跃,文献[5,6]给出了在自由空间传播的非傍轴光束的级数解,级数解的适用范围与光束的束腰宽度、发散角、传输距离以及所使用的级数解的阶次有关^[5,8]. 文献[9~11]给出了平面孔衍射场的非傍轴近似解,其有效性问题还需要作进一步研究. 本文在标量瑞利-索末菲衍射积分的基础上,进一步研究了衍射场的非傍轴近似解,详细分析了解的有效性. 以平面波圆孔衍射为例,对衍射场的精确解、非傍轴近似解以及菲涅耳近似解进行了详细的数值计算和比较研究.

1 衍射场的非傍轴近似解

光波入射到无限大不透明带孔 Σ 的平面衍射屏上时, $z > 0$ 空间的标量衍射场,一般采用第一类瑞利-索末菲衍射积分表示^[10,12]

$$E(x, y, z) = \frac{k}{2\pi i z} \iint_{\Sigma} E(x_1, y_1) \frac{\exp(ikR)}{R} \cdot (1 + \frac{i}{kR}) \frac{z}{R} dx_1 dy_1 \quad (1)$$

式中 $k = 2\pi/\lambda$, $R = [(x-x_1)^2 + (y-y_1)^2 + z^2]^{1/2}$ 为源点 $(x_1, y_1, 0)$ 与场点 (x, y, z) 之间的距离. 式(1)满足 $z=0$ 处衍射场的边界条件,通常被认为是描述整个衍射空间($z > 0$)标量衍射场的精确解. 虽然由

式(1)可以精确计算非傍轴标量衍射场,然而由于其数学与数值计算上的复杂性,在实际应用中极为不便,需要寻找有效的近似公式. 为了对式(1)作非傍轴近似,在式(1)的指数因子中将 R 近似表示为

$$R \approx r + \frac{x_1^2 + y_1^2 - 2xx_1 - 2yy_1}{2r} \quad (2)$$

式中 $r = \sqrt{x^2 + y^2 + z^2}$, 其它项中取 $R \approx r$, 并假设 $kr \gg 1$

则式(1)近似表示为

$$E(x, y, z) \approx \frac{1}{i\lambda} \cdot \frac{z}{r^2} e^{ikr} \iint_{\Sigma} E(x_1, y_1) \cdot \exp\left[\frac{ik(x_1^2 + y_1^2)}{2r}\right] \exp\left[\frac{-ik(xx_1 + yy_1)}{r}\right] dx_1 dy_1 \quad (4)$$

式(4)便是标量衍射的非傍轴近似公式,与众所周知的菲涅耳衍射公式非常相似,称之为广义菲涅耳衍射,在傍轴区过渡到菲涅耳衍射. 非傍轴近似的条件可进一步表示为

$$r \gg \lambda \quad (5a)$$

$$r^2 \gg (x_1^2 + y_1^2 - 2xx_1 - 2yy_1)_{\max} \quad (5b)$$

$$\frac{2\pi}{\lambda} \cdot \frac{(x_1^2 + y_1^2 - 2xx_1 - 2yy_1)_{\max}^2}{8r^3} < \frac{\pi}{16} \quad (5c)$$

式中不等式(5c)右边选取 $\pi/16$, 能基本满足式(4)近似成立的要求,但根据准确度要求,仍具有一定的任意性,以上近似条件为非傍轴近似的充分条件. 显然,衍射孔孔径较小时(与光波波长 λ 相比),非傍轴近似的适用范围由式(5a)、(5b)、(5c)共同决定,孔径较大时由不等式(5c)确定.

2 圆孔的非傍轴衍射及其有效性

设 $x_1 o y_1$ 面上,有一无限大不透明平面屏,屏上开有一半径为 ρ_0 的小圆孔,圆心在坐标原点, $z > 0$ 为衍射区. 设入射光为垂直照射线偏振的单位振幅单色平面波,假设

$$E(x_1, y_1, 0) = \begin{cases} 1 & x_1^2 + y_1^2 \leq \rho_0^2 \\ 0 & x_1^2 + y_1^2 > \rho_0^2 \end{cases} \quad (6)$$

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将式(6)代入式(4),可求得圆孔衍射的非傍轴近似解

$$E_a(x, y, z) \approx \frac{kz}{ir^2} e^{ikz} \int_0^{\rho_1} \exp\left(\frac{ik\rho_1^2}{2r}\right) J_0(k\rho\rho_1/r) \rho_1 d\rho_1 \quad (7)$$

式中 $\rho = \sqrt{x^2 + y^2}$, $\rho_1 = \sqrt{x_1^2 + y_1^2}$. 将式(7)应用于傍轴区便过渡到圆孔的非涅耳衍射^[13]

$$E_b(x, y, z) = \frac{2\pi}{iz\lambda} e^{ikz} \exp\left(\frac{ik\rho^2}{2z}\right) \cdot \int_0^{\rho_1} \exp\left(\frac{ik\rho_1^2}{2z}\right) J_0(k\rho\rho_1/z) \rho_1 d\rho_1 \quad (8)$$

运用与第一类瑞利-索末菲衍射积分等价的角谱衍射理论计算圆孔衍射的精确解^[14~16]

$$E_c(x, y, z) = \int_0^{\rho_0} \rho_0 J_1(\rho_0 k_\rho) J_0(\rho k_\rho) \cdot \exp(iz \sqrt{k^2 - k_\rho^2}) dk_\rho \quad (9)$$

将非傍轴近似条件式(5c)应用于圆孔衍射,可得圆孔衍射非傍轴近似解的适用范围为

$$r^3 > 4(\rho_1^2 - 2\rho_1\rho_0 \cos \theta)_{\max}^2 / \lambda = 4(\rho_0^2 + 2\rho_0\rho)^2 / \lambda \quad (10)$$

式中 θ 为矢量 $\boldsymbol{\rho}$ 与 $\boldsymbol{\rho}_1$ 之间的夹角. 在相同准确度要求下圆孔非涅耳衍射的适用范围可表示为^[17]

$$z^3 > 4(\rho_1^2 + \rho^2 - 2\rho_1\rho \cos \theta)_{\max}^2 / \lambda = 4(\rho_0 + \rho)^4 / \lambda \quad (11)$$

数值计算表明,由不等式(10)确定的圆孔衍射的非傍轴近似条件是充分的,但并非必要. 适当放宽式

(10)的要求,采用

$$r^3 > 4(\rho_0^2 + \rho_0\rho)^2 / \lambda \quad (12)$$

作为圆孔衍射的非傍轴近似条件,则更为准确. 取 $\rho_0 = 0.5\lambda, \lambda, 2\lambda$, 由式(11)、(12)计算圆孔衍射的非傍轴近似与非涅耳近似的适用范围,其范围为相应曲线与 I 轴构成的区间,如图 1. 图中 α, β, γ 为非傍轴近似; α', β', γ' 为非涅耳近似; $\alpha, \alpha' - \rho_0 = 0.5\lambda; \beta, \beta' - \rho_0 = \lambda; \gamma, \gamma' - \rho_0 = 2\lambda$. 由式(7)、(8)、(9)计算观察面上衍射场的强度分布,如图 2, 曲线 a, b, c 分别表示圆孔标量衍射场的非傍轴近似解 E_a 、非涅耳近似解 E_b 与精确解 E_c .

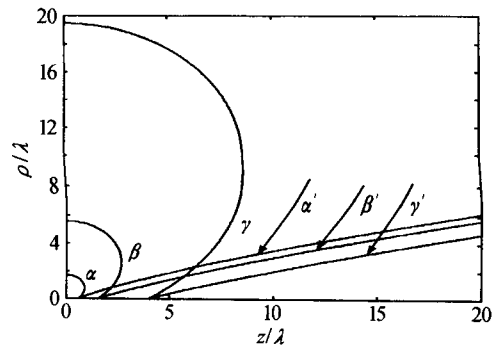


图 1 非傍轴近似与非涅耳近似适用范围比较
Fig. 1 Comparison of the applicable region between the non-paraxial approximation and fresnel approximation

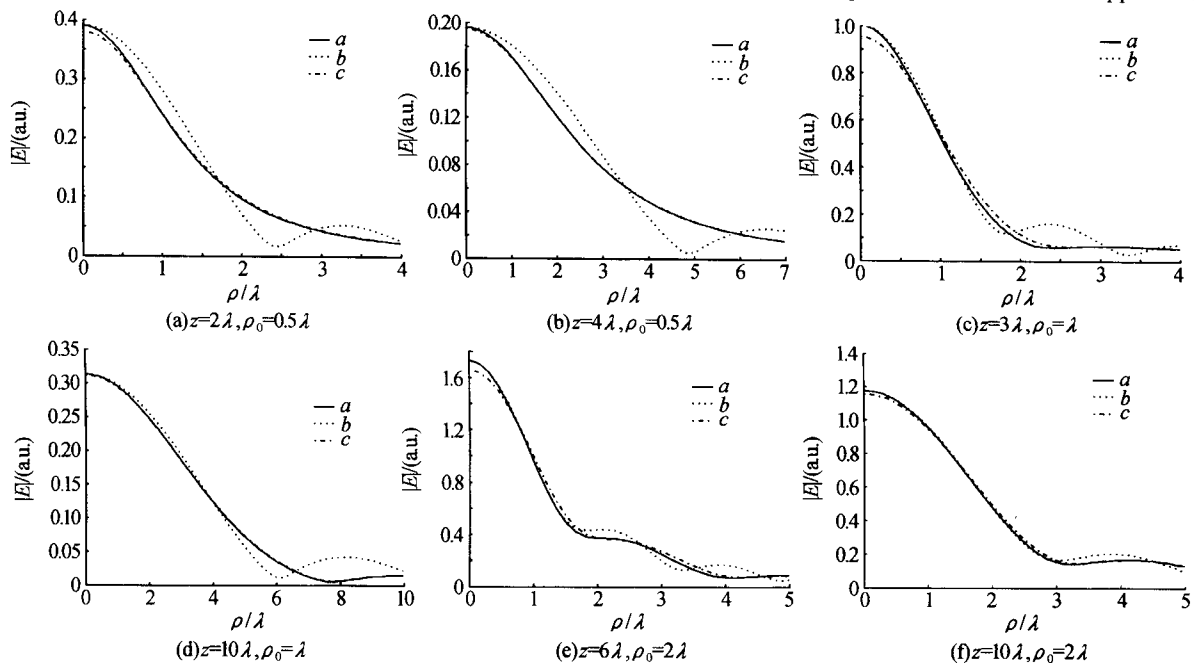


图 2 场振幅随 ρ/λ 的变化
Fig. 2 Field amplitudes versus ρ/λ

3 结论

计算结果表明,除衍射孔附近的较小区域外,非傍轴近似解(广义非涅耳衍射)与衍射场的精确积分解非常好地吻合,而且数学形式也比较简单,能节省

大量计算机时. 由式(12)确定的圆孔衍射非傍轴近似的适用范围,在近轴区,还应满足 $z \gg \lambda$ 和 $z \gg \rho_0$ 的条件(一般取 $z > 3\lambda$ 和 $z > 3\rho_0$,即可满足一定的准确度要求). 由图 1,2 可知,圆孔衍射非傍轴近似的适用范围要远远大于非涅耳近似,非涅耳近似不能

描述微小孔的非傍轴衍射. 最后必须强调指出, 研究光束的非傍轴传播, 还须考虑光场的矢量修正, 限于篇幅, 在此不作进一步讨论.

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Non-paraxial Approximation of Scalar Diffraction Theory and Its Validity

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Abstract It is well known that the paraxial approximation is no longer valid for the beams with large divergence angle and small spot size comparable with the wavelength. Thus, a rigorous non-paraxial treatment becomes necessary. In this paper, based on the scalar Rayleigh-Sommerfeld diffraction formula, the non-paraxial approximation solution and its validity are studied. Detailed numerical calculations for the exact solution, non-paraxial approximation solution and Fresnel approximation solution of circular aperture diffraction are performed and compared; It is shown that the non-paraxial approximation solution is rigorous and valid for the diffraction of a small aperture.

Keywords Physical optics; Scalar diffraction theory; Non-paraxial approximation; Circular aperture diffraction; Validity



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