

# 单片集成两段式双波长分布反馈半导体激光器

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**摘要** 双波长激光器腔内模式竞争激烈,因此输出模式的稳定性是双波长激光器的关键参数。从降低双波长激光器中两个主模之间的功率差、提高边模抑制比出发,设计了集成反射区的两段式双波长分布反馈半导体激光器。利用传输矩阵法对激光器的光栅结构进行仿真,分析了反射区光栅对激光器的阈值、主模功率差等参数的影响。根据仿真优化的结果,制作了单片集成两段式双波长分布反馈半导体激光器芯片并进行了测试。测试结果表明两段式结构能够提高双波长激光器的稳定性和边模抑制比,减小两个主模的功率差。在稳定工作的情况下,两个主模功率差可达0.3 dB,边模抑制比大于35 dB。

关键词 激光器;双波长激光器;分布反馈半导体激光器;两段式激光器;光子集成;单片集成
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# 1引言

双波长激光器在激光雷达、微波光子系统和光学 传感等领域应用广泛[17]。传统的双波长激光器通常 由两个独立的激光器组成[8-9],由于受外界温度和环境 振动等因素的影响,分立器件构成的双波长激光器的 稳定性较差,实用性不好。为了解决上述问题,人们研 制出了多种集成化的双波长分布反馈(DW-DFB)激 光器,如集成Y型波导DW-DFB激光器<sup>[10]</sup>、多段式 DW-DFB激光器和横向耦合型DW-DFB激光器<sup>[11-13]</sup>。 单片集成的DW-DFB激光器的优点是结构紧凑,性能 稳定。然而,这些单片集成的DW-DFB激光器通常需 要复杂的制造工艺和控制电路,导致其使用成本较高, 因此,实际使用中更需要一种结构简单、稳定性好的双 波长激光器。之前报道了一种基于等效啁啾双相移的 DW-DFB激光器,该激光器在啁啾光栅中引入双相移 结构,能够实现激光器两个模式在空间上的分离,减少 了模式竞争,能够实现稳定的双波长工作<sup>[14]</sup>。但是这 种结构的双波长激光器仍然存在问题:1)为了实现稳 定的双波长输出,激光器两输出端面需要镀抗反射膜, 激光器两端输出功率相同,造成功率浪费;2)此种结构 的双波长激光器在两个主模空间上实现了分离,可以 减少模式竞争,但是输出端两个模式的输出功率差较 大,影响双波长激光器的实用化。

本文报道了一种集成了反射光栅的两段式DW-DFB激光器(TS-DW-DFB)。激光器包含DW-DFB 激光器部分和反射光栅(GR)部分。两部分采用相同 的外延片结构,共用同一个波导。DW-DFB激光器部 分的光栅采用等效啁啾双相移取样光栅结构<sup>[14]</sup>,反射 光栅采用均匀取样光栅结构。两部分光栅均采用重构 等效啁啾(REC)技术设计和制作。通过设计两部分 光栅的取样周期,可以让DW-DFB激光器的两个主模 位于GR区光栅反射谱的中心,而DW-DFB激光器边 模位于GR区光栅反射谱的中心,而DW-DFB激光器的 边模抑制比(SMSR)。同时由于GR区光栅反射DW-DFB激光器主模获得更高的增益,从而提高输出端的 输出功率,利用光栅反射的调节能够有效减小激光器 两个模式的输出功率差,实现激光器稳定的双波长 激射。

利用上述结构设计并制作了单片集成 TS-DW-DFB激光器芯片。经实验测试可得,该激光器的主模 间波长差为0.62 nm,主模间功率差为0.3 dB,边模抑 制比大于 35 dB,测量了1 h内 TS-DW-DFB激光器的 光谱特性,发现激光器能够实现稳定的双波长输出。

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## 2 基本原理与仿真设计

REC 技术可以通过设计特殊结构的采样布拉格 光栅(SBG)等效实现各种复杂结构的光栅,获得不同 的反射和透射响应,下面对 REC 技术的原理进行介 绍<sup>[15-18]</sup>,SBG 的折射率可以表示为

$$\Delta n(z) = \frac{1}{2} \Delta n_0 s(z) \exp\left(\frac{j2\pi z}{\Lambda_0}\right) + \text{c.c.} \qquad (1)$$

式中: $\Delta n_0$ 为种子光栅折射率差; $\Lambda_0$ 是种子光栅周期; s(z)是采样调制函数。根据傅里叶变换,s(z)可以表达为

$$s(z) = \sum_{m} F_{m} \exp\left(j\frac{2m\pi z}{P}\right), \qquad (2)$$

式中:P是采样周期;F<sub>m</sub>为第m级信道的傅里叶展开系数,如果在z<sub>0</sub>处引入一段长度为ΔP的位移,则折射率为

$$\Delta n_m(z) = \begin{cases} F_m \exp\left(\frac{j2\pi z}{\Lambda_0} + \frac{j2\pi m z}{P}\right), & z \leq z_0 \\ F_m \exp\left(\frac{j2\pi z}{\Lambda_0} + \frac{j2\pi m z}{P} - j\theta_m\right), & z > z_0 \end{cases},$$

(3)

式中: $\theta_m$ 是等效相移, $\theta_m = 2m\pi\Delta P/P$ 。仿真时使用的 是+1级信道即 m=+1。为了实现等效π相移,需要 满足 $\Delta P/P = 0.5$ ,同时通过改变SBG的采样周期P和 占空比可以实现等效啁啾光栅和等效切趾光栅<sup>[19]</sup>。取 样光栅的+1级子光栅的光栅耦合系数 $\kappa_1$ 与占空比的 关系满足

$$\kappa_1 = \kappa_0 \frac{\sin(\gamma \pi)}{\pi}, \qquad (4)$$

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式中:γ是取样光栅占空比;κ<sub>0</sub>是种子光栅耦合系数。 可见,占空比γ=0.5时+1级子光栅的光栅耦合系数 最大,γ远离0.5时+1级子光栅的光栅耦合系数逐渐 减小,因此通过改变取样光栅占空比能够等效实现光 栅切趾。

TS-DW-DFB 激光器芯片如图 1(a) 所示, TS-DW-DFB激光器包含两部分,分别是DW-DFB激光 器区和GR区,其中DW-DFB激光器区和GR区长度 分别为L<sub>DW</sub>和L<sub>GR</sub>, DW-DFB激光器区和GR区的注入 电流分别为 IDW 和 IGR。 TS-DW-DFB 激光器的光栅结 构如图1(b)所示,DW-DFB激光器区和GR区的光栅 均采用REC技术设计的SBG,两部分光栅具有相同的 种子光栅周期。DW-DFB激光器区的光栅为线性啁 啾采样光栅,P"是采样光栅中的第n个采样周期,且在 腔内1/3和2/3处引入两个等效相移(EPS),两个等效 相移分别标记为EPS1和EPS2。GR区采用均匀采样 光栅,光栅的采样周期等于DW-DFB激光器区光栅中 心位置的采样周期。为了减小GR反射谱禁带外的边 模,GR区的采样光栅通过改变占空比实现等效切 趾<sup>[19]</sup>。这里光栅占空比是指一个采样周期内有光栅部 分的长度与采样周期的比值。当占空比沿光栅位置变 化时,可以引入不同的切趾函数曲线,从而在+1级子 光栅中实现等效切趾。如图1(c)所示,GR区的采样 光栅的占空比利用sinc函数进行切趾。本文中仿真采 用的光栅参数是:种子光栅周期 $\Lambda_0$ 为256.672 nm,有 效折射率是3.2,耦合系数是40 cm<sup>-1</sup>,GR区的采样周 期 $P_{GR}$ 是4251 nm, DW-DFB激光器区的中心采样周 期是 4251 nm, DW-DFB 激光器区的长度  $L_{DW}$  是 500 µm,光栅啁啾系数为50 nm/mm。



图 1 TS-DW-DFB激光器结构及各部分的光栅示意图。(a) TS-DW-DFB激光器结构图;(b) TS-DW-DFB激光器的光栅结构图; (c) GR区的光栅结构图

Fig. 1 Structure of TS-DW-DFB laser and partial gratings. (a) Structure of TS-DW-DFB laser; (b) grating structure of TS-DW-DFB laser; (c) grating structure in GR section

利用传输矩阵法<sup>[20-21]</sup>对GR区光栅的反射谱和 DW-DFB激光器区的光栅透射谱、时延谱进行仿真。 GR区无切趾时的反射谱和DW-DFB激光器区的透射 谱、时延谱分别如图 2(a) 及图 2(b) 所示。DW-DFB 激 光器区光栅的时延谱反映了激光器腔内不同模式获得 的增益大小,也就是说时延越大的模式能够获得的增 益越大,则越容易激射发光。由时延谱可以看出设计 的DW-DFB激光器存在两个可以激射的主模。DW-DFB激光器的两个主模(model和mode2)位于GR反 射区的反射禁带之内,激光器的边模位于反射禁带之 外,这有利于提高激光器的输出功率,提高SMSR。然 而,GR区光栅反射谱禁带两侧仍然存在较高的反射 边带,这些反射边带可为DW-DFB激光器的边模提供 额外的反射,从而影响激光器的SMSR。为了提高激 光器的SMSR,GR区反射光栅采用sinc函数实现等效 切趾<sup>[19]</sup>。经过切趾后的GR区的反射谱和DW-DFB

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激光器区的透射谱、时延谱分别如图 2(c) 及图 2(d) 所 示。经切趾后,GR区的反射谱禁带两侧反射边带消 失,有利于提高DW-DFB激光器的SMSR。双波长激 光器的波长间隔是非常重要的参数指标,TS-DW-DFB激光器可以通过改变DW-DFB激光器光栅中两 个等效相移(EPS1和EPS2)的大小来调节激光器的 波长间隔。图 3(a)是计算出的不同相移大小情况下 DW-DFB激光器光栅的时延谱,可以看出相移大小不 同,时延谱中两个峰的位置也不同,即两个主模的波长 间隔随着相移大小的变化而变化。图 3(b)是计算出 的不同相移下 TS-DW-DFB 激光器两个主模波长间 隔变化曲线。可以看出:当相移由 0.4π 增加到 1.8π 的过程中, TS-DW-DFB 激光器波长间隔先由 1.65 nm减小到 0.75 nm,之后逐渐增加到 1.61 nm。 由此可以看出通过设计 TS-DW-DFB 激光器中光栅 的等效相移大小就可以实现波长间隔的调谐。



图2 GR区和DW-DFB激光器的光栅谱图。(a)未切趾的GR区的反射谱和DW-DFB激光器的透射谱;(b)未切趾的GR区的反射 谱和DW-DFB激光器的时延谱;(c)切趾后的GR区的反射谱和DW-DFB激光器的透射谱;(d)切趾后的GR区的反射谱和 DW-DFB激光器的时延谱

Fig. 2 Spectra of gratings in GR section and DW-DFB laser. (a) Reflection spectrum in GR section without apodization and transmission spectrum of DW-DFB laser; (b) reflection spectrum in GR section without apodization and delay spectrum of DW-DFB laser; (c) reflection spectrum in GR section with apodization and transmission spectrum of DW-DFB laser; (d) reflection spectrum in GR section with apodization and delay spectrum of DW-DFB laser;

利用传输矩阵法,对具有不同光栅结构的 TS-DW-DFB激光器的输出特性进行分析<sup>[20-21]</sup>。首先分别 计算了 *I*<sub>DW</sub>=100 mA 时单段 DW-DFB激光器和 TS-DW-DFB激光器两个主模的腔内光子分布图,计算结 果如图 4(a)和图 4(b)所示。由图 4(a)可知,单段 DW-DFB激光器中两个主模的光子分布在空间上是分离 的,同一输出端面两个激光器的主模输出功率有明显 的差别,并且激光器左端面的输出总功率与右端面的 输出总功率相等。如图4(b)所示,当GR反射区的长 度为500 μm时,由于GR区光栅的反射作用,右端面 的输出功率远高于左端面的输出功率,从而提高了激 光器的输出效率。同时,由于 model在近GR区端面





Fig. 3 Delay spectra of DW-DFB section and wavelength spacing between two main-modes of TS-DW-DFB laser for different EPSs.
(a) Delay spectra of DW-DFB laser with different EPSs; (b) variation of wavelength spacing between two main-modes of TS-DW-DFB laser with different EPSs

的光子分布浓度更高,所以GR区反射光栅对 model的作用更加明显,model在激光器右端面的输出功率



图4 *I*<sub>DW</sub>=100 mA时单段式DW-DFB激光器和TS-DW-DFB激光器两个主模的功率分布。(a)单段式DW-DFB激光器两个主模的 功率分布;(b) TS-DW-DFB激光器两个主模的功率分布

Fig. 4 Power of two main modes of DW-DFB laser and TS-DW-DFB laser at  $I_{DW}$ =100 mA. (a) Power of two main modes of singlestage DW-DFB laser; (b) power of two main modes of TS-DW-DFB laser

接着分析 GR 区长度对 TS-DW-DFB 激光器的影 响,图5是计算得到的激光器两个主模的阈值增益和 主模功率差(PDM)与GR区长度的关系图。GR区的 长度 L<sub>GR</sub>从 0 增加到 1000 µm。L<sub>GR</sub>=0 是指 TS-DW-DFB激光器没有集成反射光栅部分,即普通的单段式 DW-DFB激光器。相比于单段DW-DFB激光器,TS-DW-DFB激光器的两个主模的阈值增益减小,并且两 个主模的阈值增益随着GR区长度的增加而减小。如 图1所示, EPS1相比EPS2更接近GR区, 因此GR区 对 mode1 的反射比 mode2 大,这也导致了 mode1 具有 更小的阈值增益。除此之外,GR区的反射光栅还能 够降低 model 和 mode2 两个主模之间的功率差, 如图 5(b) 所示,相比于单段 DW-DFB 激光器, TS-DW-DFB激光器的PDM随着GR区长度的增大明显减小, 也就是说两段式 DW-DFB 结构能够显著减小两个主 模的功率差。具体原因在于:没有GR区时,DW-DFB 激光器的两个主模 model 和 mode2 具有相同的阈值,

由于 model 和 mode2 在空间上是分离的,激光器的左 端面 model 的输出功率大于 mode2 的输出功率,而激 光器的右端面 model 的输出功率小于 mode2 的输出功 率[图 4(a)],因此激光器的 PDM 较大。当 DW-DFB 激光器的左端面集成了 GR 区后,激光器只能从右端 面输出,两个模式光子分布的不同导致 model 的输出 功率小于 mode2 的输出功率。但是,由于反射光栅的 作用,mode1 比 mode2 具有更低的阈值。相同电流下, mode1 能够获得比 mode2 更大的增益,从而缩小了 mode1 与 mode2 在右端面的输出功率差,有效降低激 光器的 PDM,且反射光栅越长其反射作用越强,激光 器的 PDM 越小,这与图 4(b)的仿真结果是相互吻 合的。

的提升幅度比mode2更大,因此激光器右侧输出端两

个主模的功率差明显减小。

# 3 实验结果与分析

基于上述仿真结果,设计和制作了TS-DW-DFB 激光器芯片,激光器芯片的实物图如图6所示。如图6



图 5 L<sub>DW</sub>=500 μm 时的阈值增益和输出功率差。(a)阈值增益随 GR 区长度的变化;(b)输出功率差随 GR 区长度的变化 Fig. 5 Threshold gain and output power difference for L<sub>DW</sub>=500 μm. (a) Threshold gain versus length of GR section; (b) PDM versus length of GR section

(a)所示,激光器的总腔长为1000 µm,GR 区和 DW-DFB 激光器区的长度均为500 µm,两部分之间设有电 隔离区,可保证两部分能够独立注入电流而不相互影 响,TS-DW-DFB 激光器两端镀抗反射膜。GR 反射区 和 DW-DFB 激光器采用相同的种子光栅,种子光栅的 周期为256.672 nm,GR 区的采样周期为4251 nm,GR 区光栅中心处的采样占空比为0.5,占空比向两侧逐 渐减小,采用 sinc 函数切趾。DW-DFB 激光器区的中 心采样周期是4251 nm,啁啾率为50 nm/mm,即第一 个采样周期和最后一个采样周期的差值为25 nm,两 个等效相移分别分布在腔内的 L/3 和 2L/3 处。图 6 (b)是 TS-DW-DFB 激光器芯片端面的电子扫描显微 镜(SEM)图片,TS-DW-DFB 激光器芯片采用的是脊 波导结构。

为了进行对比,还制作了单段式的DW-DFB激光器,单段式DW-DFB激光器的结构与TS-DW-DFB激光器中的发光区部分相同。图7(a)和图7(b)分别是





单段式 DW-DFB 激光器芯片在注入电流分别为 60 mA 和 90 mA 时的光谱图,从光谱图中可以发现 SMSR 均大于 30 dB,是稳定的双波长激射,然而激光 器两个主模 mode1(λ<sub>1</sub>)和 mode2(λ<sub>2</sub>)的功率差较大, mode1的输出功率小于 mode2,这与图 4(a)的仿真结 果相符。图 7(c)和图 7(d)是测得的 TS-DW-DFB 激 光器在激光器区注入电流 I<sub>bw</sub>分别为 60 mA 和 90 mA 时的光谱图,测量过程中反射区电流 I<sub>GR</sub>=12 mA。从 图中可以看出相同注入电流下,TS-DW-DFB 激光器 的 SMSR 高于单段式 DW-DFB 激光器,并且 TS-DW-DFB 激光器相比于单段式 DW-DFB 激光器具有更小 的 PDM。

TS-DW-DFB 激光器在保持 DW-DFB 激光器区 注入电流不变的情况下,可以通过调节GR区的电流 调节激光器两个主模的输出功率以及PDM和SMSR。 图 8(a) 是测得的 TS-DW-DFB 激光器两个主模输出 功率随着GR区电流的变化情况,测量过程中DW-DFB区的注入电流 Inw 设置在 80 mA 不变。GR 区电 流较小时,GR区光栅的反射作用较弱,受光子分布的 影响, mode1的输出功率小于 mode2的输出功率。随 着GR区注入电流的增加,光栅反射作用增强,model 输出功率增加, mode2输出功率减小, 当 $I_{CR}$ =12 mA时 两个模式输出功率相近。GR区注入电流继续增加, mode1的输出功率进一步增加,mode2的输出功率进 一步减小,因此通过改变GR区电流能够优化双波长 激光器的PDM。图8(b)描述了室温下TS-DW-DFB 激光器芯片在 I<sub>Dw</sub>=80 mA 时, GR 区注入电流由 0 增 加到20mA时的PDM和SMSR的变化情况。从图中 可以看出:在IGR=12 mA时激光器的SMSR最大为 36 dB, PDM 最小为 0.3 dB; 当 IGR 小于 12 mA 时反射 区吸收损耗大,反射较小,model获得的反射不够,此 时 mode2 的功率大于 mode1 的功率; 当 IGR 大于 12 mA 时,反射区的吸收损耗小,反射区对mode1的反射作用 逐渐增大, mode1的输出功率大于 mode2, PDM 开始 由小变大。



- 图7 不同注入电流下单段式 DW-DFB 激光器和 TS-DW-DFB 激光器的光谱图。(a) I<sub>Dw</sub>=60 mA 时单段式 DW-DFB 激光器的光谱 图;(b) I<sub>Dw</sub>=90 mA 时单段式 DW-DFB 激光器的光谱图;(c) I<sub>Dw</sub>=60 mA 和 I<sub>GR</sub>=12 mA 时 TS-DW-DFB 激光器的光谱图; (d) I<sub>Dw</sub>=90 mA 和 I<sub>GR</sub>=12 mA 时 TS-DW-DFB 激光器的光谱图
- Fig. 7 Optical spectra of single-section DW-DFB laser and TS-DW-DFB laser with different injection currents. (a) Optical spectrum of single-section DW-DFB laser with I<sub>DW</sub>=60 mA; (b) optical spectrum of single-section DW-DFB laser with I<sub>DW</sub>=90 mA; (c) optical spectrum of TS-DW-DFB laser with I<sub>DW</sub>=60 mA and I<sub>GR</sub>=12 mA; (d) optical spectrum of TS-DW-DFB laser with I<sub>DW</sub>=90 mA and I<sub>GR</sub>=12 mA



图 8 TS-DW-DFB激光器在不同反射区电流下的主模功率以及主模功率差和边模抑制比。(a) TS-DW-DFB激光器在 I<sub>Dw</sub>=80 mA 时不同反射区电流下两个主模的输出功率;(b) TS-DW-DFB激光器在 I<sub>Dw</sub>=80 mA 时不同反射区电流下的主模功率差和边模抑制比

Fig. 8 Output power, PDM, and SMSR of TS-DW-DFB laser versus current of GR section. (a) Output power of TS-DW-DFB laser versus current of GR section when I<sub>DW</sub>=80 mA; (b) PDM and SMSR of TS-DW-DFB laser versus current of GR section when I<sub>DW</sub>=80 mA

图 9 是测得的 DW-DFB 和 TS-DW-DFB 激光器的 SMSR 和 PDM 随着注入电流 *I*<sub>DW</sub>的变化情况,测试过程中,TS-DW-DFB 激光器的反射区电流保持为 *I*<sub>GR</sub>=12 mA。从图 9 (a)中可以看出,相同的注入电流下,TS-DW-DFB 激光器比单段 DW-DFB 激光器具有更高的 SMSR,注入电流 *I*<sub>DW</sub>大于 100 mA 后,单段 DW-DFB激光器出现多模而TS-DW-DFB仍然可以 保持稳定的双模输出。由图9(b)可知,随着注入电流 *I*<sub>Dw</sub>的增加,单段DW-DFB激光器和TS-DW-DFB激 光器的PDM都有所增大,但单段DW-DFB激光器增 大得更快,相同电流下单段DW-DFB激光器的PDM 明显高于TS-DW-DFB激光器,TS-DW-DFB激光器

的 PDM 在 *I*<sub>DW</sub> 增大过程中始终小于 1 dB。由实验结果可以看出,集成了反射光栅的 DW-DFB 激光器能够有







图 9 DW-DFB激光器和TS-DW DFB激光器的SMSR和PDM。(a)不同 I<sub>DW</sub>下DW-DFB激光器和TS-DW-DFB激光器的SMSR; (b)不同 I<sub>DW</sub>下DW-DFB激光器和TS-DW-DFB激光器的PDM

Fig. 9 SMSR and PDM of DW-DFB laser and TS-DW-DFB laser. (a) SMSR of DW-DFB laser and TS-DW-DFB laser versus I<sub>DW</sub>;
 (b) PDM of DW-DFB laser and TS-DW-DFB laser versus I<sub>DW</sub>

稳定性是激光器的重要指标,测试了TS-DW-DFB激光器在注入电流不变且激光器工作温度变化 时的光谱特性,以及保持激光器注入电流和工作温度 不变的条件下1h内的光谱特性。图10(a)是TS-DW-DFB 激光器在注入电流  $I_{DW}$ =80 mA、 $I_{GR}$ =12 mA 时, 不同温度下激光器的二维光谱图,当激光器工作温度 由 16 ℃增加到 40 ℃时, TS-DW-DFB 激光器波长向长 波长方向漂移,温度变化过程中激光器光谱基本保持 稳定。图 10(b)是 TS-DW-DFB 激光器在注入电流 *I*<sub>DW</sub>=80 mA和*I*<sub>GR</sub>=12 mA,工作温度保持24℃不变时 1h内激光器的二维光谱图,可见1h内激光器的输出 光谱基本保持稳定。图 11(a)~(c)是测得的 TS-DW-DFB 激光器在  $I_{DW}$ =80 mA、 $I_{GR}$ =12 mA 时 PDM、 SMSR、两个主模的波长以及波长间隔随温度的变化 情况,在激光器的工作温度由16℃增加到40℃过程 中,激光器的 SMSR 保持在 35 dB 以上, PDM 在 1.2 dB以下,两个主模的波长随着温度的增加同时向

(a) 40

Temperature /°C

35

30

25

20

15

长波长方向移动, model 波长的温度调谐系数为 0.086 nm/℃, mode2的温度调谐系数为0.088 nm/℃。 因此两个主模的波长间隔随着温度的升高略有增加。 其原因在于TS-DW-DFB激光器腔内的光子分布不 是均匀的,如图4(b)所示,在决定mode2波长的相移 位置 EPS2 处的光子密度高于相移位置 EPS1 处的光 子密度,导致EPS2处的温度更高,当温控系统升高相 同的温度时EPS2处的温度升高比EPS1处更快,从而 导致 mode2 波长随温度的调谐系数略大于 mode1。图 12(a)和图 12(b) 是测得的 TS-DW-DFB 激光器在 *I*<sub>DW</sub>=80 mA、*I*<sub>GR</sub>=12 mA、工作温度控制在 24 ℃的条 件下,1h内的PDM、SMSR以及波长间隔变化情况。 在激光器工作条件不变的情况,TS-DW-DFB激光器 的PDM、SMSR以及波长间隔基本保持稳定。由于在 实际测试过程中,激光器芯片并没有封装,测试系统受 外界环境的影响,所以在长时间内测试结果存在一定 的波动。



图 10 不同工作条件下测得的 TS-DW-DFB 激光器的二维光谱图。(a) *I*<sub>DW</sub>=80 mA、*I*<sub>GR</sub>=12 mA 时,不同工作温度下 TS-DW-DFB 激光器的二维光谱图;(b) *I*<sub>DW</sub>=80 mA、*I*<sub>GR</sub>=12 mA、工作温度为 24 ℃时,1 h内 TS-DW-DFB 激光器的二维光谱图

Fig. 10 2D optical spectra of TS-DW-DFB laser under different operation conditions. (a) 2D optical spectra of TS-DW-DFB laser under different temperatures when  $I_{\rm DW}$ =80 mA and  $I_{\rm GR}$ =12 mA; (b) 2D optical spectra of TS-DW-DFB laser during 1 h when  $I_{\rm DW}$ =80 mA,  $I_{\rm GR}$ =12 mA, and operation temperature is 24 °C





Fig. 11 PDM, SMSR, wavelength of two main modes, and wavelength spacing of TS-DW-DFB laser varying with temperature.
(a) SMSR and PDM of TS-DW-DFB laser varying with temperature when I<sub>DW</sub>=80 mA and I<sub>GR</sub>=12 mA; (b) wavelength of main modes varying with temperature when I<sub>DW</sub>=80 mA and I<sub>GR</sub>=12 mA; (c) wavelength spacing between two main modes varying with temperature when I<sub>DW</sub>=80 mA and I<sub>GR</sub>=12 mA;



图 12 TS-DW-DFB激光器的 PDM、SMSR 和波长间隔随时间的变化情况。(a) I<sub>Dw</sub>=80 mA、I<sub>GR</sub>=12 mA、工作温度为 24 ℃时,1 h 内 TS-DW-DFB 激光器的 PDM 和 SMSR 的变化情况;(b) I<sub>Dw</sub>=80 mA、I<sub>GR</sub>=12 mA、工作温度为 24 ℃时,1 h 内 TS-DW-DFB激光器的波长间隔变化情况

Fig. 12 PDM, SMSR, and wavelength spacing of TS-DW-DFB laser varying with time. (a) SMSR and PDM of TS-DW-DFB laser during 1 h when I<sub>DW</sub>=80 mA, I<sub>GR</sub>=12 mA, and operation temperature is 24 °C; (b) wavelength spacing between main modes of TS-DW-DFB laser during 1 h when I<sub>DW</sub>=80 mA, I<sub>GR</sub>=12 mA, and operation temperature is 24 °C;

### 4 结 论

模式稳定性对单片集成的双波长激光器芯片非常 重要,为了提高双波长激光器芯片的输出稳定性,提出 了一种单片集成TS-DW-DFB激光器。对单片集成 TS-DW-DFB激光器芯片进行了理论仿真与实验研 究。通过理论仿真分析了反射区光栅对TS-DW-DFB 激光器芯片输出稳定性的影响。根据理论分析结果, 研制了TS-DW-DFB激光器芯片并进行了实验测试。 测试结果表明TS-DW-DFB激光器芯片并进行了实验测试。 测试结果表明TS-DW-DFB激光器芯片并进行了实验测试。 利试结果表明TS-DW-DFB激光器芯片并进行了实验测试。 初读长输出,并且能够明显改善激光器两个主模之间 的功率差,提高激光器的边模抑制比。所提激光器芯 片为实现稳定的双波长激光光源提供了新的思路。

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# Monolithic Integrated Two-Section Dual-Wavelength Distributed Feedback Semiconductor Laser

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#### Abstract

**Objective** Dual-wavelength lasers are widely used in lidar, microwave photonic systems, and optical sensing systems. Traditional dual-wavelength lasers usually consist of two discrete lasers. Due to the influence of external temperature and environmental vibration, the operation stability of traditional dual-wavelength lasers is poor, which limits their application. To solve the above problems, multiple integrated dual-wavelength distributed feedback (DW-DFB) lasers have been developed, including integrated Y-waveguide DW-DFB lasers, multi-section DW-DFB lasers, and transverse coupling DW-DFB lasers. Monolithic integrated DW-DFB lasers boast a compact structure and stable performance. However, they usually require complex manufacturing processes and control circuits, which indicate high costs. Therefore, a dual-wavelength laser with a simple structure and good stability is urgently needed.

Methods In this paper, a monolithic integrated two-section DW-DFB (TS-DW-DFB) laser is proposed experimentally. The TS-DW-DFB laser consists of a DW-DFB laser section and a grating reflector (GR) section. The two sections use the same epitaxial layer structure and share the same waveguide. Both facets of the TS-DW-DFB laser are deposited with antireflection coatings. The total length of the TS-DW-DFB laser chip is 1000 µm, and the lengths of the DW-DFB laser section and GR section are both 500 µm. There is also an electrical isolator between the two sections. The gratings of the two sections are fabricated by the reconstruction-equivalent chirp (REC) technique with the same seed grating. The grating in the DW-DFB laser section is a linearly chirped sampling grating with two  $\pi$ -phase-shifts at its 1/3 and 2/3 length, and the chirp ratio of the grating is 50 nm/mm. The grating in the GR section is a uniform sampling grating, and the sampling period of the grating is equal to the sampling period at the center of the grating in the DW-DFB laser section. To eliminate the side lobes in the reflection spectrum of the GR section and to decrease their reflection to the side modes of the DW-DFB laser section, we equivalently apodize the sampling grating in the GR section by changing the sampling duty cycle. As a consequence, the two main modes of the DW-DFB laser section are symmetrically located at the stop-band center of the reflection spectrum of the sampling grating in the GR section, while the side modes lie outside of the stop-band. The reflectivity of the grating in the GR section for the main modes is much higher than that for the side modes, which can lead to a much higher side-mode suppression ratio (SMSR). Due to the reflection of the GR section, the threshold current and power difference of the main modes (PDM) of the TS-DW-DFB laser are reduced.

**Results and Discussions** Using the transfer matrix method, the output characteristics of TS-DW-DFB lasers with different grating structures are analyzed. The simulated results show that the proposed TS-DW-DFB laser has a lower threshold current and higher output power [Figs. 4 and 5(a)] compared with the single-section DW-DFB laser. The PDM can also be optimized by changes in the length of the GR section [Fig. 5(b)]. Given these results, the TS-DW-DFB laser is fabricated. In contrast, a single-section DW-DFB laser with the same structure of the DW-DFB laser section as the TW-DW-DFB laser is also fabricated. The optical spectra of the single-section DW-DFB laser and TS-DW-DFB laser are measured under the same biased currents (Fig. 7). It is obvious that the TS-DW-DFB laser has a smaller PDM and a larger SMSR than the single-section DW-DFB laser. The influence of the biased current of the GR section is adjusted to optimize the PDM and SMSR of the TS-DW-DFB laser is also discussed (Fig. 8), and the current of the GR section is adjusted to optimize the PDM and SMSR simultaneously. In addition, the operation stability is a figure of merits for dual-wavelength lasers. When the temperature of the TS-DW-DFB laser is tuned from 16  $^{\circ}$  to 40  $^{\circ}$ , both the wavelengths of the TS-DW-DFB laser section temperature is have length (Fig. 10(b)]. However, the PDM, SMSR, and wavelength spacing of the TS-DW-DFB laser are not changed significantly (Fig. 11). Their changes in one hour are observed when the operation temperature is kept at 24  $^{\circ}$ , and the bias currents of the DW-DFB laser section and GR section are set at 80 mA and 12 mA, respectively. The results reveal that the PDM, SMSR, and wavelength spacing remain stable within an hour.

**Conclusions** The in-cavity mode competition of dual-wavelength lasers is fierce, and thus, mode stability is a key figure of merit for dual-wavelength lasers. To reduce the power difference between the two main modes and improve SMSR, we propose a DW-DFB laser integrated with GR. The grating structure in the dual-wavelength laser is simulated by the transfer matrix method. The influence of GR on the threshold and PDM of the laser is analyzed. After that, a monolithic integrated TS-DW-DFB semiconductor laser is fabricated and experimentally demonstrated. The measured results show that the proposed method can improve the operation stability, enhance the SMSR, and reduce the PDM of the DW-DFB laser. In a stable operation state, the PDM of the dual-wavelength laser is less than 0.3 dB, and the SMSR is larger than 35 dB.

**Key words** lasers; dual-wavelength laser; distributed feedback semiconductor laser; two-section laser; photonic integration; monolithic integration