Damage resistance of \( \text{B}_4\text{C} \) reflective mirror irradiated by X-ray free-electron laser

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1. Introduction

With the special characteristics of high peak brightness and ultrashort pulse width, the X-ray free-electron laser (XFEL) is able to induce complicated dynamic processes when interacting with materials\(^{1-6}\). For the X-ray reflective optics used at XFEL beamlines, this interaction process can easily destroy or deteriorate these optical mirrors, thereby reducing the XFEL quality and affecting applications in user experiments. In recent years, various XFEL-induced damage mechanisms have been found, such as thermal melting\(^{7-11}\), ablation\(^{12}\), thermal stress\(^{8}\), thermal and nonthermal phase transition\(^{13,14}\). In-depth study of XFEL-induced damage process of optical film mirrors, analysis of film damage mechanisms, and preparation of X-ray mirrors with high damage resistance are of great significance.

In this paper, a simple theoretical model combining Monte Carlo simulation with the enthalpy method is provided to simulate the damage resistance of B\(_4\)C/Si-sub mirror under X-ray free-electron laser irradiation. Two different damage mechanisms are found, dependent on the photon energy. The optimum B\(_4\)C film thickness is determined by studying the dependence of the damage resistance on the film thickness. Based on the optimized film thickness, the damage thresholds are simulated at photon energy of 0.4–25 keV and a grazing incidence angle of 2 mrad. It is recommended that the energy range around the Si K-edge should be avoided for safety reasons.

Keywords: B\(_4\)C film; XFEL; damage mechanism; damage threshold; enthalpy method.

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The relationship between the damage threshold and the melting dose has been tested by free-electron laser damage experiments in the photon energy range from extreme ultraviolet (EUV) to X-ray\(^{7-12}\). Under the normal incidence condition, it was found that the measured threshold doses for B\(_4\)C, SiC are generally in agreement with the calculated melting dose in EUV range of 13.5–21.7 nm\(^{7}\). At X-ray photon energy of 0.83 and 10 keV, it was also demonstrated that the measured threshold doses for optical materials of B\(_4\)C, Si, SiO\(_2\), Pt, and Rh are quite consistent with the theoretical values\(^{8,12}\). However, at the grazing incidence condition\(^{9-11}\), the experimental damage thresholds are much higher than the calculated ones based on Eq. (1), where the electron collision escape, scattering, and other secondary processes are not easily taken into account. For simulating FEL-induced thermal and non-thermal phase transition processes, Ziaja et al.\(^{18,19}\) have developed an effective hybrid model combining tight binding molecular dynamics with Monte Carlo simulation and Boltzmann collision integrals for nonadiabatic electron–ion coupling. Since this hybrid model combines a variety of complicated theoretical calculations, it is obviously a time-consuming work to estimate the damage resistances of X-ray reflective mirrors with different film materials.

Presently, China is building Shanghai High repetition rate XFEL and Extreme light facility (SHINE), which is designed to cover the photon energy range of 0.4–25 keV\(^{20,21}\). B\(_4\)C is likely to be chosen as the main coating material for reflective
mirrors at SHINE because of its excellent characteristics of high hardness, high melting point, and optical performance. So, it is important to develop a relatively simple and convenient method to evaluate the damage resistance of B$_4$C mirrors at the working photon energy range for safety applications. In this paper, a simple theoretical model is provided to calculate the damage thresholds of B$_4$C/Si-sub mirrors based on the Monte Carlo simulation and the enthalpy method. Considering that the damage resistances are related to the film structure, the photon energy, and the grazing incidence angle, the effects of the film thickness on damage thresholds for B$_4$C/Si-sub mirror are analyzed, and the damage thresholds in the photon energy range of 0.4–25 keV are given and discussed at a typical grazing incidence angle of 2 mrad. It is found, unexpectedly, that the photon energy-dependent damage resistance is related not only with the B$_4$C film but also with the Si substrate.

2. Theoretical Model

When materials are irradiated by ultrashort femtosecond XFEL, the evolution processes of electronic excitation and relaxation will be involved. Through photoabsorption, X-ray photons interact primarily with inner-shell electrons to emit photoelectrons with high kinetic energy. Then the cascaded relaxation processes, such as the fluorescence and Auger decay as well as electron impact ionization, will be followed to exchange energy among electrons. The time scale for this electronic thermalization is about several hundred femtoseconds$^{[11,22,23]}$. Through this ultrafast process, the electrons are at a high temperature, while the crystal lattices are still at a low temperature. The thermalization process between electrons and lattices takes a longer time—about picoseconds—until reaching thermal equilibrium. Finally, the heat energy is transferred inside the material through thermal conduction. In this model, the photoabsorption and the relaxation processes, as well as the heat transfer, are considered, where the thermalization between electrons and lattices is assumed to be transient for XFEL interaction with mirrors. The Monte Carlo method is adopted to simulate X-ray photoabsorption and the cascading processes. The enthalpy method is used to simulate the heat transfer process.

2.1. Monte Carlo simulation of X-ray interaction with mirrors

Figure 1 gives a schematic of an XFEL beam incident on a B$_4$C/Si-sub mirror under grazing incidence. When the total reflection is fulfilled, most part of the incident photon energy $E_{ph}$ will be reflected due to the high reflectivity $R$. The absorbed energy $E_{ab} = (1 - R)E_{ph}$ by the mirror is only a few parts of the incident energy. The X-ray penetration depth inside the film medium is quite shallow, which is about several nanometers. However, the energetic photo- and secondary electrons generated in the X-ray interaction volume can escape from the mirror surface or travel deeply into the film or even into the substrate, which dramatically affects the energy deposition range.

In this study, X-ray energy deposited along the depth direction was simulated using Geant4 Monte Carlo simulation code$^{[24]}$, where the photoionization, Auger and fluorescence effects, and electron elastic and inelastic scattering processes were taken into account. Here the absorbed energy fraction AEF(z) represents the ratio of deposited energy per unit depth to the absorbed energy $E_{ab}$. Figure 2 gives AEF(z) for X-ray interaction with B$_4$C(50 nm)/Si-sub mirror at 1 keV and 12 keV with the grazing incidence angle of $\theta = 2$ mrad. For comparison, the theoretical results of AEF(z) considering only the exponential decay of the X-ray into the mirror are also calculated by

$$\frac{dI}{dz} = \frac{1}{d_x} \exp \left(-\frac{z}{d_x}\right).$$

(2)

Here the X-ray attenuation length $d_x$ can be calculated by

$$d_x = \frac{\lambda}{4\pi \beta} \sqrt{\frac{1}{2} ((\sin^2 \theta - 2\delta)^2 + 4\beta^2)^{1/2} + (\sin^2 \theta - 2\delta)},$$

(3)

where $\lambda$ is the X-ray wavelength, and $\delta$ and $\beta$ are the decrement of the real part and imaginary part of the complex refractive index $n = 1 - \delta - i\beta$, respectively.

At 1 keV, according to Eq. (2), most part of the absorbed energy (> 99.5%) should be localized in the surface layer of

![Fig. 1. Schematic of an XFEL beam incident on a B$_4$C/Si-sub mirror under grazing incidence.](image)

![Fig. 2. Absorbed energy fractions along the depth direction for B$_4$C(50 nm)/Si-sub irradiated by XFEL at 1 keV and 12 keV with the grazing incidence angle of 2 mrad.](image)
B₄C film, since the X-ray penetration depth is only 3.07 nm at the grazing angle of 2 mrad. However, the Monte Carlo simulation shows that ~20.2% of the absorbed energy is released and escaped from the B₄C surface by photo- and Auger electrons as well as X-ray fluorescence. Only ~79.8% of the absorbed energy is deposited into the B₄C film with the penetration depth about 30 nm, while the Si substrate absorbs almost nothing. At 12 keV, the simulation shows that 25.9% of the absorbed energy is deposited into the Si substrate, and the remaining energy is released from the mirror surface, and the remaining energy is deposited deeply into the mirror. Contrary to the case at 1 keV, the B₄C film is deposited at only 2.9% of the absorbed energy at 12 keV, while the Si substrate absorbs the most, to the deposition depth of ~1500 nm. Due to the transport by the energetic electrons, the deposition depth is dramatically increased, which is quite beneficial to enhance the damage resistance of X-ray mirrors, especially under the grazing incidence condition.

2.2. Enthalpy method

Materials irradiated by XFEL may undergo phase changes after absorbing enough energy in a short time. The methods for solving the phase transition process mainly include the equivalent heat capacity method and the enthalpy method[26,27]. The enthalpy method uses enthalpy as the dependent variable to solve the heat transfer differential equation. In this study, the enthalpy method was used to simulate the accumulation and transport of heat with a simple one-dimensional thermal diffusion model as represented by

\[
\frac{\partial h}{\partial t} = \frac{\partial}{\partial z} \left( k \frac{\partial h}{\partial z} \right) + S. \tag{4}
\]

Here, \( h \) is the enthalpy of the material with units of [J/m³], \( S \) is the heat source term with units of [W/m³], and \( z \) represents the depth direction. \( k \) is the thermal conductivity expressed in [W/(m·K)] and \( C \) is the thermal capacity expressed in [J/(K·m³)]. For B₄C and Si, the thermal conductivity and the thermal capacity as a function of temperature are provided in Refs. [28–30]. Here, a one-dimensional thermal diffusion model was considered because the projected area of the XFEL beam on the mirror is much larger than the depth of the heat-affected zone under the grazing incidence condition. For this calculation, the initial enthalpy value was assumed as 0 J/m³, which corresponds to an initial temperature of 298.15 K. The film surface was assumed to be adiabatic, and the sample bottom was kept at room temperature. The time- and depth-dependent heat source is expressed as

\[
S(z, t) = \sqrt{\frac{4 \ln 2}{\pi \tau_p} \frac{1 - R}{R} I_0 \exp(-4z^2/\tau_p^2)} \cdot \text{AEF}(z) \sin \theta, \tag{5}
\]

where \( I_0 \) is the incident laser fluence with units of [J/cm²], \( R \) is the reflectivity, \( \tau_p \) is the time duration of the heat source term, and \( \theta \) is the grazing incidence angle. AEF(\( z \)) is the absorbed energy fraction obtained by Monte Carlo simulation, as mentioned above. In this study, we assume that the absorption of XFEL gives rise to an instantaneous enthalpy or temperature rise inside materials.

3. Results and Discussion

3.1. Damage mechanisms of B₄C (50 nm)/Si-sub

The time-dependent depth distribution of enthalpy is shown in Fig. 3 for B₄C (50 nm)/Si-sub mirror irradiated under the grazing incidence angle of 2 mrad by XFEL at the fluences of 1250 J/cm² at 1 keV and 2.4 × 10⁵ J/cm² at 12 keV. These two particular fluences are chosen here, since the fluences of 1250 J/cm² at 1 keV and 2.4 × 10⁵ J/cm² at 12 keV are considered as the damage thresholds for B₄C/Si-sub mirror, as explained below. The reflectivities used in this simulation were determined by IMD software to be 0.991 and 0.999 at photon energies of 1 keV and 12 keV, respectively. The surface
The roughness of B\(_4\)C film was set as 0.3 nm, and the interface roughness of B\(_4\)C – Si was set as 0.2 nm.

According to the NIST website\(^{[30]}\), the enthalpy levels of 13.95 GJ/m\(^3\) and 3 GJ/m\(^3\) correspond to the melting temperature of B\(_4\)C and Si, respectively. At 1 keV in Fig. 3(a), the maximum enthalpy of 13.95 GJ/m\(^3\) is achieved on the surface of B\(_4\)C film at the fluence of 1250 J/cm\(^2\), while the maximum enthalpy of the Si substrate is still lower than 3 GJ/m\(^3\). It means the surface of B\(_4\)C film just reaches the melting point, but the temperature of the Si substrate is still lower than its melting point. So, the damage should occur on the top surface of B\(_4\)C film at 1 keV, with the damage threshold of 1250 J/cm\(^2\). At 12 keV in Fig. 3(b), the Si surface layer with a depth of about 1 \(\mu\)m has higher enthalpy than B\(_4\)C film, since more energy is deposited on Si-sub than on B\(_4\)C. At the fluence of 2.4 \(\times\) 10\(^5\) J/cm\(^2\), the melting damage occurs in Si-sub at \(\sim\)1.8 nm below the B\(_4\)C-sub interface. From the simulations, it can be found that there are two different damage mechanisms for B\(_4\)C/Si-sub mirror, depending on the incident photon energy, where one is caused by the melting of B\(_4\)C film surface and the other one is caused by the melting of Si-sub near the B\(_4\)C-sub interface. The simulation result at 12 keV is quite consistent with the experimental observation by Aquila \textit{et al.}\(^{[10]}\) that the damage occurred at the B\(_4\)C-sub interface for B\(_4\)C(50 nm)/Si-sub at the grazing incidence angle of 2 mrad. The experimental damage threshold for the B\(_4\)C mirror at 12 keV was also given as > 10,000 J/cm\(^2\), which is different from our result because of the different reflectivity used. In this study, the theoretical reflectivity was calculated using the density of B\(_4\)C bulk material, which leads to a higher reflectivity, since the bulk material is usually denser than the film. If the reflectivity of 0.975 given by Aquila \textit{et al.}\(^{[10]}\) is adopted in our simulation, the damage threshold of 9000 J/cm\(^2\) is obtained, which is close to the experimental result. Obviously, the actual reflectance has a great influence on the damage resistance of mirrors. Thus, it is quite important to calibrate the reflectivity accurately before using the reflective mirrors at the XFEL beamlines.

3.2. Correlation between B\(_4\)C film thickness and damage resistance

The film thickness of the reflective mirror has strong effects on the damage resistance\(^{[31,32]}\). However, no one has discussed this important question about a B\(_4\)C reflective mirror being used at XFEL beamlines. In order to answer this question, we explored the damage resistance as a function of B\(_4\)C film thickness at the photon energy of 12 keV and the grazing incidence angle of 2 mrad. Figure 4 gives the simulated AEF(z) for B\(_4\)C/Si-sub with the film thickness ranging from 10 to 200 nm. It can be seen that all the energy deposited per unit depth shows a slight upward trend along the depth direction in B\(_4\)C, while it decreases gradually with a deposition depth about 1 \(\mu\)m in the Si substrate. As the B\(_4\)C film thickness increases from 10 to 200 nm, the ratio of the escape energy to the absorbed energy decreases from 27.5% to 20.4%, and most of the absorbed energy is still deposited in the Si substrate. Therefore, the melting damage on the surface of the Si-sub is still the main damage mechanism at 12 keV for B\(_4\)C mirrors with different film thicknesses.

Using this model, the damage thresholds for B\(_4\)C mirrors with different film thicknesses are calculated as given in Fig. 5. As the B\(_4\)C film thickness increases from 10 to 50 nm, the damage threshold increases from 5.5 \(\times\) 10\(^4\) J/cm\(^2\) to 2.4 \(\times\) 10\(^5\) J/cm\(^2\), which is mainly due to the reflectivity increasing from 0.996 to 0.999. For the film thickness above 50 nm, the damage thresholds remain almost at a constant value of 2.4 \(\times\) 10\(^5\) J/cm\(^2\), where the optimal damage resistance is achieved. Considering the laser resistance and the optical performance, the typical film thickness of 50 nm is still recommended for a B\(_4\)C reflective mirror in practical applications.

![Fig. 4. Absorbed energy fraction for B\(_4\)C/Si-sub mirror with different B\(_4\)C film thicknesses at the photon energy of 12 keV and the grazing incidence angle of 2 mrad.](image)

![Fig. 5. Damage thresholds of B\(_4\)C mirrors with different B\(_4\)C film thicknesses (solid black squares) at 12 keV and the grazing angle of 2 mrad; reflectivity of B\(_4\)C mirrors as a function of film thickness calculated with IMD software (blue line).](image)
3.3. FEL damage resistance of B₄C/Si-sub at X-ray energy of 0.4–25 keV

Figure 6 gives the damage thresholds of a B₄C(50 nm)/Si-sub mirror irradiated by XFEL in the energy range of 0.4–25 keV at the grazing incidence angle of 2 mrad. These damage thresholds were determined according to the theoretical reflectivity, as shown by the solid red line.

When X-ray energy increases from 0.4 to 1.8 keV, the damage threshold increases sharply from 2.1 × 10⁻³ J/cm² to 9.6 × 10⁻³ J/cm², where the damage occurs at the surface of the B₄C film similar to the case at 1 keV as given in Fig. 3(a). At 1.839 keV, which corresponds to the Si K-edge, there is a sharp drop for the damage threshold due to the enhanced photoabsorption by the Si substrate. From 1.9 to 12 keV, the damage threshold increases gradually from 1.7 × 10⁻³ J/cm² to 2.4 × 10⁻² J/cm², where the damage occurs in the surface layer of Si-sub near the B₄C-sub interface, which is similar to the case at 12 keV, as shown in Fig. 3(b). Additionally, with the increase of photon energy, the photoelectron kinetic energy increases, which makes the deposition depth deeper. So, the increase of the damage threshold from 1.9 to 12 keV is mainly due to the deeper deposition depth inside the Si. From 12 to 15 keV, the slight drop in damage threshold is due to a slight decrease of reflectivity. For photon energy above 15 keV, the damage threshold of B₄C/Si-sub mirror drops sharply, which is caused by the low reflectivity. From the simulation, it is demonstrated that the energy range around the Si K-edge at 1.839 keV should be avoided for safety reasons when applying the B₄C/Si-sub mirror at XFEL beamlines. Additionally, the working angle for the B₄C mirror should be less than 2 mrad when applying it at a photon energy range above 15 keV.

4. Conclusion

In this study, a simple model based on Monte Carlo simulation and the enthalpy method is proposed to evaluate the damage resistance of B₄C/Si-sub reflective mirrors irradiated by XFEL at a typical grazing incidence angle of 2 mrad. It is found that two different mechanisms are responsible for XFEL damage of B₄C mirrors. In the photon energy range of 0.4–1.8 keV, the melting damage happens on the surface layer of B₄C film because the absorbed energy is mainly deposited inside the B₄C film. At the photon energy above 1.9 keV, the melting damage occurs in the surface layer of the Si substrate close to the B₄C-sub interface because the absorbed energy is deposited deeply into the Si substrate by the transport of the energetic electrons. For B₄C/Si-sub mirror with an optimum film thickness of 50 nm, the damage thresholds are determined in the photon energy range of 0.4–25 keV using the theoretical reflectivity. According to the simulations, it is suggested that the energy range around the Si K-edge should be avoided for safety reasons when using a B₄C/Si-sub mirror at XFEL beamlines. It is expected this research will be helpful for the design and operation of reflective mirrors at XFEL beamlines.

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References


