

# High-quality-factor optical microresonators fabricated on lithium niobate thin film with an electro-optical tuning range spanning over one free spectral range [Invited]

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We demonstrate high-quality (intrinsic  $Q$  factor  $\sim 2.8 \times 10^6$ ) racetrack microresonators fabricated on lithium niobate thin film with a free spectral range (FSR) of  $\sim 86$  pm. By integrating microelectrodes alongside the two straight arms of the racetrack resonator, the resonance wavelength around 1550 nm can be red shifted by 92 pm when the electric voltage is raised from  $-100$  V to 100 V. The microresonators with the tuning range spanning over a full FSR are fabricated using photolithography assisted chemo-mechanical etching.

**Keywords:** microresonators; lithium niobate; electro-optical tuning; chemo-mechanical etching.

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## 1. Introduction

Whispering gallery mode (WGM) optical microresonators play a crucial role in both photonic research and applications owing to the strong confinement of light fields resulting from the characteristic high quality ( $Q$ ) factors<sup>[1]</sup>. Currently, most optical microresonators fabricated on chip have a footprint below 1 mm. Such small microresonators can be fabricated on various types of materials such as silica<sup>[2]</sup>, semiconductors<sup>[3]</sup>, crystals<sup>[4]</sup>, and polymers<sup>[5]</sup> using lithographic techniques. The integrated microresonators have enabled a broad range of functionalities including filtering, wavelength conversion, sensing, optomechanics, and comb generation<sup>[6,7]</sup>. However, for some applications such as Brillouin lasers<sup>[8,9]</sup>, optical-frequency synthesizers<sup>[10]</sup>, microwave photonics<sup>[11]</sup>, spectroscopy<sup>[12]</sup>, gyroscope sensors<sup>[13]</sup>, optical atomic clocks<sup>[14]</sup>, and high-resolution spectrometers<sup>[15]</sup>, WGM microresonators with

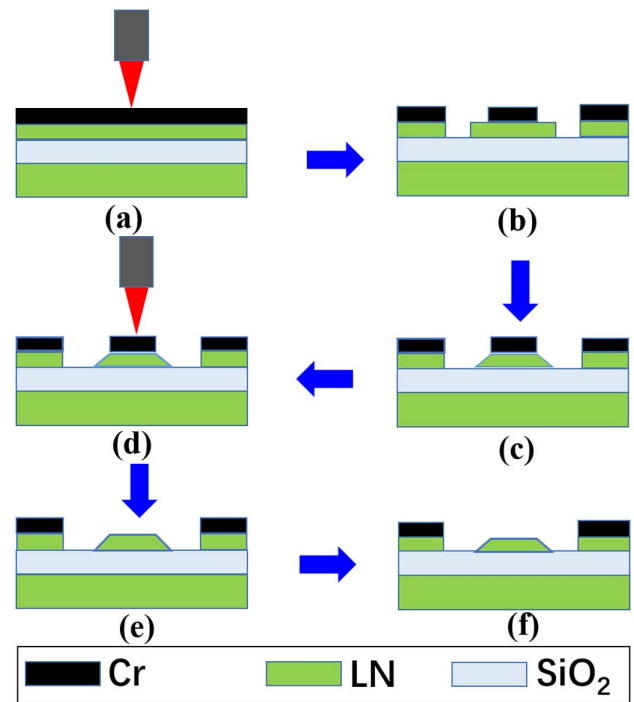
perimeters up to the centimeter level are desirable. Fabrication of such large microresonators using lithographic techniques is challenging. Usually, electron-beam lithography (EBL) and deep ultraviolet (DUV) lithography cannot directly define such large microresonators, and thus stitching is frequently employed to meet the size requirement<sup>[16,17]</sup>. Nevertheless, stitching can inevitably cause fabrication errors, which will spoil the  $Q$ . Ultraviolet (UV) lithography can directly expose such large microresonators easily one time, so it does not need stitching. However, UV lithography does not have high fabrication resolution to expose such a narrow waveguide<sup>[18]</sup>. On the other hand, focused ion beam (FIB) milling is a low-throughput fabrication technique compared to optical lithography technologies<sup>[19]</sup>. It is time consuming to fabricate the large WGM microresonators using EBL and FIB, although high fabrication resolutions are readily achievable with the two techniques.

Here, we demonstrate high- $Q$  optical microresonators fabricated on lithium niobate (LN) thin film (900-nm-thick) with an electro-optical (EO) tuning range spanning over one free spectral range (FSR). The advantage of choosing LN as the substrate material is the strong EO property. In particular, our fabrication technique based on photolithography assisted chemo-mechanical etching (PLACE) allows us to define the mask pattern of a racetrack WGM resonator with a footprint size of  $2.5\text{ mm} \times 6\text{ mm}$  in only  $\sim 1\text{ h}$ . The fabricated device shows an intrinsic  $Q$  factor of  $2.8 \times 10^6$ . We also integrate microelectrodes (600-nm-thick) alongside the two straight arms of the racetrack resonator. We examine the EO tunability and tuning range of the microresonator. We observe that when the electric voltage is raised from  $-100\text{ V}$  to  $100\text{ V}$ , the resonance wavelength around the  $1550\text{ nm}$  can be red tuned by  $92\text{ pm}$ , which exceeds the FSR ( $\sim 86\text{ pm}$ ) of the fabricated microresonator.

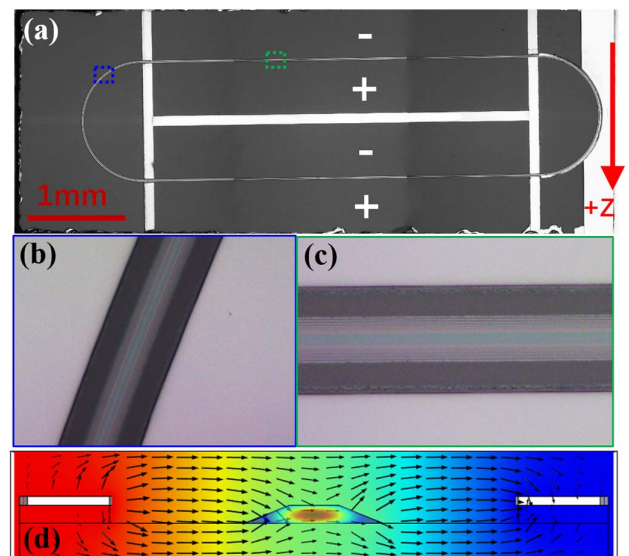
## 2. Materials and Methods

In our experiment, the on-chip LN racetrack resonator integrated with Cr electrodes was fabricated on a commercially available X-cut LN on insulator (LNOI) wafer (NANOLN, Jinan Jingzheng Electronics Co., Ltd.). The LN thin film with a thickness of  $900\text{ nm}$  was bonded onto a silica layer with a thickness of  $\sim 2\text{ }\mu\text{m}$ , and the silica layer was grown on a  $0.5\text{-mm-thick}$  LN substrate. A  $600\text{-nm-thick}$  chromium (Cr) film was further deposited on the top surface of the LNOI by magnetron sputtering. The fabrication process includes four steps, as illustrated in Fig. 1. Firstly, the Cr film on the LNOI sample was patterned into a stripe-shaped mask using a space-selective femtosecond laser (PHAROS, LIGHT CONVERSION Inc.). Subsequently, the chemo-mechanical polishing (CMP) process was performed to fabricate the LN waveguide using a wafer polishing machine (NUIPOL802, Kejing Inc.). The LN thin film protected by the Cr mask was preserved after the CMP process, whereas the remaining LN in the opening area was completely removed. This step allows us to create LN waveguides with extremely smooth sidewalls, ensuring high- $Q$  factors for the fabricated microresonators<sup>[20–22]</sup>. Next, femtosecond laser ablation was carried out again to remove the Cr mask left behind on the LN waveguides. Finally, a post-CMP process was performed for thinning the LN disk and smoothing the top surface. In the current experiment, the fabricated ridge waveguides have a top width of  $\sim 2\text{ }\mu\text{m}$ , and the Cr electrodes on the two sides of the waveguide are separated by a distance of  $\sim 20\text{ }\mu\text{m}$ . It takes about  $1\text{ h}$  in total to produce the whole racetrack microresonator, as shown in Fig. 2(a). The microelectrodes were fabricated by patterning the Cr film using femtosecond laser micromachining.

Figure 2(a) presents the top view of the on-chip LN racetrack resonator integrated with the Cr electrodes. The diameter of the two half-circles of the racetrack resonator is  $1.2\text{ mm}$ , and the length of the straight arms of the racetrack is  $4\text{ mm}$ . Therefore, the perimeter of the microresonator is  $\sim 11.77\text{ mm}$ . Figures 2(b) and 2(c) present the zoom-in images of the curved and straight waveguides of the LN racetrack resonator, respectively.



**Fig. 1.** Process flow of fabricating an on-chip LN racetrack resonator integrated with Cr microelectrodes. (a), (b) Patterning the Cr thin film into a stripe mask using femtosecond laser microfabrication. (c) Etching of the LNOI layer by chemo-mechanical polishing. (d), (e) Selective removal of the stripe Cr mask on the LN racetrack resonator using femtosecond laser ablation. (f) Reducing the thickness of LN racetrack resonator by a post chemo-mechanical polishing, which leads to the smooth top surface on the LN waveguide.



**Fig. 2.** (a) Top view optical micrograph of the on-chip LN racetrack resonator integrated with Cr electrodes. Zoom-in optical micrographs of the (b) curved and (c) straight waveguides. (d) Distributions of the optical (TE) field and electrical field overlapping each other simulated using COMSOL. The arrows indicate the direction of the electric field.

The surface roughness on the waveguides was determined with an atomic force microscope (AFM) to be subnanometer<sup>[21,22]</sup>. Figure 2(d) shows the numerically simulated optical and electric fields overlapping each other. Since the device was fabricated on an X-cut LN wafer, the transverse-electric (TE) optical modes can be efficiently tuned thanks to the highest EO tensor component ( $r_{33}$ ) of LN.

### 3. Results

Figure 3(a) presents the digital-camera-captured picture of the racetrack resonator as compared with a one-yuan Renminbi coin. The zoom-in optical micrograph in Fig. 3(b) further shows the details of the racetrack resonator and the Cr electrodes as well as the aluminum wire (diameter  $\sim 25 \mu\text{m}$ ) connecting the Cr electrodes to the voltage generator. To characterize the tunability of the LN racetrack resonator, we used an experimental setup, as illustrated in Fig. 3(c). Here, a tunable laser (TLB-6728, Newport Inc.) was used as the light source. To enable efficient coupling of the light into the racetrack microresonator, we first fabricated a straight waveguide on another LNOI wafer. The straight waveguide has the same geometric parameters such as the cross sectional shape, thickness, and width as that of the waveguide in the LN racetrack resonator. The light from the tunable laser was coupled into the straight waveguide using a lensed fiber. The straight waveguide was then brought close to

the racetrack resonator from the top to allow for evanescent coupling, as shown in Fig. 3(c). Due to the same geometric parameters of the straight waveguide and the racetrack waveguide, phase matching can be ensured to enable efficient coupling between them. The polarization state of the input light was adjusted using an in-line fiber polarization controller. The output of the LN waveguide was connected to a photodetector (New Focus 1811-FC-AC, Newport Inc.) using a lensed fiber for measuring the transmission spectrum and the Q factor of the racetrack resonator. A programmable linear direct current (DC) stabilized power source (IPMP500-0.6L, INTERLOCK Technologies Inc.) was used as the voltage generator, which provided a variable voltage ranging from  $-100 \text{ V}$  to  $100 \text{ V}$ .

Figure 4(a) shows the measured transmission spectrum in the wavelength range between 1544 and 1552 nm. The resonance lines appeared regularly spaced, indicating that mostly the fundamental mode is excited in the LN racetrack resonator. This should be a result of the same geometric parameters of the upper straight waveguide and the lower racetrack waveguide, as shown in Fig. 3(c). In this case, the spatial mode profiles in the two waveguides can most efficiently overlap. Thus, as the fundamental mode is excited in the upper waveguide with the lensed fiber, the same fundamental mode will be excited in the racetrack resonator as well. As shown in Fig. 4(b), the FSR of the microresonator was determined to be 86 pm. One of the WGMs at the resonant wavelength of 1547.93 nm was chosen for the measurement of the loaded Q factor  $Q_L$ . By fitting with the Lorentz function,  $Q_L$  is determined to be  $1.4 \times 10^6$ , which corresponds to an intrinsic Q factor  $Q_i \sim 2.8 \times 10^6$ , as evidenced by the critical coupling characteristic (i.e., the very deep dips in the transmission spectrum) in Fig. 4(b).

Numerical simulations were performed using COMSOL to reveal the optical mode profile in our LN waveguide [see Fig. 2(d)]. In particular, the calculated effective index  $n_{\text{eff}}$  and group index  $n_g$  as functions of the wavelength are shown in Figs. 5(a) and 5(b), respectively. At the wavelength of 1550 nm,  $n_{\text{eff}}$  and  $n_g$  are calculated to be  $\sim 2.098$  and  $\sim 2.329$ , respectively. In combination with the experimentally determined perimeter of our LN racetrack of  $\sim 11.77 \text{ mm}$ , it can be derived that the FSR  $\approx 87.6 \text{ pm}$ . The theoretical simulated FSR agrees well with our experimentally measured FSR (86 pm).

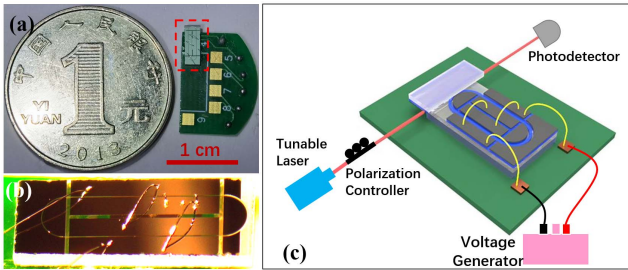


Fig. 3. (a) Picture of the LN racetrack resonator integrated with Cr electrodes. (b) Zoom-in image of the LN racetrack resonator integrated with Cr electrodes. (c) Schematic of the experimental setup for characterizing the Q factor and tunability of the device.

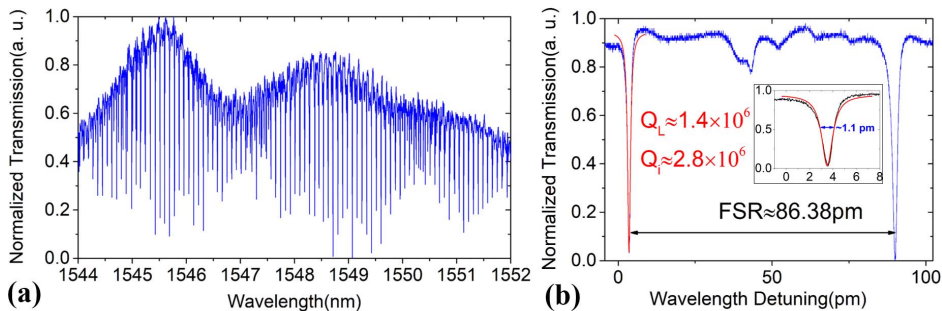


Fig. 4. (a) Transmission spectrum of the LN racetrack resonator. (b) The Lorentz fitting [red curve] reveals a loaded Q factor of  $1.4 \times 10^6$ , corresponding to an intrinsic Q factor of  $2.8 \times 10^6$  as measured at the 1547.93 nm wavelength; the linewidth of 1.1 pm is shown in inset.

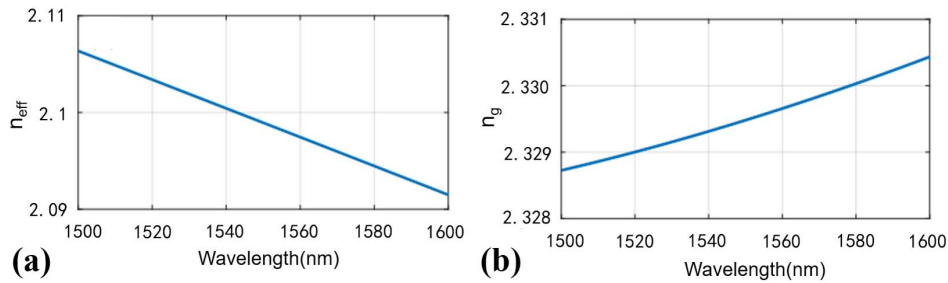


Fig. 5. (a) Calculated effective index  $n_{\text{eff}}$  and (b) group index  $n_g$  of the optical mode in Fig. 2(d) for the wavelength range between 1550 and 1600 nm.

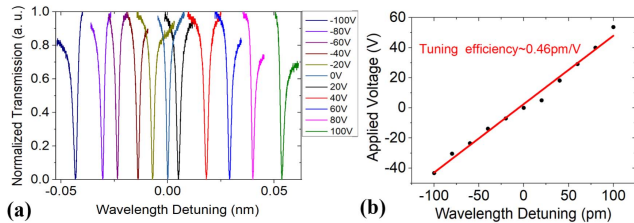


Fig. 6. (a) Resonance wavelength continuously red shifts with the increasing voltage. (b) The linear fit reveals an electrical tuning rate of  $\sim 0.46$  pm/V and indicates that the tuning range spans over a full FSR.

Finally, we carried out the real-time tuning of the racetrack microresonator by adding a tunable electric voltage on the Cr microelectrodes. Figure 6(a) shows a group of transmission spectra recorded near the 1550 nm wavelength at various electric voltages in the range from  $-100$  V to  $+100$  V with a voltage tuning step of 20 V. We observe that by increasing the electric voltage from  $-100$  V to  $+100$  V, the resonance wavelength can be continuously red tuned by  $\sim 92$  pm, which spans over one FSR ( $\sim 86$  pm). The fitting line in Fig. 6(b) indicates a linear dependence of the resonance wavelength on the applied electric voltage and a tuning rate of  $\sim 0.46$  pm/V.

#### 4. Discussions

This is the first attempt of fabricating a large WGM microresonator on an LNOI using the PLACE fabrication technique. Before, we have shown that freestanding microdisks with diameters of  $\sim 100$   $\mu\text{m}$  fabricated on the LNOI substrates using the PLACE technique can easily achieve  $Q$  factors above  $10^7$ <sup>[20–23]</sup>. The racetrack microresonator shows a  $Q$  factor nearly one order of magnitude lower. Most likely, this is caused by the scratches and defects that can be more easily generated on large photonic structures during the CMP. This suggests that more stringent control on the polishing and sample cleaning processes is necessary for maintaining the high  $Q$  factors in the fabrication of the large microresonators. In principle, the sub-nanometer surface roughness provided by the PLACE fabrication technique is sufficient to support  $Q$  factors above  $10^7$  for the large microresonators. We also notice that the EO tuning rate of 0.46 pm/V is lower in comparison with the results reported in the recent literature<sup>[24,25]</sup>. This is because the waveguide designed in this

work has been optimized for dispersion engineering, which features a relatively broad width at the bottom of the waveguide ( $\sim 7$   $\mu\text{m}$ ). Thus, the two electrodes on the two sides of the waveguides are separated from each other by  $\sim 20$   $\mu\text{m}$ , which is the main cause of the relatively low tuning efficiency. However, in the current work, our goal is to achieve the large tuning range covering the full FSR but not the high-speed EO modulation. The demonstrated tuning efficiency is sufficiently high for such application.

#### 5. Conclusions

To conclude, we have demonstrated high- $Q$  optical microresonators fabricated on LN thin film with an EO tuning range spanning over a full FSR. The EO tunable racetrack resonator of the perimeter above 1 cm features an FSR of  $\sim 86$  pm, an intrinsic  $Q$  factor of  $\sim 2.8 \times 10^6$ , and an EO tuning rate of 0.46 pm/V. The device provides a technological platform for a range of applications such as filtering, microwave photonics, sensing, and information processing.

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