

Research on Volume Scattering Phase Function under Ultraviolet Non-Line-of-Sight Single Scattering Link

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Abstract Based on ultraviolet single scattering link communication, the scattering particle is regarded as a differential volume, whose phase function after scattering is integrated under the correct scattering angles on the different positions within the effective scattering volume. Furthermore, the process is analyzed and formulas are deduced, which provide good support and solid foundation for ultraviolet communications.

Key words optical communication; ultraviolet communication; non-line-of-sight; single scattering; phase function

OCIS codes 290.1310; 290.5825; 290.4210

紫外光非视距单次散射链路中体散射相函数的研究

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摘要 针对椭球坐标系下的紫外光单次散射链路模型,将共同散射体中的各个散射粒子看作一个体积微分,对散射体中不同位置上的微分粒子散射之后造成的能够落入接收范围的散射角对应的相函数进行体积积分和运算,并进行了过程的分析与公式的推导,为研究紫外光通信提供了良好的支持与扎实的基础。

关键词 光通信;紫外光通信;非视距;单次散射;相函数

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1 Introduction

The aim of any communication system is transferring information from one location to another. Information shift is almost completed through modulating information to electromagnetic carrier frequency. Modulated carrier is then transmitted to the destination and demodulated after received. In radio system, carrier is selected from the radio frequency. Microwave and millimeter wave are selected as carriers in the microwave and millimeter communication system. Optical frequencies are chosen in optical communication, such as infrared, visible and ultraviolet (UV)^[1]. Among these communications, optical communication has superiority due to the broad unlicensed frequency and wide bandwidth.

Infrared and ultraviolet are both good carriers in optical communication^[2]. Infrared can work only when transmitter correctly faces to the receiver. But this term can not be satisfied in harsh conditions^[3]. Since the wavelength of ultraviolet is shorter than infrared, ultraviolet is scattered more strongly than infrared by constituents in atmosphere. Thus, ultraviolet taken as carrier can achieve non-line-of-sight (NLOS) communication, which has wider application than infrared communication requiring strict pointing between transmitter and receiver.

This paper mainly studies the scattering phase function, which shows the different scattering intensity with the change of scattering angle. The scattering particle in the common volume is seen as a differential volume and the scattering phase function can be achieved through integrating under the correct scattering angles on the different positions within the effective scattering volume.

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2 Ultraviolet single scattering link model

The model is shown in Fig.1, in which T is transmitter position, R is receiver position, θ_1 and ϕ_1 are transmitter apex and divergence, θ_2 and ϕ_2 are receiver apex and field of view (FOV), V is overlapping volume, r is distance between T and R , r_1 is distance between T and V , r_2 is distance between V and R , scattering angle $\theta = \theta_1 + \theta_2$. Ultraviolet is transmitted with apex and divergence as θ_1 and ϕ_1 at $t=0$. When $t=r_1/c$, ultraviolet reach overlapping volume V which is full of kinds of particles. These particles scatter ultraviolet to every direction. Those dropping in receiving range can be received by destination receiver at $t=(r_1+r_2)/c$.

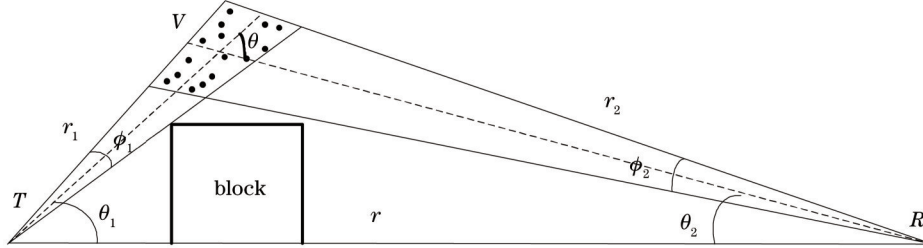


Fig.1 Single scattering link model of NLOS UV communication

The whole communication can be seen as NLOS process, which can also be seen as two line-of-sight processes, like $T \rightarrow V$ then $V \rightarrow R$. According to Ref.[4], receiving power at receiver is defined as^[3]

$$P = \left(\frac{P_1}{\Omega_1} \right) \left[\frac{\exp(-K_e r_1)}{r_1^2} \right] \left[\frac{K_s}{4\pi} p(\theta) V \right] \left(\frac{\lambda}{4\pi r_2} \right)^2 \exp(-K_e r_2) \frac{4\pi A_r}{\lambda^2}, \quad (1)$$

where P_1 is transmitting power, Ω_1 is transmitting solid angle, K_e is extinction coefficient, A_r is receiving aperture, λ is wavelength, $P(\theta)$ is scattering phase function. From Eq.(1) it's found that receiving power depends on the same scattering angle θ wherever the particle lies within the overlapping volume.

3 Volume scattering phase in ultraviolet communication

Ultraviolet scattering is mainly divided into two kinds. When particle size is shorter than UV wavelength, Rayleigh scattering happens. When particle size is longer than UV wavelength, Mie scattering happens^[5]. Rayleigh scattering is diffusion scattering, which characterizes atmosphere channel as quiet and clean one. However, Mie scattering is raised by big particles, i.e., fog, smoke, and little globule, which is even severe. Therefore, study of Mie scattering is of important significance for ultraviolet communication.

3.1 Mie scattering phase function

The ratio of scattering energy on a given direction in unit solid angle to average energy on all directions in unit solid angle is called scattering phase function. Scattering phase function is an important parameter in particle scattering character, which reflects scattering ability of atmospheric particles for different scattering angles between 0° and 180° .

Mie scattering process is very complex. Except for traditional Mie theory, empirical Mie phase function, such as Henyey-Greenstein function, revised Henyey-Greenstein function, and Henyey-Greenstein function defined by Cronette and Shanks are used in analoguing atmospheric transmission^[6-7].

Henyey-Greenstein phase function (H-G) is defined as^[6]

$$p_{\text{HG}}(\theta, g) = \frac{1 - g^2}{(1 + g^2 - 2g \cos \theta)^{3/2}}, \quad (2)$$

Henyey-Greenstein function (H-G1) is defined by Cronette and Shanks as^[6]

$$p_{\text{HG1}}(\theta, g) = \frac{3}{2} \times \frac{1 - g^2}{2 + g^2} \frac{1 + \cos^2 \theta}{(1 + g^2 - 2g \cos \theta)^{3/2}}, \quad (3)$$

Revised Henyey-Greenstein function (H-G2) is shown as^[6]

$$p_{\text{HG2}}(\theta, g) = \frac{1 - g^2}{4\pi} \left[\frac{1}{(1 + g^2 - 2g \cos \theta)^{3/2}} + f \frac{0.5(3 \cos^2 \theta - 1)}{(1 + g^2)^{3/2}} \right], \quad (4)$$

where g is asymmetry factor and f is scattering factor.

Seen from Eqs.(2)~(4), change of θ will change phase function no matter which function is employed to compute phase function.

3.2 Scattering phase function in UV communication

In section 2, the receiving power is defined as Eq.(1). Seen from Eq.(1), the overlapping volume can be imagined as countless differential volume dv . Each dv takes responsibility of scattering ultraviolet. Receiving power is computed through integrating effect of all dv .

Each dv is taken as a scattering particle. Due to each dv at different location within overlapping volume, the apex angles from T to dv and dv to R for each dv are different, which is to say, for each dv on different location, θ_t and θ_r are different from others. After scattering incident wave, all scattering angles which can drop in receiving range for each dv are $\theta = \theta_t + \theta_r$. Thus, the scattering angle for each dv is different. φ_t is transmitting divergence angle associated with scattering volume, φ_r is receiving divergence angle associated with scattering volume, θ_t is the angle between baseline and line of scattering volume and transmitter, θ_r is the angle between baseline and line of scattering volume and receiver, ξ_{max} and ξ_{min} are the maximum and minimum limits of X axis. As shown in Fig.2, for example, three dv on different positions have different scattering angles $\theta_{s1}, \theta_{s2}, \theta_{s3}$. Due to $r > D^2/\lambda$ (D is the size of scattering particle), scattering of each particle in the overlapping volume is irrelevant. Therefore, we can add the contribution of every particle directly. Based on this condition, we should integrate the different phase functions raised by each dv .

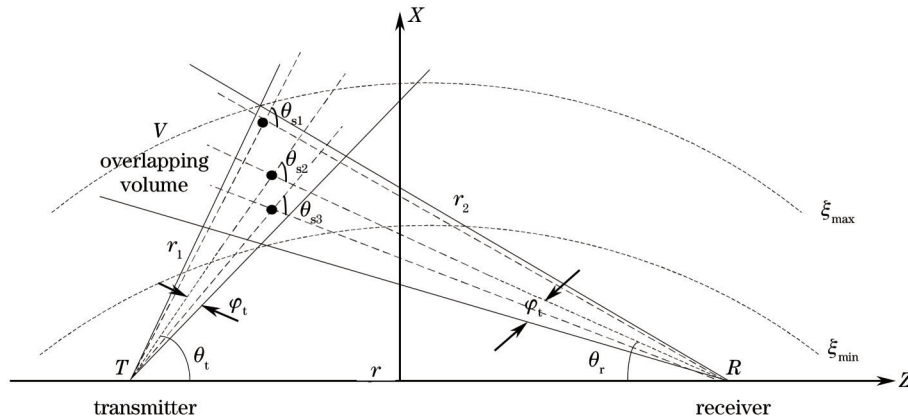


Fig.2 Different scattering angles within receiving range in the effective scattering volume

3.3 Volume scattering phase function based on prolate-spheroid coordinate

Prolate-spheroid coordinate is shown in Fig.3, in which ellipse is rotated around the main axis to complete a prolate-spheroid surface. Arbitrary point on prolate-spheroid is determined with a radial coordinate ξ , an angular coordinate η and an azimuthal coordinate φ ^[8]. Fig.4 is the application of prolate-spheroid coordinate in UV single scattering link model. Integration of phase function $P(\theta)$ is to integrate ξ, η, φ within overlapping volume, which can be defined as

$$\int_{\xi_{min}}^{\xi_{max}} \int_{\eta_1, \psi_1}^{\eta_2, \psi_2} P(\cos \theta) d\varphi d\xi \quad (5)$$

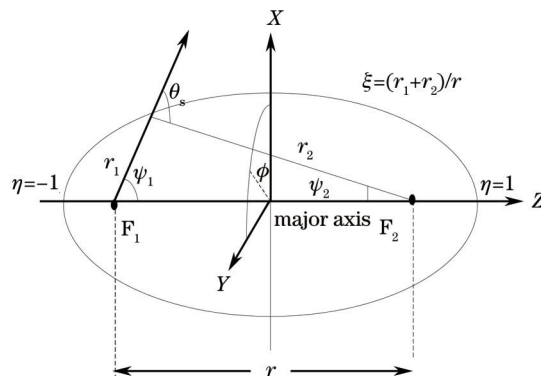


Fig.3 Prolate-spheroid coordinate

where scattering angle $\theta = \psi_1 + \psi_2$. When parameters relevant to ψ_1 and ψ_2 change, phase function will change. Three ranges of ξ, η, φ are discussed as follows.

3.3.1 Radial coordinate

In Fig.3, η_1 is a function related to ξ and ψ_1 and η_2 is a function related to ξ and ψ_1 ^[8], F_1 is the transmitter, F_2 is the receiver, ψ_1 is transmitting angle, and ψ_2 is receiving angle. There are functions shown as below:

$$\cos(\psi_1) = (1 + \xi\eta) / (\xi + \eta), \quad (6)$$

$$\cos(\psi_2) = (1 - \xi\eta) / (\xi - \eta). \quad (7)$$

According to Eqs.(6) and (7),

$$\eta_1 = \frac{\xi \cos(\psi_1) - 1}{\xi - \cos(\psi_1)} = \frac{\xi \cos(\theta_t) - 1}{\xi - \cos(\theta_t)}, \quad (8)$$

$$\eta_2 = \frac{1 - \xi \cos(\psi_2)}{\xi - \cos(\psi_2)} = \frac{1 - \xi \cos(\theta_r)}{\xi - \cos(\theta_r)}, \quad (9)$$

where $\psi_1 = \theta_t$, $\psi_2 = \theta_r$. According to Eqs.(8) and (9), change of η_1 and η_2 is mainly relevant to ξ .

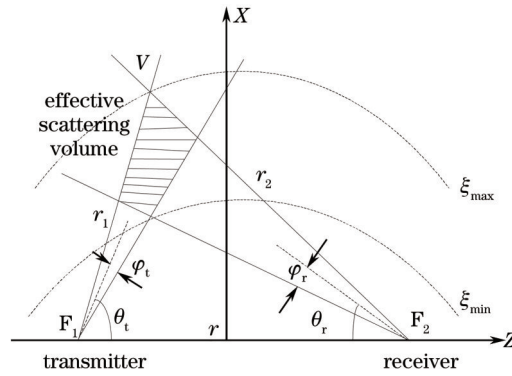


Fig.4 UV link under prolate-spheroid coordinate

3.3.2 Angular coordinate

When ξ adopt minimum ξ_{\min} , $\eta_1 = \eta_2$. Make Eq.(8) equivalent to Eq.(9), where $\psi_1 = \theta_t - \varphi_t$, $\psi_2 = \theta_r - \varphi_r$, and ξ_{\min} is defined as

$$\xi_{\min} = \frac{1 + \cos(\psi_1) \cdot \cos(\psi_2)}{\cos(\psi_1) + \cos(\psi_2)} + \sqrt{\left[\frac{1 + \cos(\psi_1) \cdot \cos(\psi_2)}{\cos(\psi_1) + \cos(\psi_2)} \right]^2 - 1} = \frac{1 + \cos(\theta_t - \varphi_t) \cdot \cos(\theta_r - \varphi_r)}{\cos(\theta_t - \varphi_t) + \cos(\theta_r - \varphi_r)} + \sqrt{\left[\frac{1 + \cos(\theta_t - \varphi_t) \cdot \cos(\theta_r - \varphi_r)}{\cos(\theta_t - \varphi_t) + \cos(\theta_r - \varphi_r)} \right]^2 - 1}. \quad (10)$$

when ξ adopts maximum ξ_{\max} , $\eta_1 = \eta_2$. Make Eq.(8) equivalent to Eq.(9), where $\psi_1 = \theta_t + \varphi_t$, $\psi_2 = \theta_r + \varphi_r$, and ξ_{\max} is defined as

$$\xi_{\max} = \frac{1 + \cos(\psi_1) \cdot \cos(\psi_2)}{\cos(\psi_1) + \cos(\psi_2)} + \sqrt{\left[\frac{1 + \cos(\psi_1) \cdot \cos(\psi_2)}{\cos(\psi_1) + \cos(\psi_2)} \right]^2 - 1} = \frac{1 + \cos(\theta_t + \varphi_t) \cdot \cos(\theta_r + \varphi_r)}{\cos(\theta_t + \varphi_t) + \cos(\theta_r + \varphi_r)} + \sqrt{\left[\frac{1 + \cos(\theta_t + \varphi_t) \cdot \cos(\theta_r + \varphi_r)}{\cos(\theta_t + \varphi_t) + \cos(\theta_r + \varphi_r)} \right]^2 - 1}. \quad (11)$$

According to Eqs.(10) and (11), scattering angle mainly depends on changing ξ , which embodies the change of $\psi_1 + \psi_2$ from ($\psi_1 = \theta_t - \varphi_t$, $\psi_2 = \theta_r - \varphi_r$) to ($\psi_1 = \theta_t + \varphi_t$, $\psi_2 = \theta_r + \varphi_r$). At the same time, η_1 and η_2 are also interfered.

3.3.3 ϕ_1 and ϕ_2

Transmitting cone of light is of symmetry with receiving cone according to $X-Z$ plane, thus, the overlapping volume is of symmetry according to $X-Z$ plane.

$$\phi_2 = -\phi_1, \quad (12)$$

$$\phi_2 = \arctan[r_1 \sin(\theta_t), r_1 \tan(\varphi_t)] = -\phi_1. \quad (13)$$

4 Conclusion

Aiming at phase function at different position within overlapping volume, phase function with different scattering angles dropping in receiving range is integrated, and integration limit is analyzed and deduced, which gives receiving power more correct quantity.

参 考 文 献

- 1 R M Gagliardi, S Karp. Optical Communication [M]. Beijing: People Telecommunication Press, 1982: 1.
- 2 Hongwei Yin, Shengli Chang, Honghui Jia, *et al.*. Non-line-of-sight multiscatter propagation model [J]. Journal of Optical Society of America, 2009, 26(11): 2466-2469.
- 3 Ke Xizheng. Ultraviolet Self-Organizing Network Theory [M]. Beijing: Science Press, 2011: 6-9.
柯熙政. 紫外光自组织网络理论[M]. 北京: 科学出版社, 2011: 6-9.
- 4 Zhao Taifei, Zhang Aili, Jin Dan, *et al.*. Research on the inter-link interference model in wireless ultraviolet non-line-of-sight communication [J]. Acta Optica Sinica, 2013, 33(7): 0706023.
赵太飞, 张爱丽, 金 丹, 等. 无线紫外光非视距通信中链路间干扰模型研究[J]. 光学学报, 2013, 33(7): 0706023.
- 5 Zhao Taifei, Feng Yanling, Ke Xizheng, *et al.*. Research on the coverage area of communication in the solar-blind UV communication network [J]. Acta Optica Sinica, 2010, 30(8): 2229-2235.
赵太飞, 冯艳玲, 柯熙政, 等. “日盲”紫外光通信网络中节点覆盖范围研究[J]. 光学学报, 2010, 30(8): 2229-2235.
- 6 Zhu Mengzhen, Zhang Hailiang, Jia Honghui, *et al.*. Study of ultraviolet scattering phase function based on Mie scattering theory [J]. J Light Scattering, 2007, 19(3): 225-228.
- 7 Dominique Toubanc. Henyey-Greenstein and Mie phase functions in Monte Carlo radiative transfer computations [J]. Appl Opt, 1996, 35(18): 3270-3274.
- 8 Mark R Luetgen, Jeffrey H Shapiro, David M Reilly. Non-line-of-sight single-scatter propagation model [J]. Journal of Optical Society of America A, 1991, 8(12): 1964-1972.

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