

激光空心 TIG 同轴复合电弧特性数值模拟

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摘要 利用 GAMBIT 软件建立了光纤连续激光和内径为 2 mm 空心 Tungsten Inert Gas(TIG)电弧在稳态条件下的激光同轴复合电弧数学模型,对其进行分区域网格划分,最小网格长度为 0.1 mm;用 FLUENT 软件中的用户自定义函数(UDF)功能添加了激光热源、动量方程源项、能量方程源项和材料物理性能设置,计算得到了 150 W 功率激光和 118 A 电流下同轴复合电弧的温度场、速度场及电弧空间压力、电场和磁场分布,并与相同条件下的空心单 TIG 电弧作了对比。结果表明:激光同轴复合电弧温度场变化明显,最高温度和高温区面积显著增加,电弧能量密度显著提高,熔透能力大大增强;由于激光同轴复合电弧温度升高和钨极腔内压力增大,电弧中上部等离子体流动速度明显加快;激光同轴复合电弧的电弧压力降低,电磁场强度减小,电弧静压力减小。最后用实验验证了模拟结果的可靠性。

关键词 激光技术; 同轴复合焊; 激光电弧复合焊; 空心 TIG 电弧; 电弧特性; 电弧模拟

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1 引言

激光电弧复合焊以焊接效率高、接头质量好等优点受到越来越多的关注,尤其是激光电弧旁轴复合焊,由于复合方式容易实现而被广泛关注和研究^[1-3]。而激光电弧同轴复合因为焊枪头加工难度大研究较少,但同轴复合焊枪头结构紧凑,在复杂狭小空间完成异形焊缝焊接的优势明显。20 世纪陈彦斌等^[4]就指出空心 Tungsten Inert Gas(TIG)电弧中心电流密度低,有利于激光穿过,同轴复合焊接熔深会大于旁轴复合。Doi^[5]利用功率为 4 kW 的 YAG 激光对内径为 4 mm 的空心 TIG 电弧进行同轴复合实验,熔深达到单激光焊的 2 倍,原因是同轴复合匙孔正上方的电弧压力扩张了匙孔的开口半径,电弧放电效率得到提高且匙孔变得稳定。Cho 等^[6]计算 YAG 激光沿 TIG 电弧轴线穿过时等离子体对激光的逆韧致吸收,发现每 1 cm 长电弧的等离子体只吸收 0.2% 的激光能量,可以忽略等离子体对 YAG 激光传输的影响;研究发现随着激光功率升高,复合等离子体温度升高,当激光功率超过

400 W 时,工件阳极上方等离子体流动并出现漩涡,改变了熔池表面剪切力分布。Startsev 等^[7]在不考虑光致等离子体的影响下,研究发现 CO₂ 激光垂直穿过 TIG 电弧时,电弧中心温度明显升高。丁玲芳^[8]模拟了 YAG 激光和实心 TIG 同轴复合电弧,发现激光照射点上方电流密度和温度大幅升高,而阴极温度和电流密度变化不大,阳极表面上方等离子体同样出现涡流特征。Yan 等^[9]研究了电弧形态为钟罩形的空心钨极,发现电弧弧柱区的压力分布随弧长的变化规律跟传统 TIG 弧变化规律类似,且由于工件表面是圆环,电弧顶部中心的压力变得比较均匀。

国内外学者关于激光和 TIG 电弧同轴复合的模拟研究主要集中在激光和传统实心 TIG 电弧的模拟,然而要想实现真正意义上的同轴复合,让激光穿过空心钨极电弧同轴复合是较好的方式。激光穿过空心钨极电弧到达阳极表面,必然引起材料蒸发和电离,产生激光等离子体。激光等离子体与电弧等离子体作用,改变电弧空间温度分布和离子流动状态,影响电弧在阳极表面的压力分布等,但是有关

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激光和空心 TIG 电弧复合后电弧特性的研究未见报道。因此,本文利用 FLUENT 软件对激光空心钨极同轴复合电弧等离子体进行了 Computational Fluid Dynamics(CFD)计算,得到了电弧温度场、等离子体流动状态及电弧压力分布,并用实验对电弧形态和电弧熔透能力进行了验证,为激光电弧同轴复合焊接技术开发提供了理论依据。

2 电弧模型和假设条件

2.1 物理模型和基本假设

空心钨极和激光同轴复合电弧数学模型如图 1 所示,焊接电流为直流 118 A,弧长为 4 mm,电弧满足轴对称条件,计算选择二维模型。空心钨极

直径为 8 mm,空腔为圆锥形,电极尖端内腔直径为 2 mm,尖端呈圆台状,凸台厚度为 0.2 mm,锥角为 76°。空心钨极内气体和外保护气均为高纯氩气,气流速度相同。激光为光纤激光,波长为 1060 nm,激光束有效半径为 0.6 mm。工件阳极材料选用 Q345 钢,激光的吸收率设为 0.3^[6]。对激光同轴复合电弧物理场的计算进行了如下假设:1)电弧为稳态;2)不考虑电弧对激光传输的影响,激光能量在阳极表面呈面热源分布;3)电弧空间等离子体局部处于热平衡状态;4)电弧等离子体流动状态为层流;5)电弧为光学厚度小的介质;6)氩气各项物理性能只是温度的函数;7)钨极尖端电流密度均匀分布。

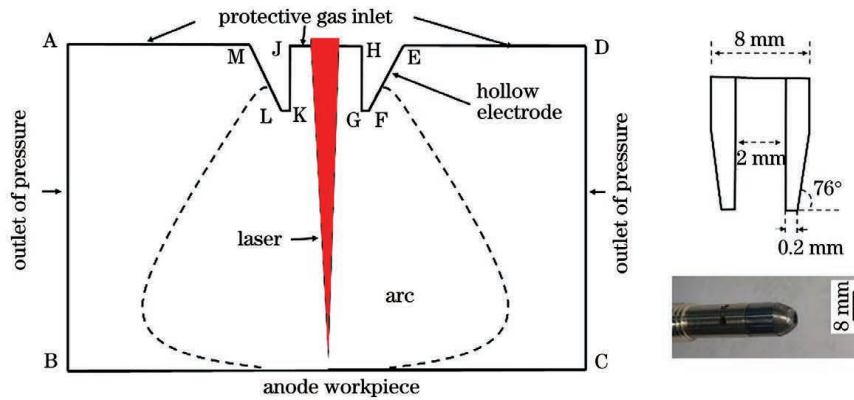


图 1 激光同轴复合电弧数学模型示意图

Fig. 1 Schematic diagram of mathematical model of laser coaxial composite arc

2.2 控制方程

质量连续方程为

$$\frac{\partial \rho}{\partial t} + \frac{1}{r} \frac{\partial (r\rho u)}{\partial r} + \frac{\partial (\rho v)}{\partial z} = 0. \quad (1)$$

径向动量方程为

$$\frac{\partial (\rho v)}{\partial t} + \frac{1}{r} \frac{\partial (r\rho v u)}{\partial r} + \frac{\partial (\rho u v)}{\partial z} = F_r - \frac{\partial P}{\partial r} + \frac{1}{r} \frac{\partial}{\partial r} \left(2\mu r \frac{\partial u}{\partial r} \right) - 2\mu \frac{v}{r^2} + \frac{\partial}{\partial z} \left(\mu \frac{\partial v}{\partial z} + \mu \frac{\partial u}{\partial r} \right). \quad (2)$$

轴向动量方程为

$$\frac{\partial (\rho u)}{\partial t} + \frac{1}{r} \frac{\partial (r\rho u v)}{\partial r} + \frac{\partial (\rho u u)}{\partial z} = F_z - \frac{\partial P}{\partial z} + \frac{1}{r} \frac{\partial}{\partial r} \left(\mu r \frac{\partial v}{\partial z} + \mu r \frac{\partial u}{\partial r} \right) + \frac{\partial}{\partial z} \left(2\mu \frac{\partial u}{\partial z} \right). \quad (3)$$

能量方程为

$$\frac{\partial (\rho c_p T)}{\partial t} + \frac{1}{r} \frac{\partial (r\rho v c_p T)}{\partial r} + \frac{\partial (\rho u c_p T)}{\partial z} = \frac{1}{r} \frac{\partial}{\partial r} \left(k r \frac{\partial T}{\partial r} \right) + \frac{\partial}{\partial z} \left(k \frac{\partial T}{\partial z} \right) + Q, \quad (4)$$

式中: v 为径向速度; u 为轴向速度; T 为温度; P 为压力; ρ 为密度; μ 为粘度; c_p 为比热容; k 为导热系数;体积力在 r 方向的分量 $F_r = (J \times B)_r$, 体积力在 z 方向的分量 $F_z = (J \times B)_z + \rho g$, 其中 J 为电流密度, B 为磁感应强度, g 为重力加速度。

麦克斯韦方程组如下。电流连续方程为

$$\frac{\partial}{\partial z} \left(\sigma \frac{\partial \varphi}{\partial z} \right) + \frac{1}{r} \frac{\partial}{\partial r} \left(\sigma r \frac{\partial \varphi}{\partial r} \right) = 0. \quad (5)$$

欧姆定律为

$$J_r = -\sigma \frac{\partial \varphi}{\partial r}, J_z = -\sigma \frac{\partial \varphi}{\partial z}. \quad (6)$$

安培环流定理为

$$B = \frac{\mu_0}{r} \int_0^r J_z r dr, \quad (7)$$

式中： μ_0 为真空磁导率； σ 为电导率； φ 为电势。

激光热源模型为

$$q_{\text{laser}} = \frac{3\eta M}{\pi r_L^2} \exp\left(-\frac{3r^2}{r_L^2}\right), \quad (8)$$

表 1 空心 TIG 电弧模型边界条件

Table 1 Boundary condition of hollow TIG arc model

Area	Boundary type	Velocity $V / (\text{m} \cdot \text{s}^{-1})$	Temperature T / K	Potential φ / V	Magnetic vector $A / (\text{Wb} \cdot \text{m}^{-1})$
BC	Wall	0	5000	0	$\partial A / \partial z = \partial A / \partial r = 0$
AB, CD	Pressure-outlet		1000	$\partial \varphi / \partial z = \partial \varphi / \partial r = 0$	0
DE, HJ, MA	Velocity-inlet	2	1000	$\partial \varphi / \partial z = \partial \varphi / \partial r = 0$	$\partial A / \partial z = \partial A / \partial r = 0$
EF, LM	Wall	0	1000	$\partial \varphi / \partial z = \partial \varphi / \partial r = 0$	$\partial A / \partial z = \partial A / \partial r = 0$
FG, KL	Wall	0	3000	$-\sigma \partial \varphi / \partial z = I / S_c$	$\partial A / \partial z = \partial A / \partial r = 0$
GH, JK	Wall	0	1000	$\partial \varphi / \partial z = \partial \varphi / \partial r = 0$	$\partial A / \partial z = \partial A / \partial r = 0$

动量方程和能量方程的源项加载使用用户自定义函数(UDF)进行程序编写并加载,电弧辐射散热用 S_R 表示^[10],能量方程源项表达式为

$$Q = \frac{J_z^2 + J_r^2}{\sigma} + \frac{5k_B}{2e} \left(J_z \frac{\partial T}{\partial z} + J_r \frac{\partial T}{\partial r} \right) - S_R, \quad (9)$$

式中： J_r 和 J_z 分别为电流密度在径向和轴向的分量； k_B 为玻尔兹曼常数； e 为电子电荷量。

3 结果与讨论

3.1 温度场

激光同轴复合电弧温度场云图如图 2(a)所示,电弧呈典型的钟罩形状,激光照射点上方和电极附近出现了高温区域,相比空心 TIG 电弧温度云图 2(b),高温区的温度明显升高,图 2(b)模拟结果与文献[9]一致,说明结果的正确性;同轴复合电弧上方中间温度下凹,电弧中间高温区形状呈心形;同轴复合电弧上下宽度均略大于空心 TIG 电弧;激光

的加入使钨极温度升高,发射电子能量增强,复合电弧上部变宽,电弧高度上升,电弧出现“上爬”;电弧温度跟电流密度有关,阳极材料在激光照射下产生的金属蒸气被电离成电子和正离子,电子密度增大,电流密度升高,电弧温度升高;高密度电子的流动换热是激光电弧复合焊阳极工件温度升高的主要机制^[6]。图 3 是同轴复合电弧和空心 TIG 电弧温度对比曲线,同轴复合电弧最高温度 19240 K 出现在电极附近阴极区,激光照射点上方阳极区温度达到 18737 K,空心 TIG 电弧温度峰值为 15512 K;激光的加入使得同轴复合电弧最高温度升高了 24%;从同轴复合电弧中轴线温度曲线可以看到,电弧中心温度从电极平面到电极下方 1.5 mm 一直增加到 13228 K,然后到 3 mm 区间稍有减小到 12550 K,最后一直到工件表面快速增加,可见电弧上部中心温度凹陷 1.5 mm;相比之下空心 TIG 电弧中心轴线温度也从电极平面到 1.5 mm 位置先增加到

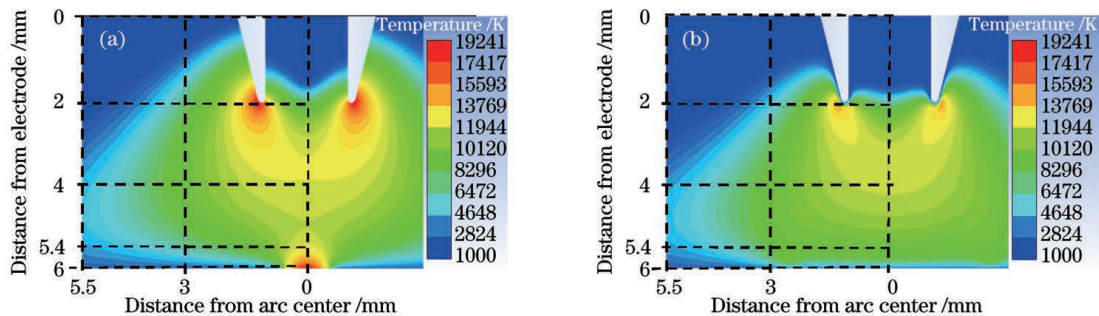


图 2 电弧温度分布云图。(a)激光同轴复合电弧;(b)TIG 电弧

Fig. 2 Cloud diagram of arc temperature distribution. (a) Laser coaxial composite arc; (b) TIG arc

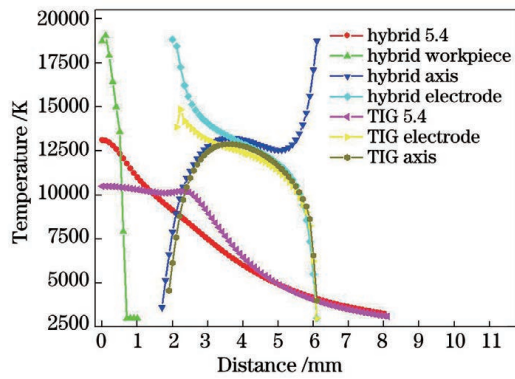


图 3 电弧温度对比

Fig. 3 Comparison of arc temperature

12900 K, 随后 1.5 mm 到 2.6 mm 区间温度稍有降低, 到 12127 K, 随后温度连续降低, 并且降温速度越来越快。两电弧中上部温度下凹原因是空心钨极圆环放电, 环中心电流密度小, 而钨极尖端即圆环电流密度大, 电流密度差异大造成了中心温度偏低; 随着电子向阳极移动, 到电弧中部中心区域电流密度趋于均匀, 因此形成了电弧中部的高温等温区。Jiang 等^[11]把电弧从中心向外分为中心原子电离的高温区、原子和电子复合的中温区及外围低温区 3 个区域。从图 2 可以看到, 复合电弧的高温区宽度和长度均大于空心 TIG 电弧, 激光的加入使得电

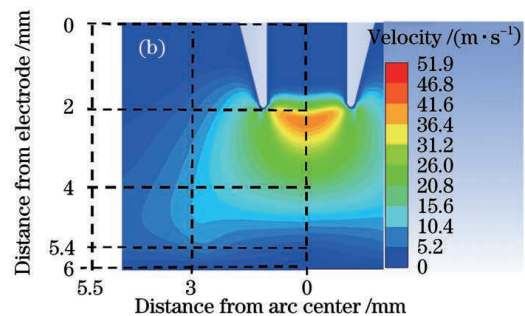
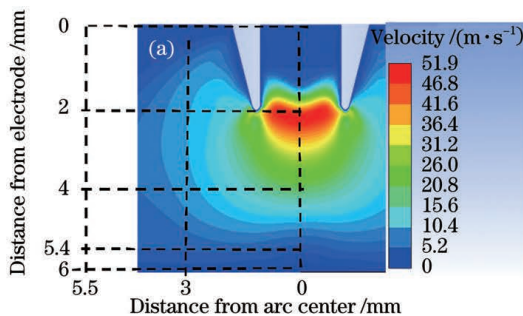


图 4 等离子体速度分布云图。(a)激光同轴复合电弧;(b)TIG 电弧

Fig. 4 Cloud diagram of plasma velocity distribution. (a) Laser coaxial composite arc; (b) TIG arc

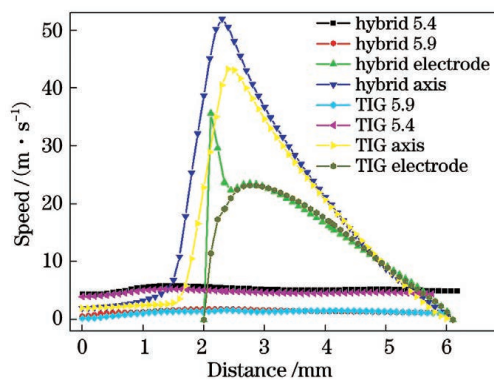


图 5 等离子体速度对比

Fig. 5 Comparison of plasma velocity

弧中心高温区温度和面积均增加; 从电弧 3.4 mm 处径向温度曲线可以看到, 复合电弧中心温度明显高于 TIG 电弧, 但是当距离中轴线超过 1.3 mm, 复合电弧温度低于空心 TIG 电弧。

3.2 速度场

电弧空间等离子体在力和热的作用下做高速运动。激光同轴复合电弧等离子体速度分布云图如图 4(a)所示, 可以看到空心阴极出口处离子流动速度较快, 最高速度为 52 m/s, 在出口到 1.3 mm 运动速度维持在 30 m/s 以上, 随后速度明显减慢, 2 mm 位置速度降为 20 m/s。空心 TIG 电弧速度场云图如图 4(b)所示, 等离子体运动趋势与同轴复合电弧相同, 但是速度减小, 速度峰值为 43 m/s, 同轴复合电弧等离子体最大速度相比空心 TIG 电弧提高了 21%。温差越大, 等离子体运动速度也越快, 由图 2 可知同轴复合电弧阴极区温度大于空心 TIG 电弧, 因此复合电弧离子流动速度也更快; 另一方面, 激光使得空心钨极腔内压力升高, 保护气体和粒子运动速度加快。图 5 是电弧空间离子运动速度对比。从轴线速度曲线可以看到, 粒子在钨极腔内的加速度相同, 但是由于同轴复合电弧上爬, 粒子加速的时间更早, 在腔内加速时间更长, 结果到达钨

极出口的速度更高; 从同轴复合电弧电极下方离子速度曲线可以看到, 激光的加入使得电极下方离子运动速度先迅速增大, 随后减小, 到电极下方 0.4 mm 处速度与空心 TIG 电弧离子相同; 同轴复合电弧与空心 TIG 电弧等离子体在工件表面运动速度相近。

3.3 压力场

电弧压力主要源于电磁力产生的静压力和等离子体流动产生的动压力。同轴复合电弧和钨极腔内压力分布云图如图 6(a)所示, 钨极腔内压力较大, 电弧空间压力分布比较均匀, 激光照射点周围压力

略有升高。空心 TIG 电弧压力如图 6(b) 所示, 相比之下二者压力分布规律近似, 但同轴复合电弧在阳极表面的压力更低。电弧空间压力对比如图 7 所示。从轴线压力分布曲线可以看到, 同轴复合电弧压力低于空心 TIG 电弧, 但是激光使得钨极温度升高, 腔内压力更大; 从工件表面压力分布曲线可以看到, 同轴复合电弧压力比较均匀, 中心激光照射点压力最大值为 22 Pa, 低于空心 TIG 电弧的压力; TIG 电弧压力在工件表面中心呈矩形均匀分布, 最大值为 32 Pa; 激光的加入使空心 TIG 电弧阳极表面压

力降低了 31%; 电弧在工件表面的压力源于等离子流压力和电磁收缩力, 离子运动速度越快, 遇到工件产生的动压力越大。由图 5 可以看到, 靠近阳极表面, 离子运动速率相近, 因此等离子流产生的动压力也接近; 电磁收缩力跟电弧形状有关, 从图 2 可以看到, 同轴复合电弧上部明显变宽, 电弧下部和上部宽度比值减小, 上下产生的压力差值减小, 电磁收缩力减小; 空心 TIG 电弧压力比传统实心 TIG 低, 有利于抑制薄板焊接烧穿缺陷的出现, 同轴复合电弧压力降低可以进一步提升焊接质量。

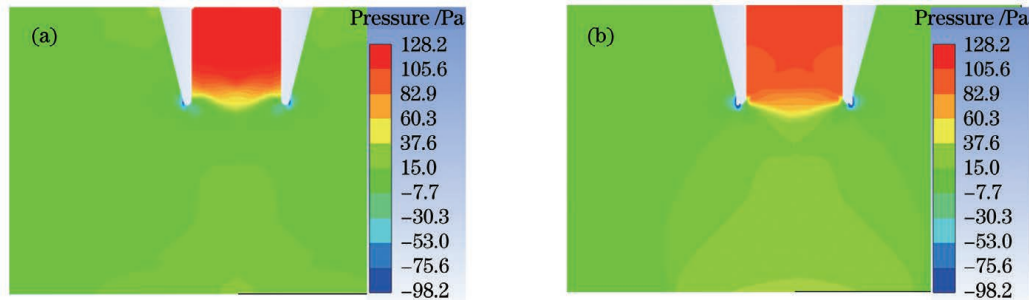


图 6 电弧压力分布云图。(a)激光同轴复合电弧;(b)TIG 电弧

Fig. 6 Cloud diagram of arc pressure distribution. (a) Laser coaxial composite arc; (b) TIG arc

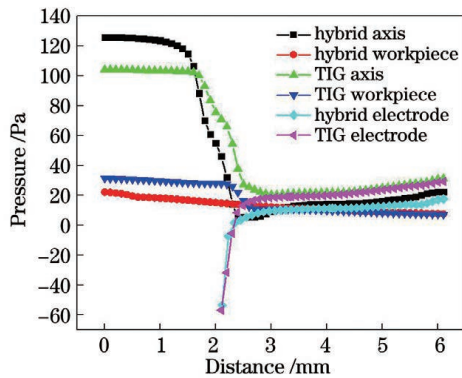


图 7 电弧压力对比

Fig. 7 Comparison of arc pressure

3.4 电场

同轴复合电弧电场分布云图如图 8(a) 所示。由于工件阳极电势设为 0, 从阴极到阳极电势逐渐

升高, 因此电弧空间电势为负值。可以看到, 由于阴极周围电流密度高, 电场强度较大, 离阴极越远, 场强越小; 阳极激光照射点周围电子密度增大, 电势升高。图 9 是同轴复合电弧与空心 TIG 电弧空间电势对比图, 可以看到, 同轴复合电弧空间电场强度更小; 从电极下方电势分布曲线可以看到, 同轴复合电弧和空心 TIG 电弧的阴极区电压降都非常明显, 复合电弧的弧柱区电势变化很小, 空心 TIG 电弧弧柱区电势变化稍大; 从电弧轴线电势分布可以看到, 同轴复合电弧从阴极腔内的 -2.16 V 逐渐升高到阴极区的 -1.7 V , 再到阳极面上方 0.6 mm 位置的 -0.6 V , 而空心 TIG 电弧阴极区电势为 -6 V , 阳极面上方 0.6 mm 位置电势为 -4.2 V , 同轴复合电弧阴极区电场强度降低了 72%, 在阳极表面降低

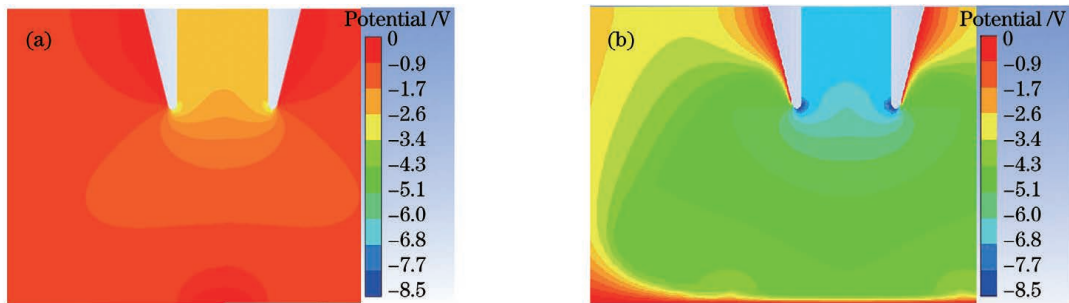


图 8 电弧电势分布云图。(a)激光同轴复合电弧;(b)TIG 电弧

Fig. 8 Cloud diagram of arc potential distribution. (a) Laser coaxial composite arc; (b) TIG arc

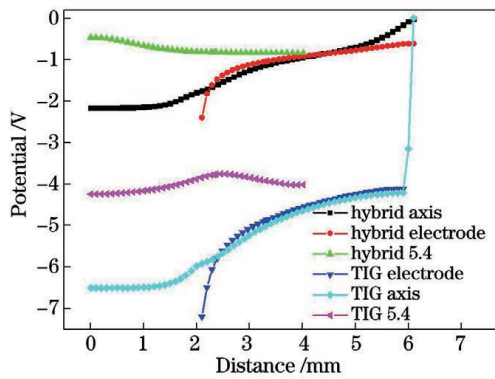


图 9 电弧电势对比

Fig. 9 Comparison of arc potential

了 90%；激光作用位置电势从周围的 -0.6 V 升高到了 -0.4 V。

3.5 磁场

电子流从阴极向阳极移动,形成了载流导体并在周围产生了磁场。同轴复合电弧空间轴向磁场分布云图如图 10(a)所示,阴极尖端电流密度高,磁场强度也较大,而空心阴极中心电流密度小,电弧顶部中心磁场强度低;从电弧阴极区向阳极区,磁场强度逐渐减小,电弧中部中心区域磁场均匀,激光照射点

磁场强度降低。图 10(b)是空心 TIG 电弧轴向磁场云图,同样阴极尖端下方磁场强度高,电弧中部和下部磁场强度差异不大。激光的加入使得复合电弧轴向磁场强度低于空心 TIG 电弧,轴向磁场强度梯度增大。图 10(c)是同轴复合电弧径向磁场分布云图,电弧顶部磁场中心低而阴极尖端周围高,中心区域磁场均匀;激光照射点磁场强度降低。图 10(d)是空心 TIG 电弧径向磁场云图,空心阴极出口下方有一个低磁场强度区域,电弧中部和下部中心区域磁场比较均匀。激光同轴复合电弧与空心 TIG 电弧的径向磁场方向相反,复合电弧径向磁场强度梯度增大。图 11 是电弧空间磁场强度分布曲线。同轴复合电弧纵向磁场强度最大值为 5.1262 mT,激光照射点磁场强度最低为 5.1228 mT;横向磁场强度最大值为 1.879 mT,工件表面激光照射点磁场强度最低为 1.8741 mT;空心 TIG 电弧纵向磁场强度最大值为 30.3804 mT,工件表面磁场强度最大值为 30.3798 mT;横向磁场强度最大值为 65.698 mT。激光同轴复合电弧的纵向磁场强度最大值和工件表面磁场强度相比空心 TIG 电弧均降低了 83%,横向磁场方向相反,磁场强度降低 97%。

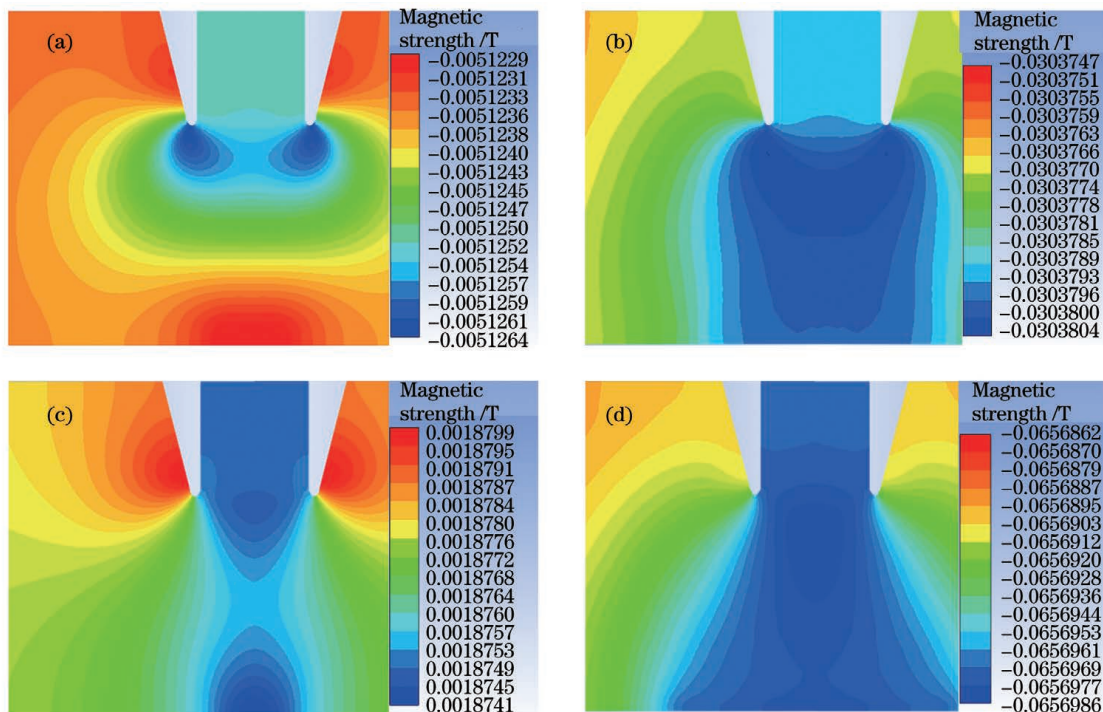


图 10 电弧磁场分布云图。(a)复合电弧纵向;(b)TIG 电弧纵向;(c)复合电弧横向;(d)TIG 电弧横向

Fig. 10 Cloud diagram of arc magnetic field distribution. (a) Horizontal of composite arc; (b) horizontal of TIG arc; (c) vertical of composite arc; (d) vertical of TIG arc

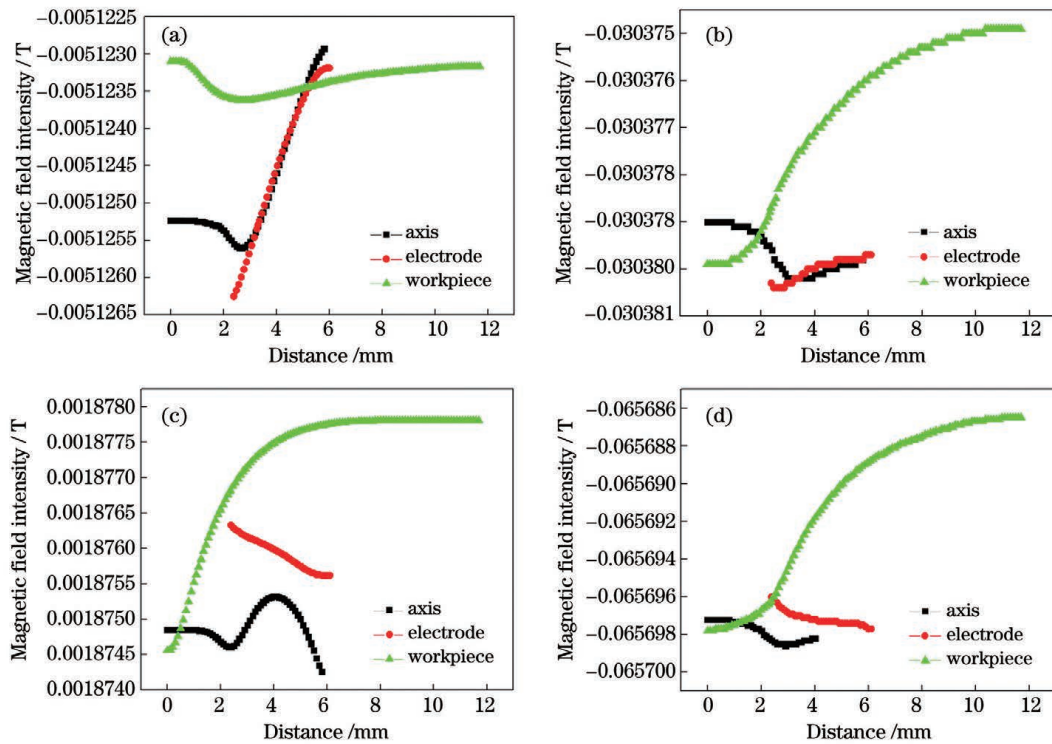


图 11 磁场对比。(a)复合电弧纵向;(b)TIG 纵向;(c)复合电弧横向;(d)TIG 横向

Fig. 11 Comparison of magnetic field. (a) Horizontal of composite arc; (b) horizontal of TIG arc; (c) vertical of composite arc; (d) vertical of TIG arc

4 实验验证

图 12(a)是实际焊接实验中激光同轴复合电弧形态,钨极离工件高度 4 mm,实际焊接电弧和模拟电弧均呈典型钟罩形状。图 12(b)是相同电流下空心 TIG 电弧,复合电弧宽度略大于空心 TIG 电弧。图 13(a)是同轴复合电弧在 Q345 钢表面堆焊焊缝截面形貌,熔深为 0.8 mm,熔宽为 3.8 mm,焊缝宽度为 4.3 mm。图 13(b)是空心 TIG 电弧堆焊焊缝截面,熔深为 0.5 mm,熔宽为 3.4 mm,焊缝宽度为 4.0 mm。激光同轴复合电弧与空心 TIG 电弧相比,堆焊焊缝熔宽增加 12%,焊缝宽度增加

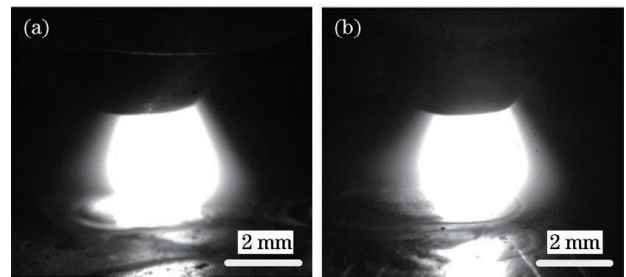


图 12 实验电弧。(a)激光同轴复合电弧;(b)TIG 电弧

Fig. 12 Arc in experiment. (a) Laser coaxial composite arc; (b) TIG arc

7%,熔深增加 60%;激光同轴复合焊熔深增加与同轴复合电弧温度升高、能量密度增大有关。

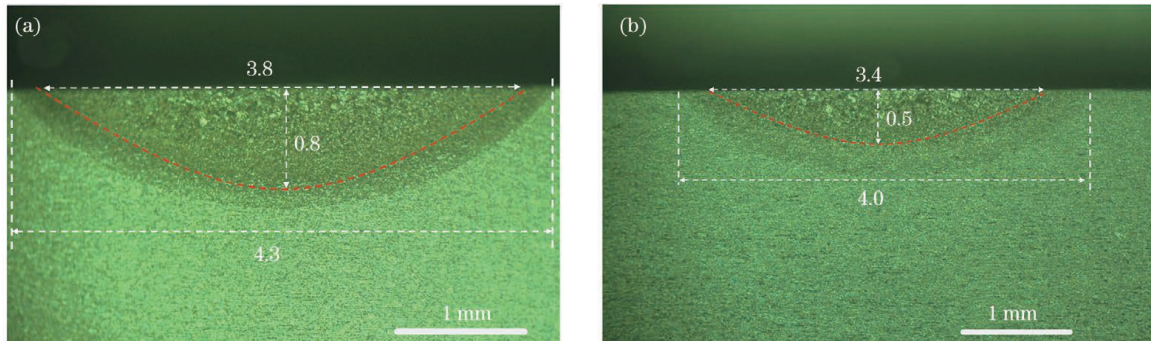


图 13 实验焊缝。(a)激光同轴复合电弧;(b)TIG 电弧

Fig. 13 Welding joint in experiment. (a) Laser coaxial composite arc; (b) TIG arc

5 结 论

运用 FLUENT 软件对激光和空心 TIG 电弧同轴复合进行模拟。激光同轴复合电弧峰值温度和高温区面积相比空心 TIG 电弧均明显增大,峰值温度出现在阴极尖端,温度升高了 24%,电弧上部宽度明显增大。激光同轴复合电弧中心轴线温度从阴极区到阳极区先升高后保持不变然后再升高,顶部中心有 1.5 mm 下凹。激光同轴复合电弧中上部等离子体运动速度相比空心 TIG 电弧明显加快,最大速度提高了 21%,电弧温度升高和钨极腔内压力增大是速度加快的主要原因,阳极表面等离子体运动速度变化不大。激光同轴复合电弧压力相比空心 TIG 电弧减小,在工作阳极表面的压力降低了 31%,激光照射斑点处压力升高。激光同轴复合电弧空间电磁场分布比较均匀,相比空心 TIG 电弧电磁场强度减小,横向磁场方向发生反向。激光同轴复合电弧相比空心 TIG 电弧熔透能力增强,相同电流条件下,焊缝熔深增加 60%,熔宽增加 12%,焊缝宽度增加 7%。

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Numerical Simulation of Characteristics of Laser-Hollow Tungsten Inert Gas Coaxial Composite Welding Arc

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Abstract

Objective Laser-arc hybrid welding technology has advantages of high welding efficiency and good joint quality. However, it is currently based on the paraxial composite, which has high requirements for welding direction, and the keyhole instability causes welding defects and other problems. Laser-arc coaxial composite welding has advantages of good heat source symmetry, stable welding process, high joint quality, and good reachability of the welding torch joint. However, research on laser-arc coaxial composite welding is limited because of the difficulty of welding torch head processing. There are two kinds of coaxial recombination, and the structure of passing a laser through a hollow tungsten electrode is relatively simple and easier to achieve than splitting the laser into two beams and then converging them at the focus. However, research on the composite arc characteristics of lasers and hollow tungsten inert gas (TIG) arc has not been reported. Therefore, the study on the characteristics of laser coaxial composite arc in this paper can provide theoretical support for technology development.

Methods The mathematical model of laser coaxial composite arc of a fiber continuous laser and 2 mm hollow TIG welding arc under steady-state conditions was established using GAMBIT software. The regional mesh was divided and the minimum mesh was 0.1 mm. The laser heat source, momentum source term, and energy source term were added, and reasonable material physical property parameters were set using user-defined functions (UDF) FLUENT software. The temperature field, velocity field, spatial pressure, electric field, and magnetic field distribution of the coaxial composite arc under 118-A current and 150-W laser power were calculated and compared with those of the hollow single TIG arc under the same conditions. Finally, experiments verify the penetration ability of laser coaxial composite arc.

Results and Discussions The calculation results show that compared with the hollow TIG arc, the maximum temperature and high-temperature area of the laser coaxial composite arc increase significantly. The arc's upper part becomes wider and the arc height increases. The middle temperature above the arc is concave, and the arc's middle high-temperature area is heart-shaped (Fig. 2). The laser coaxial composite arc's maximum temperature is 19240 K in the cathode region near the electrode, the temperature of the laser irradiation point reaches 18737 K, and the hollow TIG arc's peak temperature is 15512 K (Fig 3). The flow peak velocity of laser coaxial composite arc plasma is 52 m/s, whereas that of hollow TIG arc is 43 m/s. The plasma flow velocity in the middle and upper parts of the laser coaxial composite arc is faster than that of the hollow TIG arc, but the surface velocity of the workpiece changes little (Fig. 4 and Fig. 5). The maximum pressure of the laser coaxial composite arc on the workpiece surface is 22 Pa, whereas that of the hollow TIG arc is 32 Pa (Fig. 6 and Fig. 7). The pressure of the laser irradiation spot increases. The temperature of the hollow tungsten electrode and pressure in the cavity increase under the action of laser; therefore, the flow of protective gas and plasma flow in the upper part of the laser coaxial composite arc speed up. The reduction in arc pressure can suppress welding defects, such as burning through and gnawing, but the arc stiffness might be reduced. The maximum voltage of the laser coaxial composite arc is -1.7 V, and that of the hollow TIG arc is -6 V. The electric field intensity is reduced by 72% (Fig. 8 and Fig. 9). The longitudinal magnetic field intensity of the laser coaxial composite arc decreases from 30 mT to 5 mT in the hollow TIG arc, and the transverse magnetic field decreases from 65.6 mT to 1.87 mT, and the transverse magnetic field changes direction (Fig. 10 and Fig. 11). Under the same current conditions, the surfacing welding penetration of laser coaxial composite arc and hollow TIG arc on Q345 steel is 0.8 mm and 0.5 mm and the widths are 4.3 mm and 4.0 mm, respectively (Fig. 13).

Conclusions In this study, FLUENT software was used to simulate the coaxial recombination of laser and hollow TIG arc. Compared with the hollow TIG arc, the peak temperature of the laser coaxial composite arc is increased by

24%, the width of the upper arc is increased, and the top center has a concave of 1.5 mm. The plasma velocity in the middle and upper parts of the laser coaxial composite arc is significantly faster, and the maximum velocity is increased by 21%. The increase in arc temperature and the pressure in the tungsten cavity is the main reason for plasma velocity acceleration, whereas the surface plasma velocity in the anode changes little. The pressure of the laser coaxial composite arc decreases by 31% on the anodic surface of the workpiece, and the pressure at the spots exposed by the laser increases. The electromagnetic field intensity of the laser coaxial composite arc decreases, and the transverse magnetic field direction is reversed. Compared with hollow TIG arc, laser coaxial composite arc has a better penetration ability under the same current conditions, welding depth, and melt width, and welding width increases by 60%, 12%, and 7%.

Key words laser technique; coaxial hybrid welding; laser-arc hybrid welding; hollow TIG arc; arc characteristic; arc simulation

OCIS codes 140.3390; 350.3850; 350.3390