

类钠锗离子软 X 射线光谱的理论计算

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摘要 本文将原子结构的 HFR 方法与关于能量参数的最小二乘优化方法相结合, 详细计算了类钠锗离子 $1s^2 2s^2 2p^6 nl$ 组态 ($n = 3 \sim 6, l = 0 \sim 5$) 电偶极跃迁的软 X 射线光谱 (波长 $< 30 \text{ nm}$)。本文的波长计算值在误差范围内与实验十分符合, 振子强度与 Wiese 等人的结果一致。

关键词 软 X 射线光谱, 类钠锗离子

A numerical computation of the soft X-ray spectra of Na-like Ge ion

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Abstract The soft X-ray spectra of Na-like Ge ion of $1s^2 2s^2 2p^6 nl$ ($n = 3 \sim 6, l = 0 \sim 5$) configurations for dipole-allowed transitions are computed with the aid of the HFR program with the optimization of Slater parameters on the basis of the known experimental levels. The excellent agreement between the calculated wavelengths with the observed ones is obtained and these results are discussed.

Key words soft X-ray spectra, Na-like Ge ion

1 引言

在复合机制软 X 射线激光研究方面, 类钠离子由于具有与类锂离子相似的电子结构(闭合壳层之外带一个价电子), 其电离能比类 H、类 He 离子的电离能低很多, 因而也同样具有低泵浦、高效率等优点; 而且可以比类 H 离子 ($\lambda \propto Z^{-2}$) 和类 Li 离子 ($\lambda \propto Z^{-2.5}$) 更快地推向短波长 ($\lambda \propto Z^{-2.9}$), 因此对类钠离子的研究已得到越来越多的注意^[1~3]。对类钠铜离子初步软 X 射线激光研究表明^[4]: 利用复合机制得到自发辐射放大是有希望的, 沿类钠系列将软 X 射线激光波长推至“水窗”波段是完全有可能的。因此, 理论上提供类 Na 离子软 X 射线光谱数据对这个系列的软 X 激光研究具有重要意义。

利用类 Na 锗离子可能得到比类钠铜离子波长更短的软 X 射线激光。关于类钠锗离子的光谱已有几个实验及理论研究^[5~7]。但主要集中于 $3l \rightarrow n'l'$ ($n' \geq 4$) 及 $4l \rightarrow 5l'$ 的电偶极跃迁, 对高激发态之间的跃迁光谱研究很少, 而且缺乏振子强度、辐射跃迁几率等光谱数据的详细计算, 因而有必要进行进一步全面的计算以满足软 X 射线激光研究的需要。

本文应用以 HFR 方法为基础的原子结构和光谱计算程序^[8], 结合关于原子能量参数的最小二乘优化方法, 计算了类钠锗离子 $1s^2 2s^2 2p^6 nl$ ($n = 3 \sim 6, l = 0 \sim 5$) 组态电偶极跃迁的软 X

射线光谱。我们的波长计算值在误差范围内与实验观测值十分符合,振子强度也与文献[9]的结果一致。

2 计算方法

本文计算采用的程序是 R. D. Cowan 博士的原子结构和光谱计算程序库 RCN34/RCN2/RCG9/RCE19^[8],其中 RCN34 计算原子的单电子轨函数和 Slater 积分即 E_{av} 、 F^K 、 G^K 和 ξ_i ,RCN2 计算组态相互作用积分 R^K 和电偶极跃迁积分,RCG9 计算原子能级本征值,本征波函数及波长,振子强度和辐射跃迁几率等;在已知部分实验能级值的情况下^[5],RCE19 用最小二乘法优化方法调整 Slater 径向积分(能量参数),使能级计算值与实验值之差达最小。运用优化后的 Slater 径向积分重复运行程序 RCG9,就得到最后的原子光谱数据。

有关 HFR 方法和最小二乘优化方法的详述可参见文献[10],这里不再赘述。

3 结果和分析

表 1 列出了类钠离子 $1s^2 2s^2 2p^6 nl$ ($n = 3 \sim 6$, $l = 0 \sim 5$) 电偶极跃迁的软 X 射线光谱计算结果。

Table 1 The soft X-ray spectra of Na-like Ge ion of $2P^6 nl$ configurations with dipole-allowed transitions (wavelengths in 0.1 nm)

(nl)	Transitions		$J - J'$	λ	λ_1	λ_2	gf	gA (s ⁻¹)
	${}^2L - (nl)$	2L		(calc.)	(obs.)	(obs.)		
(3s)	$2S - (3p)$	$2P$	$0.5 - 0.5$	261.515	261.52	261.500	0.198	$1.931E + 10$
(3s)	$2S - (3p)$	$2P$	$0.5 - 1.5$	226.505	226.505	226.497	0.457	$5.943E + 10$
(3d)	$2D - (3p)$	$2P$	$1.5 - 1.5$	194.437	194.435	194.435	0.082	$1.447E + 10$
(3d)	$2D - (3p)$	$2P$	$2.5 - 1.5$	190.614	190.614	190.608	0.753	$1.383E + 11$
(3d)	$2D - (3p)$	$2P$	$1.5 - 0.5$	174.395	174.396	174.391	0.457	$1.003E + 11$
(5d)	$2D - (6p)$	$2P$	$1.5 - 0.5$	162.129			0.393	$9.982E + 10$
(5d)	$2D - (6p)$	$2P$	$2.5 - 1.5$	161.055			0.712	$1.833E + 11$
(5d)	$2D - (6p)$	$2P$	$1.5 - 1.5$	160.440			0.079	$2.060E + 10$
(6d)	$2D - (5f)$	$2F$	$1.5 - 2.5$	160.041			0.234	$6.109E + 10$
(6d)	$2D - (5f)$	$2F$	$2.5 - 3.5$	159.138			0.337	$8.876E + 10$
(6d)	$2D - (5f)$	$2F$	$2.5 - 2.5$	158.933			0.016	$4.455E + 09$
(5g)	$2G - (6f)$	$2F$	$3.5 - 2.5$	154.935			0.063	$1.757E + 10$
(5g)	$2G - (6f)$	$2F$	$4.5 - 3.5$	154.892			0.082	$2.279E + 10$
(5g)	$2G - (6f)$	$2F$	$3.5 - 3.5$	154.743			0.002	$6.530E + 08$
(5g)	$2G - (6h)$	$2H$	$4.5 - 4.5$	153.306			0.305	$8.677E + 10$
(5g)	$2G - (6h)$	$2H$	$4.5 - 5.5$	153.256			6.516	$4.690E + 12$
(5g)	$2G - (6h)$	$2H$	$3.5 - 4.5$	153.161			3.465	$3.829E + 12$
(6g)	$2G - (5f)$	$2F$	$3.5 - 3.5$	152.079			0.260	$7.509E + 10$
(6g)	$2G - (5f)$	$2F$	$4.5 - 3.5$	152.005			9.118	$2.632E + 12$
(6g)	$2G - (5f)$	$2F$	$3.5 - 2.5$	151.892			7.039	$2.035E + 12$
(6s)	$2S - (5p)$	$2P$	$0.5 - 1.5$	149.854			0.561	$1.666E + 11$
(6s)	$2S - (5p)$	$2P$	$0.5 - 0.5$	147.489			0.285	$8.740E + 10$
(5d)	$2D - (6f)$	$2F$	$2.5 - 2.5$	143.455			0.193	$6.279E + 10$

Transitions				λ	λ_1	λ_2	gf	gA (s^{-1})
(nl)	$^2L - (nl)$	2L	$J - J$	(calc.)	(obs.)	(obs.)		
(5d)	$2D - (6f)$	$2F$	$2.5 - 3.5$	143.291			3.879	$1.260E + 12$
(5d)	$2D - (6f)$	$2F$	$1.5 - 2.5$	142.967			2.721	$8.881E + 11$
(6d)	$2D - (5p)$	$2P$	$1.5 - 1.5$	130.412			0.127	$5.005E + 10$
(6d)	$2D - (5p)$	$2P$	$2.5 - 1.5$	129.675			1.155	$4.582E + 11$
(6d)	$2D - (5p)$	$2P$	$1.5 - 0.5$	128.617			0.647	$2.609E + 11$
(5s)	$2S - (6p)$	$2P$	$0.5 - 0.5$	125.764			0.199	$8.430E + 10$
(5s)	$2S - (6p)$	$2P$	$0.5 - 1.5$	124.745			0.403	$1.728E + 11$
(4d)	$2D - (5p)$	$2P$	$1.5 - 0.5$	87.626			0.240	$2.087E + 11$
(4d)	$2D - (5p)$	$2P$	$2.5 - 1.5$	87.154			0.434	$3.818E + 11$
(5d)	$2D - (4f)$	$2F$	$1.5 - 2.5$	86.876			0.095	$8.422E + 10$
(5d)	$2D - (4f)$	$2F$	$2.5 - 3.5$	86.827			0.136	$1.205E + 11$
(4d)	$2D - (5p)$	$2P$	$1.5 - 1.5$	86.812			0.048	$4.292E + 10$
(5d)	$2D - (4f)$	$2F$	$2.5 - 2.5$	86.697			0.006	$6.053E + 09$
(5g)	$2G - (4f)$	$2F$	$3.5 - 3.5$	83.100			0.297	$2.871E + 11$
(5g)	$2G - (4f)$	$2F$	$4.5 - 3.5$	83.057	83.058		0.408	$1.006E + 13$
(5g)	$2G - (4f)$	$2F$	$3.5 - 2.5$	82.981	82.975		8.036	$7.785E + 12$
(5s)	$2S - (4p)$	$2P$	$0.5 - 1.5$	79.398			0.394	$4.178E + 11$
(5s)	$2S - (4p)$	$2P$	$0.5 - 0.5$	77.975			0.201	$2.205E + 11$
(4d)	$2D - (5f)$	$2F$	$2.5 - 2.5$	77.558			0.215	$2.384E + 11$
(4d)	$2D - (5f)$	$2F$	$2.5 - 3.5$	77.509	77.510		4.303	$4.777E + 12$
(4d)	$2D - (5f)$	$2F$	$1.5 - 2.5$	77.287	77.287		3.020	$3.373E + 12$
(5d)	$2D - (4p)$	$2P$	$1.5 - 1.5$	69.550			0.132	$1.828E + 11$
(5d)	$2D - (4p)$	$2P$	$2.5 - 1.5$	69.435	69.435		1.195	$1.653E + 12$
(5d)	$2D - (4p)$	$2P$	$1.5 - 0.5$	68.455			0.673	$9.584E + 11$
(4s)	$2S - (5p)$	$2P$	$0.5 - 0.5$	66.198			0.184	$2.807E + 11$
(4s)	$2S - (5p)$	$2P$	$0.5 - 1.5$	65.732			0.371	$5.735E + 11$
(6d)	$2D - (4f)$	$2F$	$1.5 - 2.5$	54.839			0.016	$3.568E + 10$
(6d)	$2D - (4f)$	$2F$	$2.5 - 3.5$	54.760			0.023	$5.118E + 10$
(6d)	$2D - (4f)$	$2F$	$2.5 - 2.5$	54.708			0.001	$2.567E + 09$
(6g)	$2G - (4f)$	$2F$	$3.5 - 3.5$	53.899			0.040	$9.398E + 10$
(6g)	$2G - (4f)$	$2F$	$4.5 - 3.5$	53.890			1.433	$3.291E + 12$
(6g)	$2G - (4f)$	$2F$	$3.5 - 2.5$	53.849			1.106	$2.545E + 12$
(4d)	$2D - (6p)$	$2P$	$1.5 - 0.5$	53.675			0.043	$9.999E + 10$
(4d)	$2D - (6p)$	$2P$	$2.5 - 1.5$	53.618			0.077	$1.806E + 11$
(4d)	$2D - (6p)$	$2P$	$1.5 - 1.5$	53.489			0.008	$2.021E + 10$
(4d)	$2D - (6f)$	$2F$	$2.5 - 2.5$	51.514			0.051	$1.299E + 11$
(4d)	$2D - (6f)$	$2F$	$2.5 - 3.5$	51.493			1.034	$2.601E + 12$
(4d)	$2D - (6f)$	$2F$	$1.5 - 2.5$	51.395			0.725	$1.831E + 12$
(6s)	$2S - (4p)$	$2P$	$0.5 - 1.5$	49.731			0.080	$2.159E + 11$
(6s)	$2S - (4p)$	$2P$	$0.5 - 0.5$	49.169			0.040	$1.117E + 11$
(6d)	$2D - (4p)$	$2P$	$1.5 - 1.5$	47.387			0.041	$1.238E + 11$
(6d)	$2D - (4p)$	$2P$	$2.5 - 1.5$	47.289			0.375	$1.121E + 12$
(6d)	$2D - (4p)$	$2P$	$1.5 - 0.5$	46.876			0.210	$6.393E + 11$
(4s)	$2S - (6p)$	$2P$	$0.5 - 0.5$	44.793			0.056	$1.874E + 11$

Transitions				λ	λ_1	λ_2	gf	gA (s^{-1})
(nl)	$^2L - (nl)$	2L	$J - J$	(calc.)	(obs.)	(obs.)		
(4s)	$2S - (6p)$	$2P$	$0.5 - 1.5$	44.663			0.113	$3.781E + 11$
(3d)	$2D - (4p)$	$2P$	$1.5 - 0.5$	39.929	39.929		0.104	$4.357E + 11$
(3d)	$2D - (4p)$	$2P$	$2.5 - 1.5$	39.728	39.728		0.188	$7.963E + 11$
(3d)	$2D - (4p)$	$2P$	$1.5 - 1.5$	39.565			0.021	$8.957E + 10$
(3d)	$2D - (4f)$	$2F$	$2.5 - 2.5$	35.665			0.269	$1.410E + 12$
(3d)	$2D - (4f)$	$2F$	$2.5 - 3.5$	35.643	35.643	35.643	5.383	$2.826E + 13$
(3d)	$2D - (4f)$	$2F$	$1.5 - 2.5$	35.534	35.534	35.534	3.779	$1.996E + 13$
(4s)	$2S - (3p)$	$2P$	$0.5 - 1.5$	34.917	34.916		0.233	$1.278E + 12$
(4s)	$2S - (3p)$	$2P$	$0.5 - 0.5$	34.211	34.213		0.119	$6.793E + 11$
(4d)	$2D - (3p)$	$2P$	$1.5 - 1.5$	30.926	30.926		0.148	$1.035E + 12$
(4d)	$2D - (3p)$	$2P$	$2.5 - 1.5$	30.885	30.885		1.337	$9.353E + 12$
(4d)	$2D - (3p)$	$2P$	$1.5 - 0.5$	30.373	30.373		0.755	$5.463E + 12$
(3s)	$2S - (4p)$	$2P$	$0.5 - 0.5$	28.899	28.900		0.172	$1.373E + 12$
(3s)	$2S - (4p)$	$2P$	$0.5 - 1.5$	28.709	28.709		0.346	$2.802E + 12$
(3d)	$2D - (5p)$	$2P$	$1.5 - 0.5$	25.905			0.017	$1.757E + 11$
(3d)	$2D - (5p)$	$2P$	$2.5 - 1.5$	25.902			0.031	$3.163E + 11$
(3d)	$2D - (5p)$	$2P$	$1.5 - 1.5$	25.833			0.003	$3.542E + 10$
(3d)	$2D - (5f)$	$2F$	$2.5 - 2.5$	24.984			0.048	$5.153E + 11$
(3d)	$2D - (5f)$	$2F$	$2.5 - 3.5$	24.979	24.979		0.964	$1.031E + 13$
(3d)	$2D - (5f)$	$2F$	$1.5 - 2.5$	24.919	24.920		0.676	$7.270E + 12$
(5s)	$2S - (3p)$	$2P$	$0.5 - 1.5$	23.249			0.047	$5.839E + 11$
(5s)	$2S - (3p)$	$2P$	$0.5 - 0.5$	22.934			0.024	$3.042E + 11$
(5d)	$2D - (3p)$	$2P$	$1.5 - 1.5$	22.323			0.043	$5.819E + 11$
(5d)	$2D - (3p)$	$2P$	$2.5 - 1.5$	22.311	22.312		0.391	$5.245E + 12$
(5d)	$2D - (3p)$	$2P$	$1.5 - 0.5$	22.033	22.033		0.220	$3.026E + 12$
(3d)	$2D - (6p)$	$2P$	$2.5 - 1.5$	21.842			0.011	$1.616E + 11$
(3d)	$2D - (6p)$	$2P$	$1.5 - 0.5$	21.824			0.006	$9.001E + 10$
(3d)	$2D - (6p)$	$2P$	$1.5 - 1.5$	21.793			0.001	$1.808E + 10$
(3d)	$2D - (6f)$	$2F$	$2.5 - 2.5$	21.485			0.017	$2.542E + 11$
(3d)	$2D - (6f)$	$2F$	$2.5 - 3.5$	21.481	21.481		0.351	$5.087E + 12$
(3d)	$2D - (6f)$	$2F$	$1.5 - 2.5$	21.437	21.438		0.246	$3.583E + 12$
(3s)	$2S - (5p)$	$2P$	$0.5 - 0.5$	20.764	20.764		0.051	$7.906E + 11$
(3s)	$2S - (5p)$	$2P$	$0.5 - 1.5$	20.718	20.718		0.102	$1.592E + 12$
(6s)	$2S - (3p)$	$2P$	$0.5 - 1.5$	19.792	19.790		0.018	$3.163E + 11$
(6s)	$2S - (3p)$	$2P$	$0.5 - 0.5$	19.563			0.009	$1.638E + 11$
(6d)	$2D - (3p)$	$2P$	$1.5 - 1.5$	19.409	19.410		0.019	$3.411E + 11$
(6d)	$2D - (3p)$	$2P$	$2.5 - 1.5$	19.393			0.173	$3.078E + 12$
(6d)	$2D - (3p)$	$2P$	$1.5 - 0.5$	19.189			0.097	$1.765E + 12$
(3s)	$2S - (6p)$	$2P$	$0.5 - 0.5$	18.057			0.023	$4.715E + 11$
(3s)	$2S - (6p)$	$2P$	$0.5 - 1.5$	18.036			0.046	$9.464E + 11$

Note : λ_1 — Observed wavelengths from reference [5]; λ_2 — Observed wavelengths from reference [6]; gf — Weighted oscillator strengths; gA — Weighted transition probability.

为便于比较, 表中同时列出了 Kononov 等人^[5]和 Reader 等人^[6]的波长观测值。由表可见,

我们的波长计算值与实验观测值十分符合,振子强度也与 Wiese 等人的结果一致。这主要是因为我们考虑了影响原子结构计算的两个主要效应:相对论效应和相关效应。在解 HF 方程中我们计入了主要的相对论效应修正(Darwin 修正和质量速度修正),在关于 Slater 能量参数的最小二乘优化过程中考虑了 HF 方法所无法计入的电子间的相关效应,因此得到了准确的计算结果。

从表中计算结果还可以看出,类钠储离子的 $5g \rightarrow 4f, 5f \rightarrow 4d, 6g \rightarrow 4f$ 及 $6f \rightarrow 4g$ 跃迁具有相当大的振子强度,上述跃迁的组态平均振子强度分别为 0.72, 0.38, 0.10 和 0.11,而下能级具有较大的辐射跃迁几率。例如 $4f \rightarrow 3d$ 跃迁几率为: $2.88 \times 10^{12} \text{ s}^{-1}$, 振子强度为 0.47, $4d \rightarrow 3p$ 的跃迁几率为 $1.13 \times 10^{12} \text{ s}^{-1}$, 振子强度为 0.16, 与其它跃迁相比,振子强度和跃迁几率都相对较大,因而有可能利用复合机制形成粒子数反转而产生上述跃迁的软 X 射线激光,。其中 $6g \rightarrow 4f, 6f \rightarrow 4d$ 跃迁的波长已接近“水窗”波段,相信这些结果在今后的实验中会得到进一步的证实。

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