Enhancement of second harmonic generation in LiNbO₃ crystals with polysynthetic domains

Hong Jingfen, Ming Naiben, Yang Yongshun, Zhu Jinsong, Yang Zhen, Wang Yening, Feng Duan

(Institute of Solid State Physics, Nanjing University)

The largest non-linear optical coefficient in LiNbO₃ is d_{33} (about 7.5 times larger than d_{31} which is the coefficient ordinarily used) which can not be phase-matched. During Czochralski growth, doping with about 1% of yttrium, and intentionally making the rotation axis noncoincident with the symmetry axis of the temperature field, so crystals with pronounced striae are grown, then cooling through Curie point, LiNbO₃ crystals with periodic ferro-electric domains in alternate signs are obtained. Growth axis is along a-axis which is perpendicular to the polar axis of crystals. In genera l_p , the width of the positive domains and l_n , those of negative domains, are unequal, however by adjusting rotation periods and pulling rates, so that $l_p + l_n \sim 2l_c$ may be achieved (l_c is the coherence length, $l_c \sim 3.4 \,\mu$ for $1.06 \,\mu$ fundamental wave).

Using the output of a continuously pumped acousto-optically Q-switched YAG laser as fundamental wave, we have measured the intensity of S. H. G. out of cavity due to quasi-phase-matching for b_{33} in crystals with polysynthetic domains. For comparison, we have also measured the intensity of S. H. G. in mono-domained crystals with same thickness and similar optical quality under angle phase-matched condition for d_{31} . Enhancement of S. H. G. output relative to ordinary phase-matched crystals is observed. The enhancement factors for some crystals are found in the range 1.5–12, while the maximum theoretical value is $(7.5 \times 2/\pi)^2 \sim 23$.

We also try to verify the relationship between the intensity of S. H. G. and number of domains, i.e. I $(2\omega)=I_o$ $(2\omega)N^2$ \cos^2 $(|l_p-l_n|\pi/2l_c)$. The results are shown in table 1. The deviations may be due to the irregularity of spacings and extra scattering of domain walls.

Table 1 $I(2\omega)/I_o(2\omega)$

N (No. of domains)	40	48	58	64	76
Theoretical value	1550	2240	3260	3960	5600
Experimental value	1500	1900	2200	3000	3200

聚片多畴 LiNbO3 晶体的倍频增强效应

洪静芬 闵乃本 杨永顺 朱劲松 杨 震 王业宁 冯 端

(南京大学固体物理研究所)

在 LiNbO₃ 晶体中最大的非线性光学系数为 d₃₃ (其数值为通用的d₃₁的7.5倍)是不能位相 匹配的。我们在提拉法生长中掺入 1 %左右的钇,并使籽晶转轴与温场对称轴相偏离,长出的 晶体具有明显的生长条纹,在通过居里点降温时诱发出正负交替的周期性聚片多畴 LiNbO₃晶体。生长轴沿 a 轴,极化矢量与之垂直。通常正畴厚度 l_0 和负畴厚度 l_n 不等,但可通过调节 转速和拉速使 $l_0+l_n\sim 2l_c$ (l_c 为相干长度,对应于 1.06μ 的基波 $l_c=3.4\mu$)。

应用连续泵浦声光调 Q的 YAG激光器输出的激光为基波,在腔外测量了聚片多畴LiNbO₃ 晶体由于 d_{33} 准位相匹配所产生的倍频光强。为了对比,我们也测量了厚度相同、光学质量相仿的单畴 $LiNbO_3$ 晶体在 d_{31} 角度位相匹配条件下的倍频光强。观测到了准位相匹配相对于正常位相匹配的倍频增强效应,实测到的增强因子在 $1.5\sim12$ 之间,增强因子的最大理论值为 $(7.5\times\frac{2}{\pi})^2\sim23$ 。

我们也对聚片多畴晶体的倍频光强和片数N的关系式 $I(2\omega)=I_0(2\omega)N^2\cos^2(\frac{\pi|l_p-l_n|}{2l_o})$ 进行了检验,结果见表 1,偏差是随片数增长的,可能是畴间距不规则及畴界 散射 所 引起的。

表 1 $I(2\omega)/I_0(2\omega)$ 的实验值和理论值

片	数	N	40	48	58	64	76
理	论	值	1550	2240	3260	3980	5600
实	验	值	1500	1900	2200	3000	3200